

# Real-time non-perturbative dynamics in Schwinger model: jet production, chiral magnetic wave, and more

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## References:

2301.11991: Florio, Frenklakh, [Ikeda](#), Kharzeev, Korepin, SS, Yu

2305.00996: [Ikeda](#), Kharzeev, Meyer, SS

2305.05685: [Ikeda](#), Kharzeev, SS

(池田一毅)

- motivation: real time dynamics in QFT
- model set up
- jet production
  - jet production
  - vector and axial charge transport
  - phase structure
- summary and outlook

“first principle” microscopic theory: quantum field theory  $L \leftrightarrow H$

perturbative calculation: scattering process, thermodynamics, transport

non-perturbative calculation(lattice QFT): thermodynamics, transport

*real time dynamics of non-perturbative theory?*

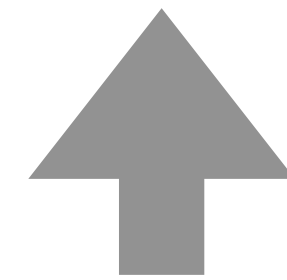
time evolution of a quantum (field) state:

$$|\psi(t)\rangle = \mathcal{T} e^{-i \int \hat{H} dt} |\psi(0)\rangle$$

Ideally, *quantum simulation* for *full QCD in 3+1 D*, but ...

1+1D Schwinger model with time-dependent external source

$$H(t) = \int \left( \frac{E^2}{2} - \bar{\psi}(i\gamma^1 \partial_z - g\gamma^1 A - m)\psi \right) dx.$$



$E$ : electric field

$A$ : electric potential

$\psi, \bar{\psi}$ : fermion field

$$L(t) = \int \left( -\frac{F^{\mu\nu} F_{\mu\nu}}{4} + \bar{\psi}(i\gamma^\mu \partial_\mu - g\gamma^\mu A_\mu - m)\psi \right) dx.$$

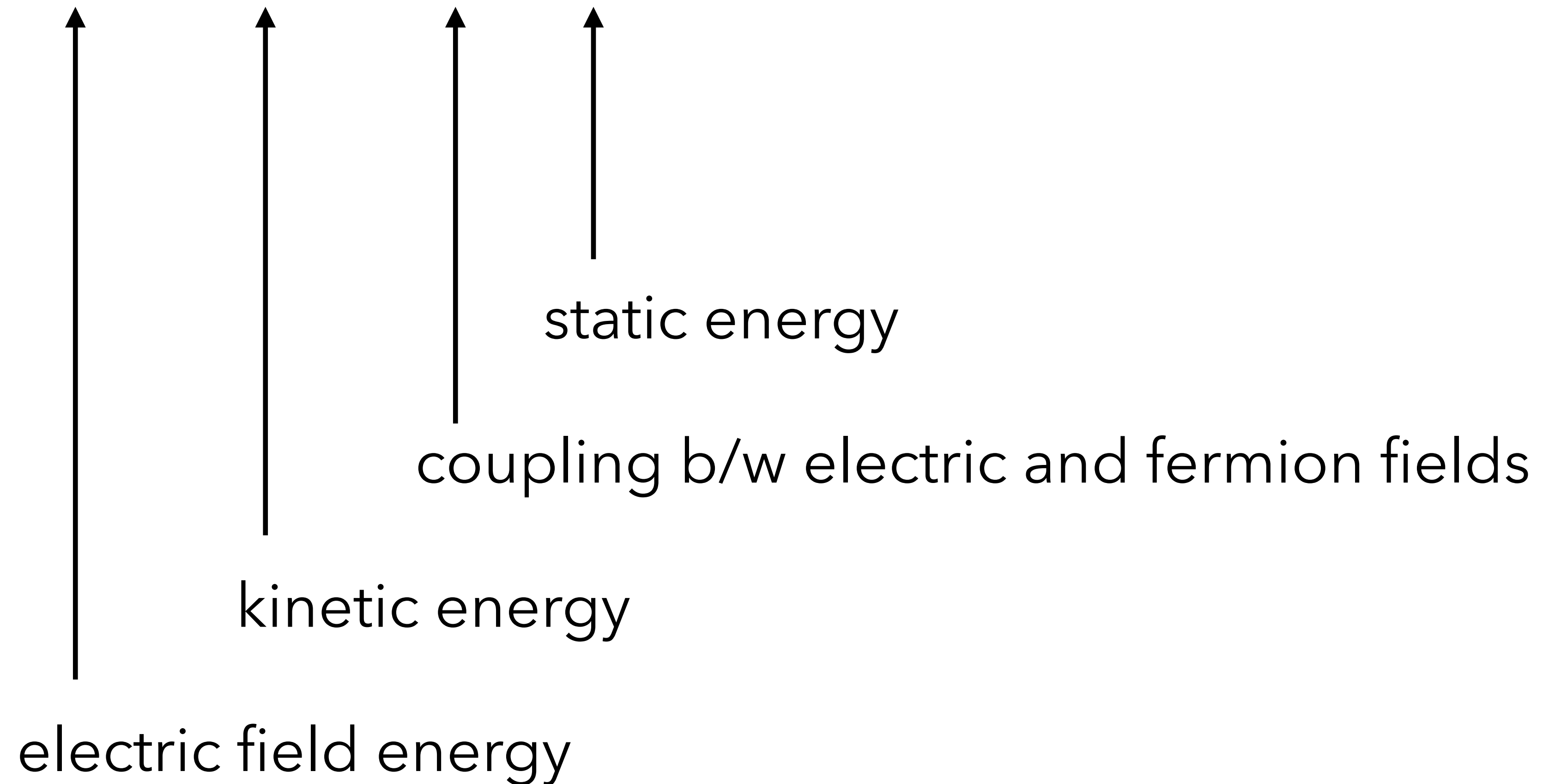
1+1D Schwinger model with time-dependent external source

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1+1D Schwinger model with time-dependent external source

$$H(t) = \int \left( \frac{E^2}{2} - \bar{\psi}(i\gamma^1 \partial_z - g\gamma^1 A - m)\psi - j_{\text{ext}}^1(t)A \right) dx.$$

$E$ : electric field  
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$$L(t) = \int \left( -\frac{F^{\mu\nu}F_{\mu\nu}}{4} + \bar{\psi}(i\gamma^\mu \partial_\mu - g\gamma^\mu A_\mu - m)\psi - j_{\text{ext}}^\mu(t)A_\mu \right) dx.$$

coupling w/ external source (jets)

$$j_{\text{ext}}^1(x, t) = g [\delta(x - t) + \delta(x + t)] \theta(t)$$

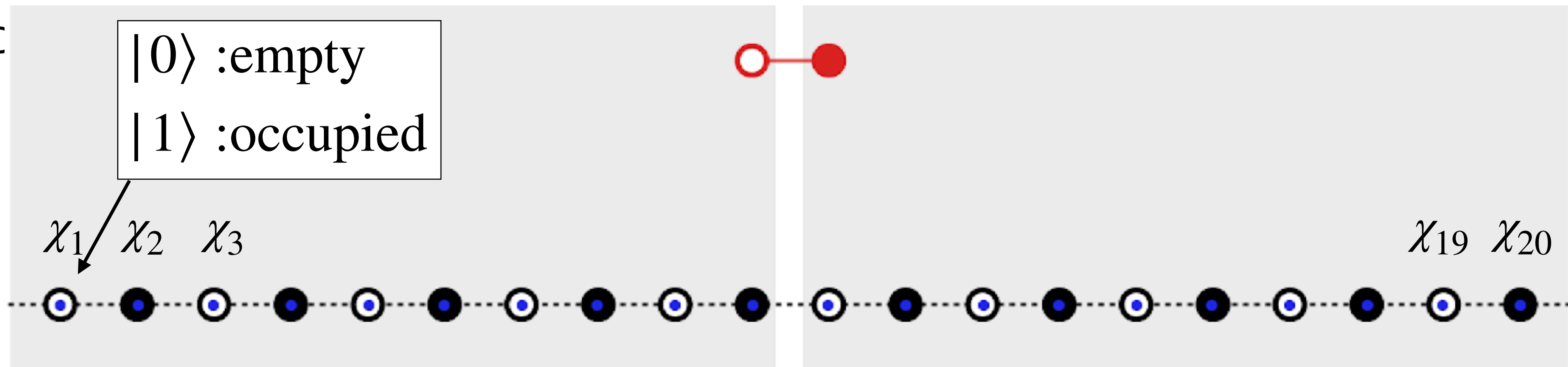
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discretize and matrix(gate) representation:



1+1D Sc



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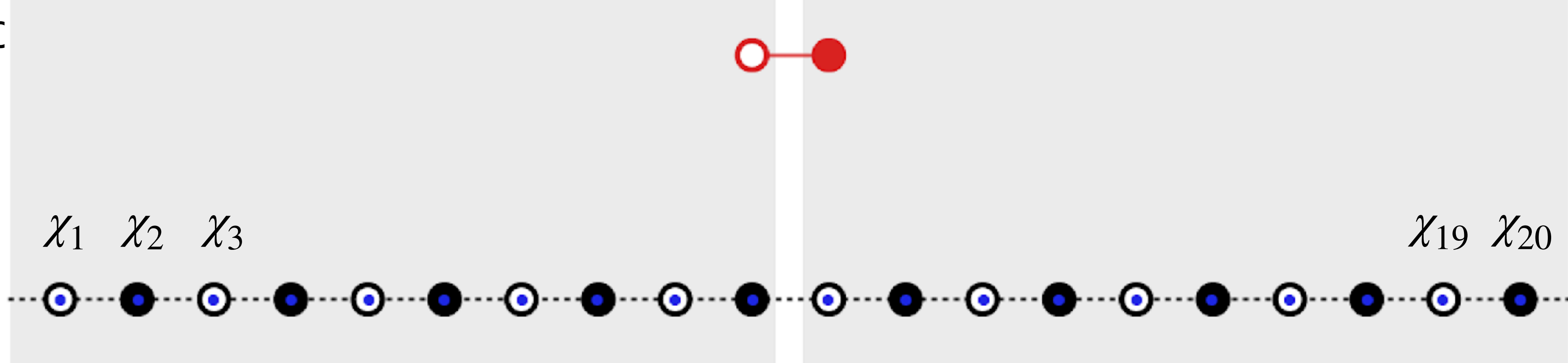
staggered fermion that satisfied anti-commutation:  $\{\psi_a(x), \psi_b^\dagger(y)\} = \delta_{a,b}\delta(x - y)$

$$\psi(x = a n) \quad \leftrightarrow \quad \frac{1}{\sqrt{a}} \begin{pmatrix} \chi_{2n} \\ \chi_{2n-1} \end{pmatrix}$$

Kogut-Susskind



1+1D Sc



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Pauli matrices:  $X, Y, Z$

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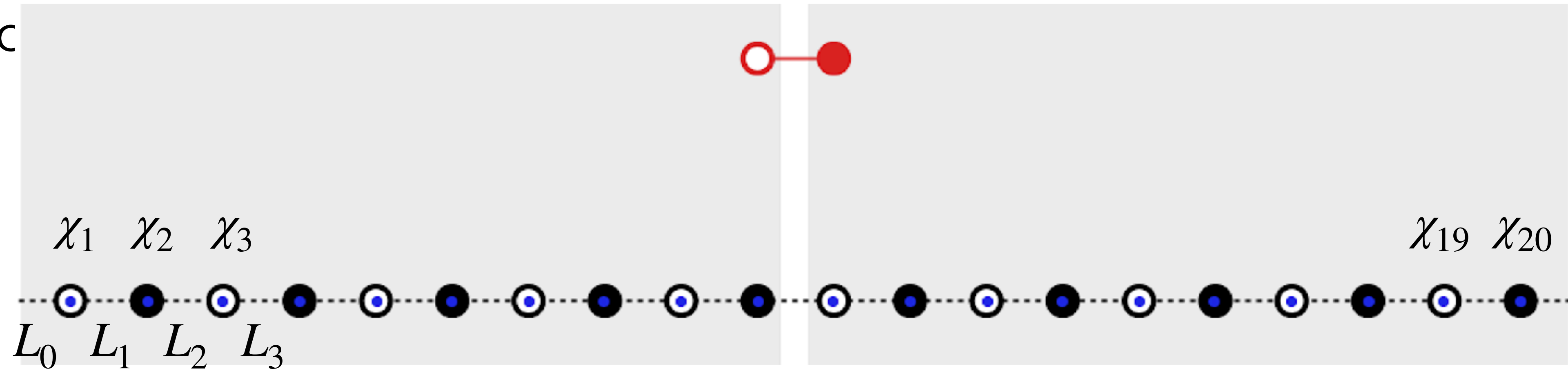
Kogut-Susskind

$$\chi_n = \frac{X_n - iY_n}{2} \prod_{m=1}^{n-1} (-iZ_m)$$

Jordan-Wigner

$$\{\chi_n^\dagger, \chi_m\} = \delta_{nm}, \quad \{\chi_n^\dagger, \chi_m^\dagger\} = \{\chi_n, \chi_m\} = 0.$$

1+1D Sc



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gauge field fixed by Gauss' law:  $\partial_1 E - g \bar{\psi} \gamma^0 \psi = j_{\text{ext}}^0$

$$E(x = an) \quad \leftrightarrow \quad L_n \quad L_n - L_{n-1} - \frac{Z_n + (-1)^n}{2} = \frac{1}{g} \int_{(n-1/2)a}^{(n+1/2)a} dx j_{\text{ext}}^0(x, t) ,$$

1+1D Schwinger model with time-dependent external source

$$H(t) = \int \left( \frac{E^2}{2} - \bar{\psi}(i\gamma^1 \partial_z - g\gamma^1 A - m)\psi - j_{\text{ext}}^1(t)A \right) dx.$$

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$$H(t) = \frac{1}{4a} \sum_{n=1}^{N-1} (X_n X_{n+1} + Y_n Y_{n+1}) + \frac{m}{2} \sum_{n=1}^N (-1)^n Z_n + \frac{a g^2}{2} \sum_{n=1}^{N-1} L_n^2(t).$$

# why quantum computer?

dimension of state vector =  $2^N$

$N$  : number of lattice sides

dimension of Hamiltonian =  ~~$2^N \times 2^N$~~  sparse  $\sim 2N \times 2^N$

$N$	dimension	memory of Hamiltonian	# of qubit (N)
8	256	~ 131 kB	8
12	4,096	~ 3.1 MB	12
16	65,536	~ 67 MB	16
20	1,048,576	~ 1.3 GB	20
24	16,777,216	~ 26 GB	24
28	268,435,456	~ 481 GB	28

unrealistic in a "classical" computer,  
but plausible in the state-of-art quantum computer?

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28	<del>268,435,456</del>	<del>~ 481 GB</del>	28

performance not satisfying...



1+1D Schwinger model with time-dependent external source

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time-dependent Schroedinger equation:

$$\frac{\partial}{\partial t} |\psi(t)\rangle = -i H |\psi(t)\rangle$$

$$q_{n,t} \equiv \langle \psi^\dagger(a n) \psi(a n) \rangle_t = \frac{\langle Z_n \rangle_t + (-1)^n}{2a},$$

$$Q_t \equiv \int \langle \psi^\dagger(x) \psi(x) \rangle_t dx = a \sum_{n=1}^N q_{n,t},$$

$$\nu_{n,t} \equiv \langle \bar{\psi}(a n) \psi(a n) \rangle_t = \frac{(-1)^n \langle Z_n \rangle_t}{2a},$$

$$\nu_t \equiv \int \langle \bar{\psi}(x) \psi(x) \rangle_t dx = a \sum_{n=1}^N \nu_{n,t},$$

$$\Pi_{n,t} \equiv \langle E(a n) \rangle_t = g \langle L_n \rangle_t,$$

$$E_{\text{ele},t} \equiv \frac{1}{2} \int \langle E^2(x) \rangle_t dx = \frac{a g^2}{2} \sum_{n=1}^{N-1} \langle L_n^2 \rangle_t.$$

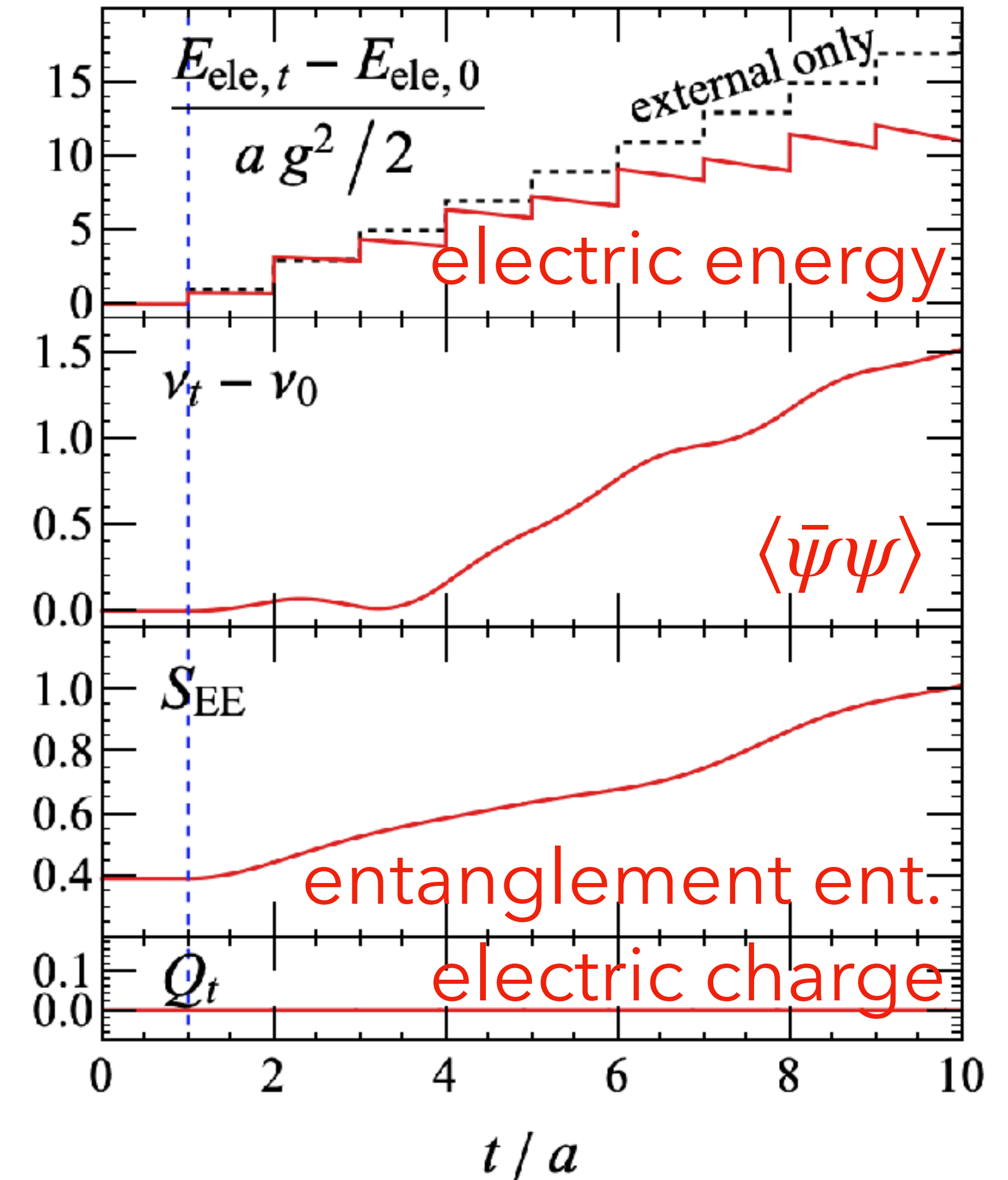
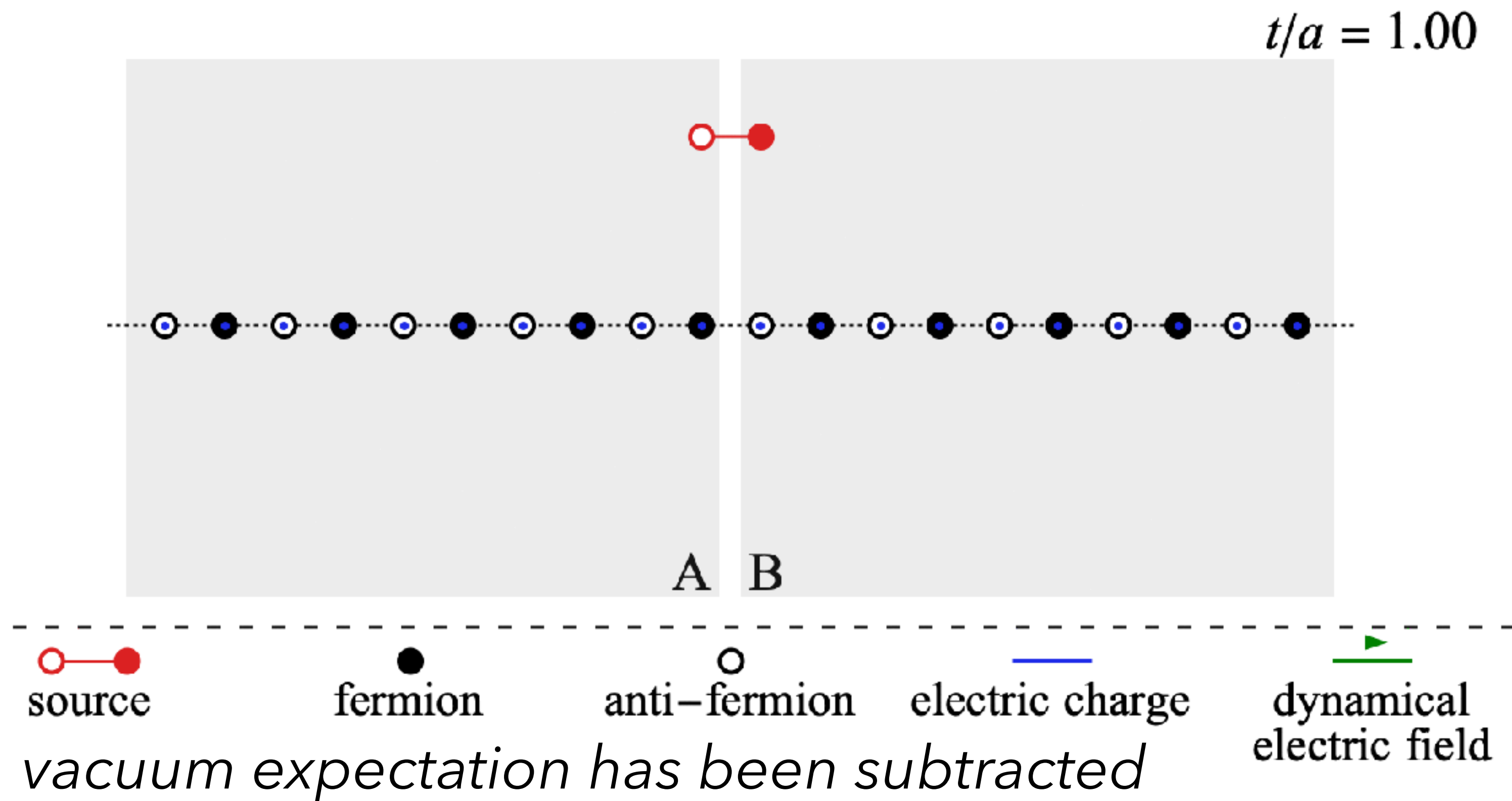
# *I. jet production*

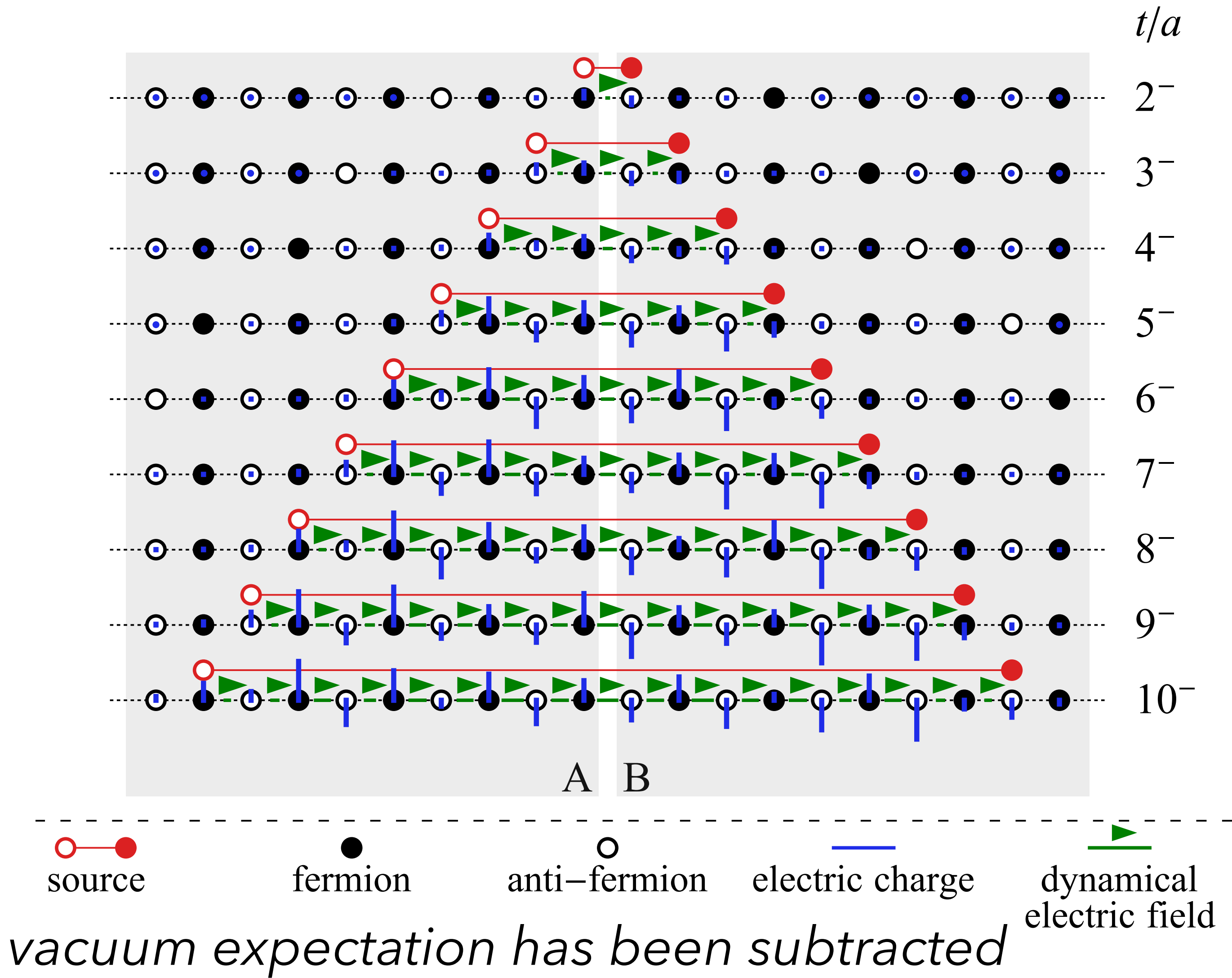
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A. Florio, D. Frenklakh, K. Ikeda, D. Kharzeev, V. Korepin, SS, K. Yu  
PhysRevLett.131.021902 (arXiv: 2305.05685)

initial state: vacuum  $H(t = 0) |\psi(t = 0)\rangle = E_0 |\psi(t = 0)\rangle$

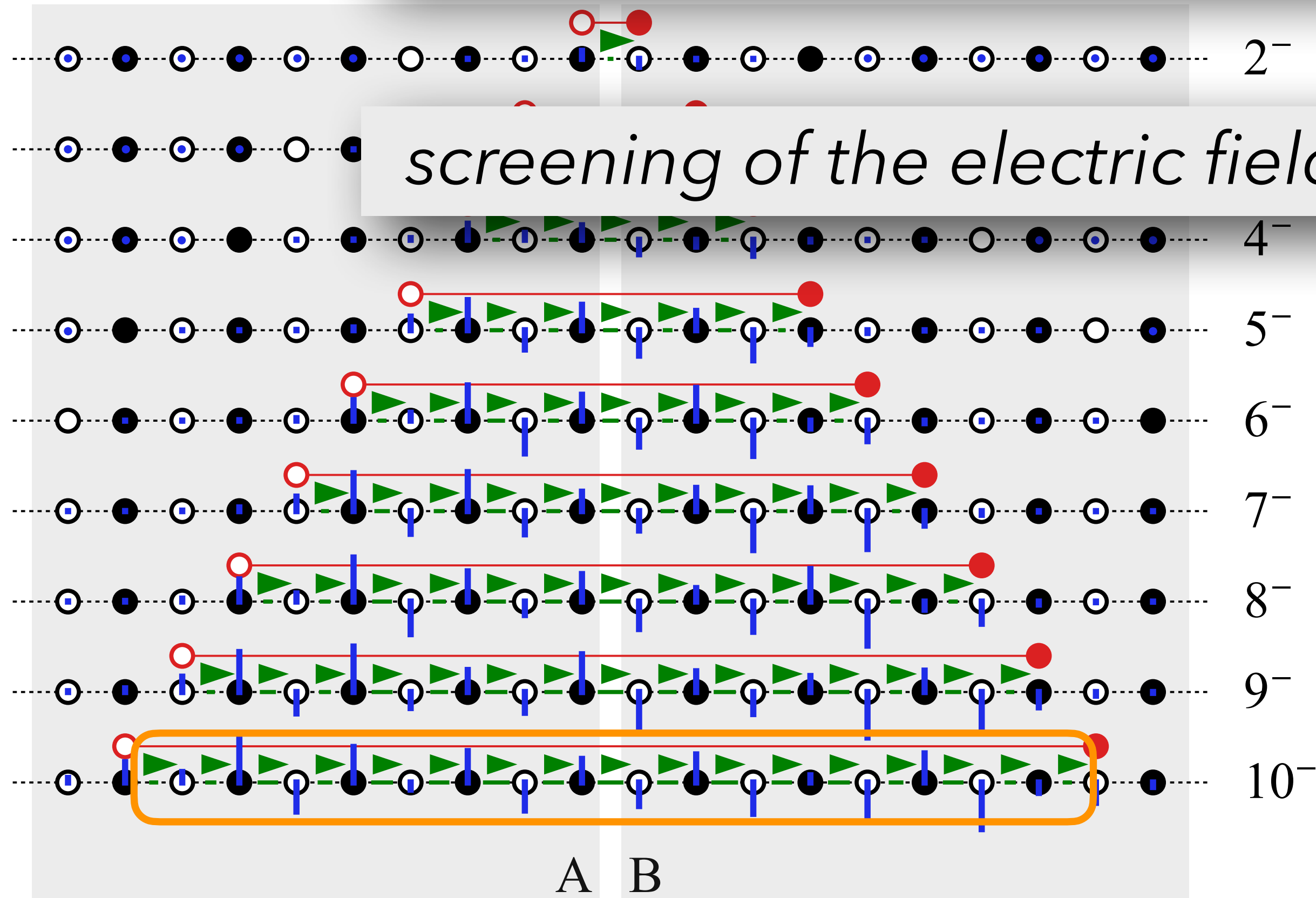
$$\frac{\partial}{\partial t} |\psi(t)\rangle = -i H(t) |\psi(t)\rangle$$





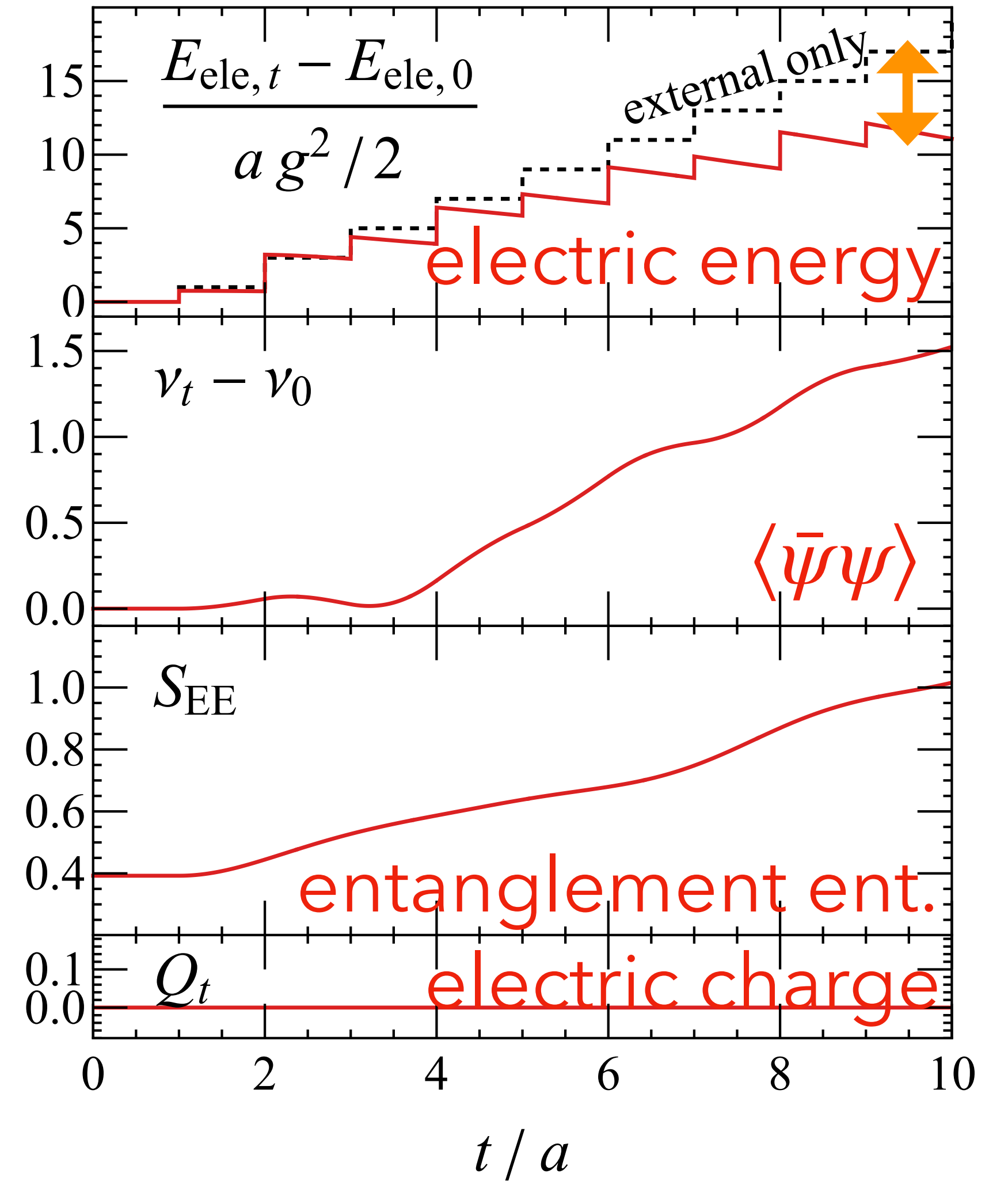
effects of pair production:

screening of the electric field

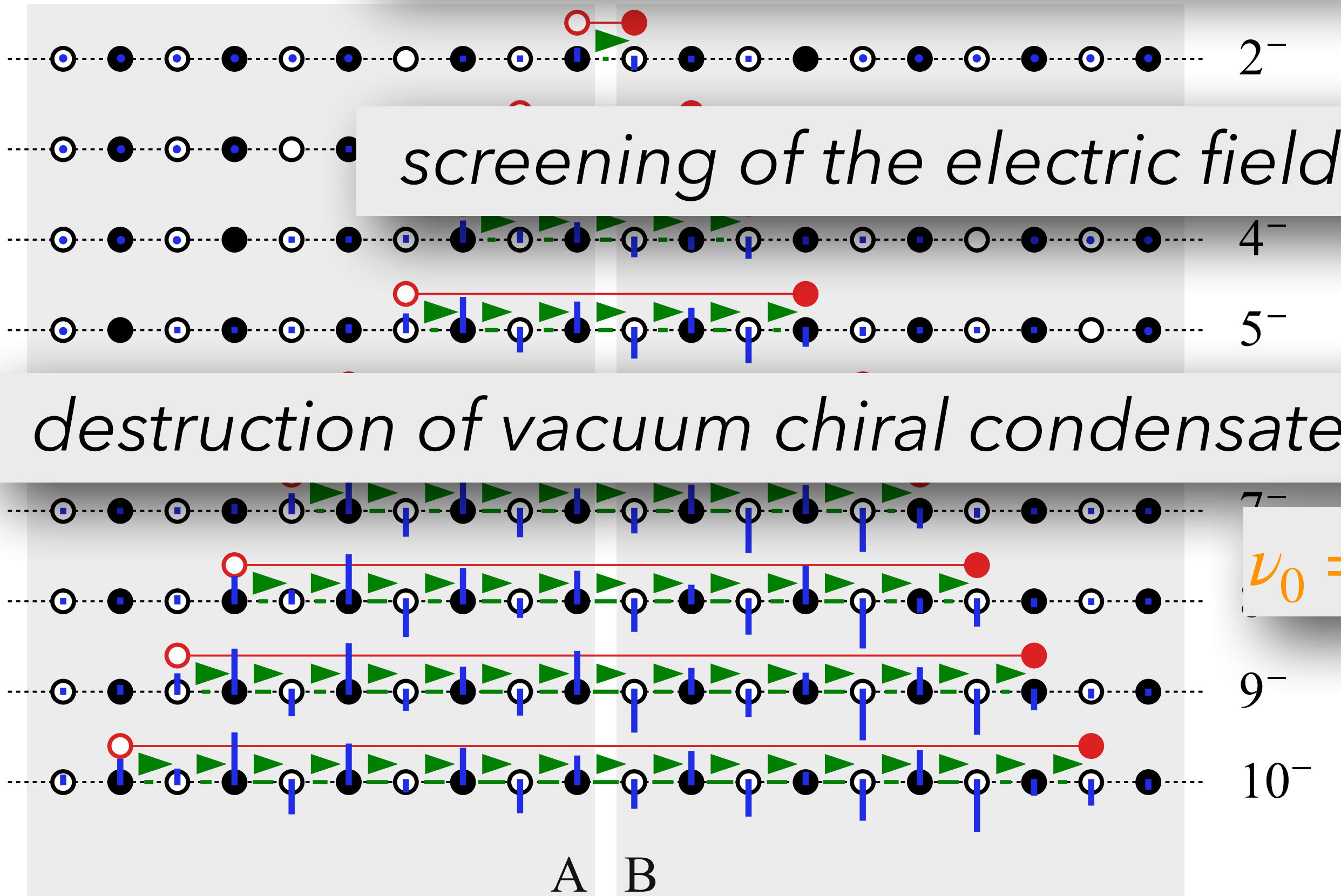


source  
 fermion  
 anti-fermion  
 electric charge  
 dynamical electric field

vacuum expectation has been subtracted



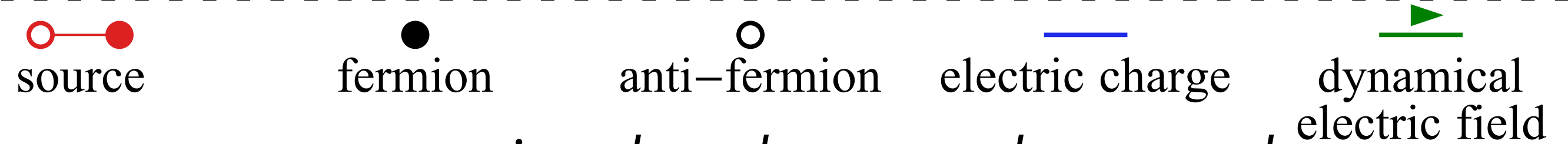
effects of pair production:



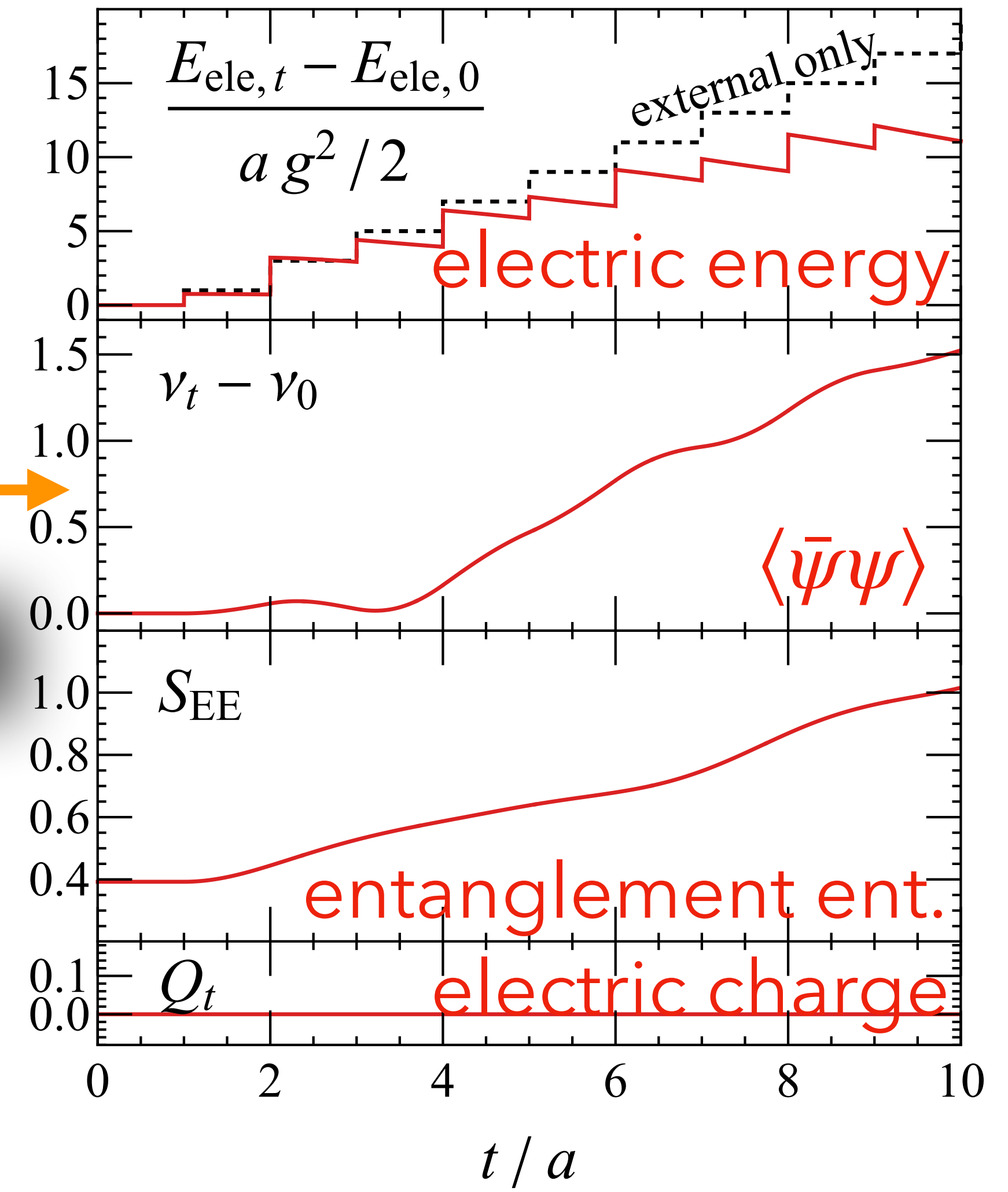
screening of the electric field

destruction of vacuum chiral condensate

$\nu_0 = -5.16$

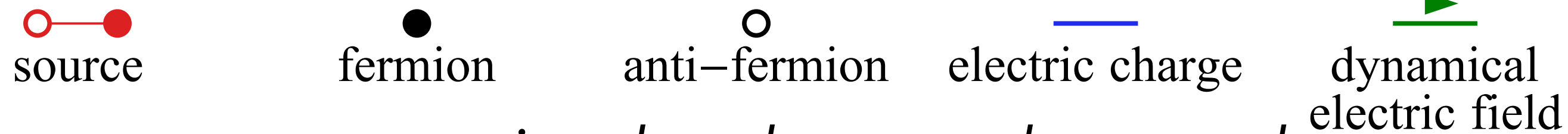
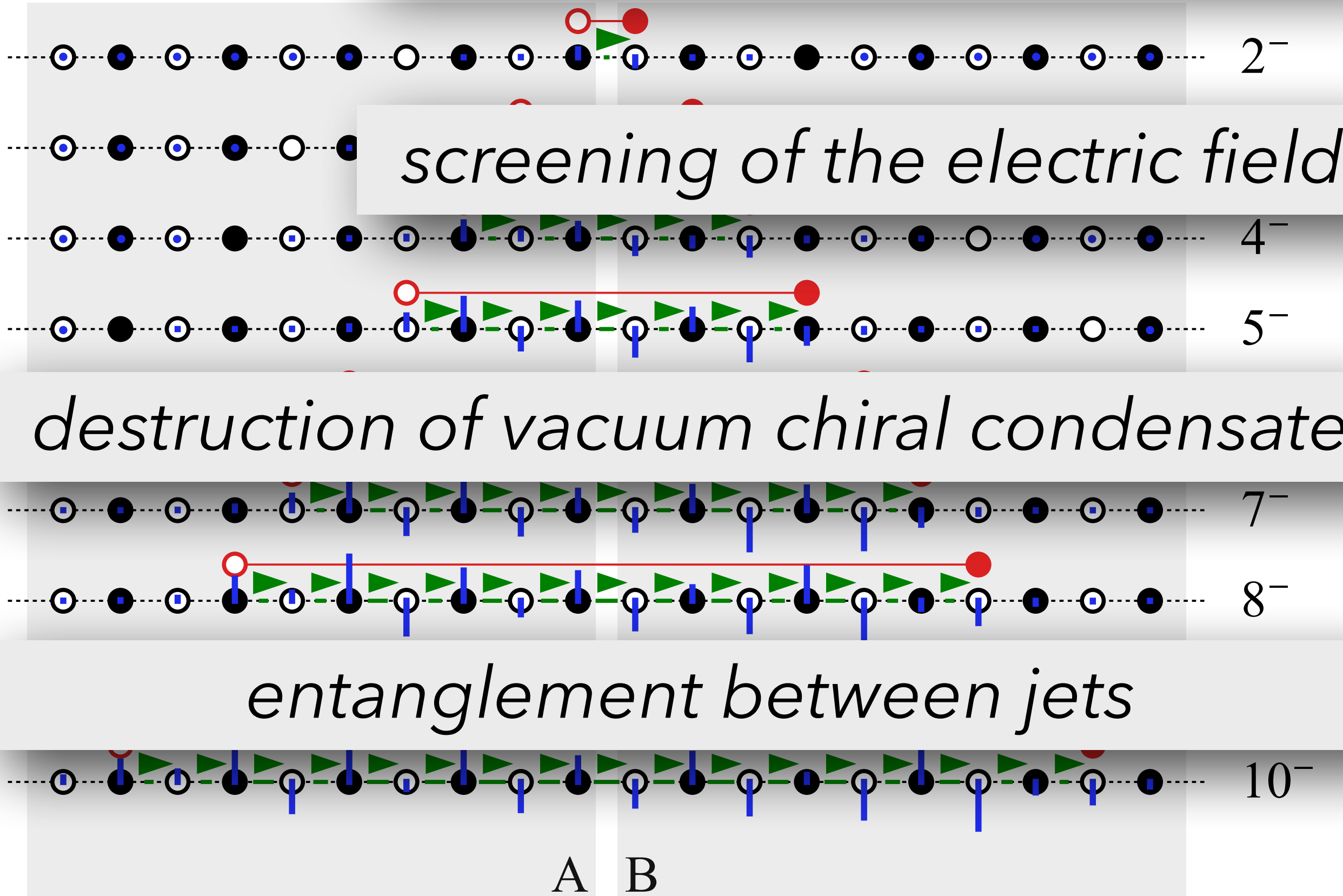


vacuum expectation has been subtracted

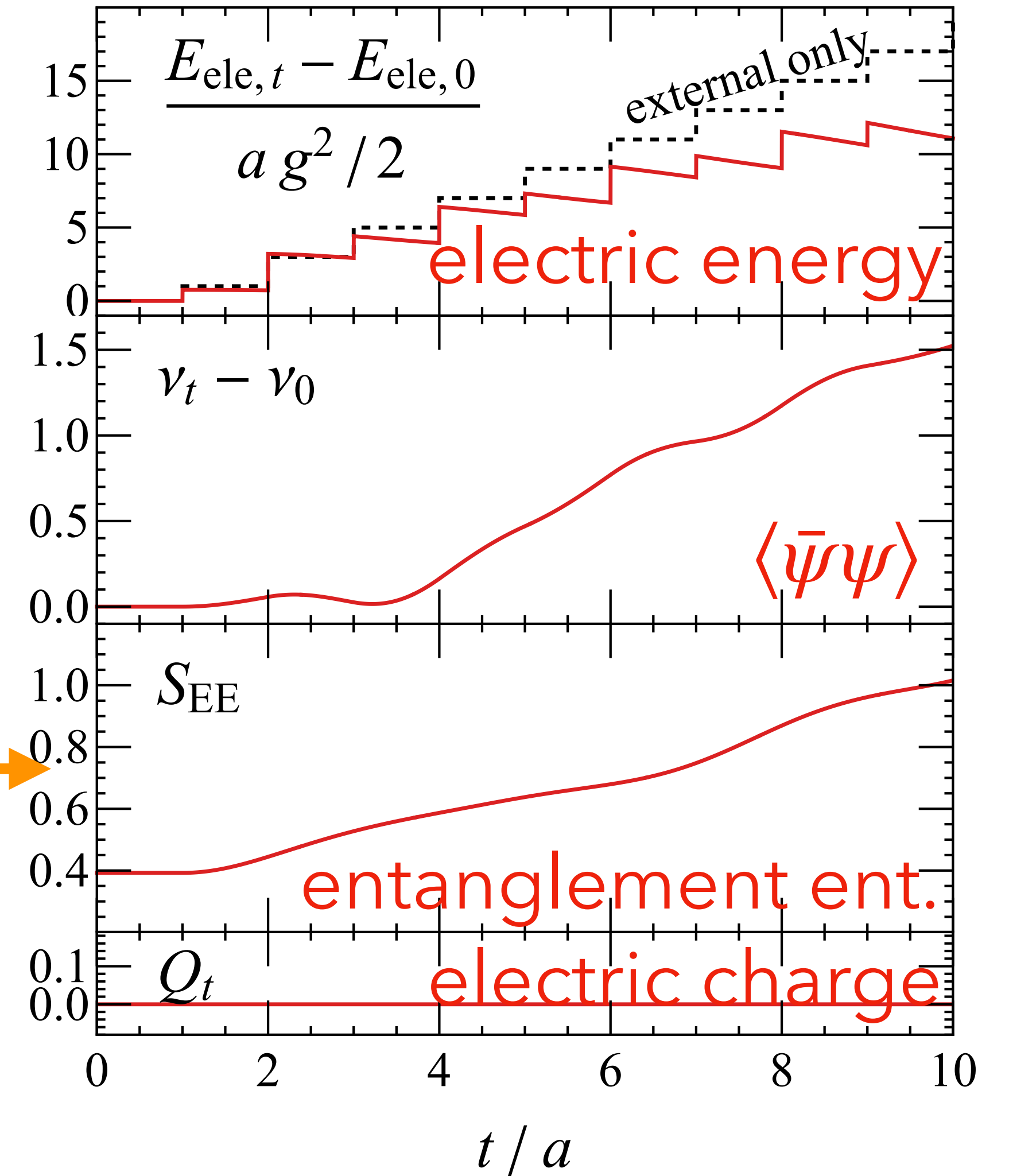




effects of pair production:

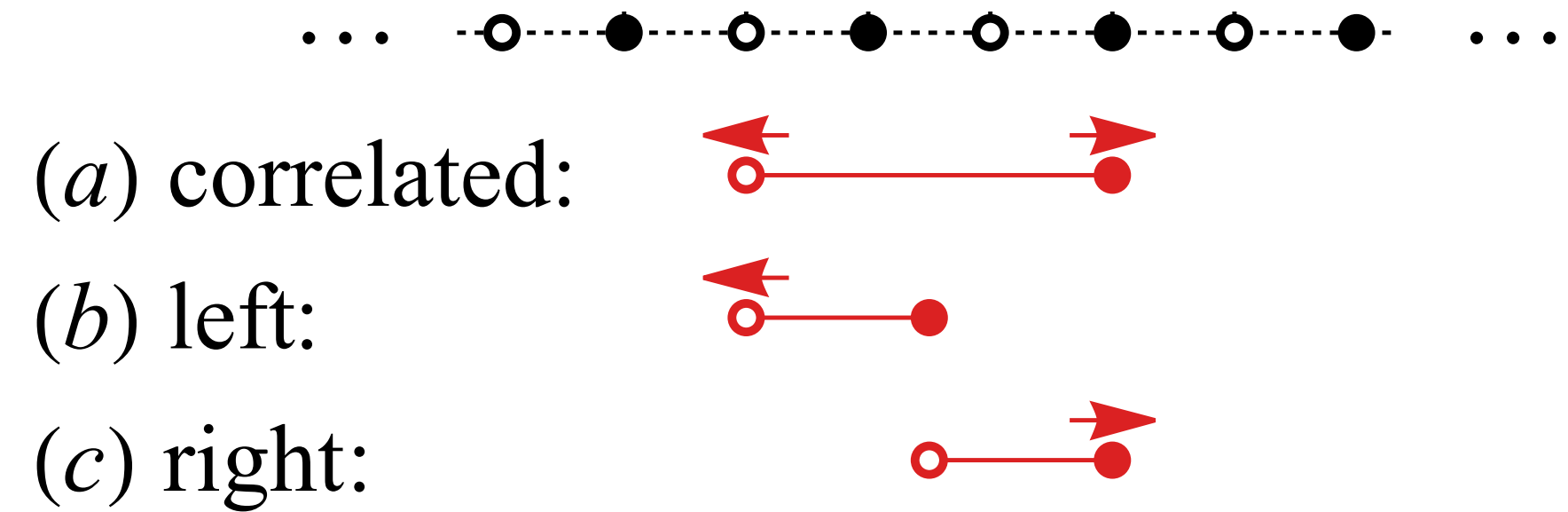


vacuum expectation has been subtracted

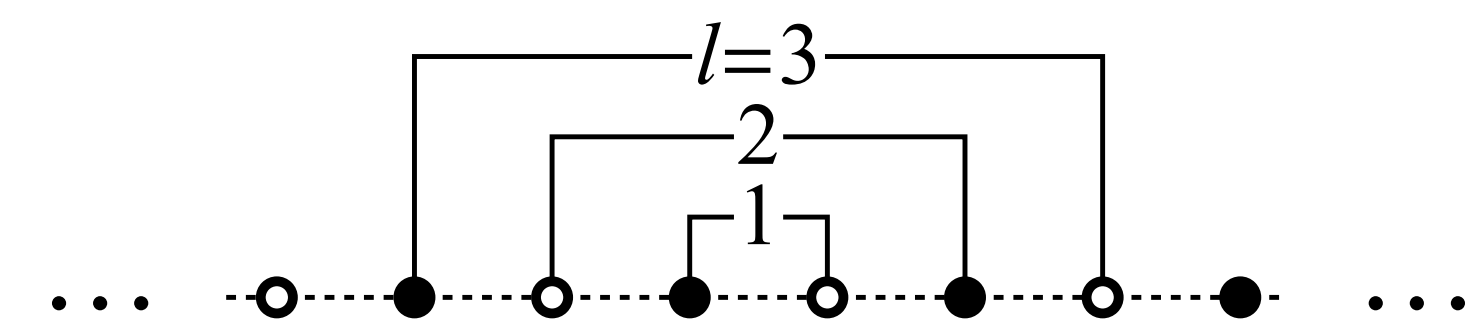
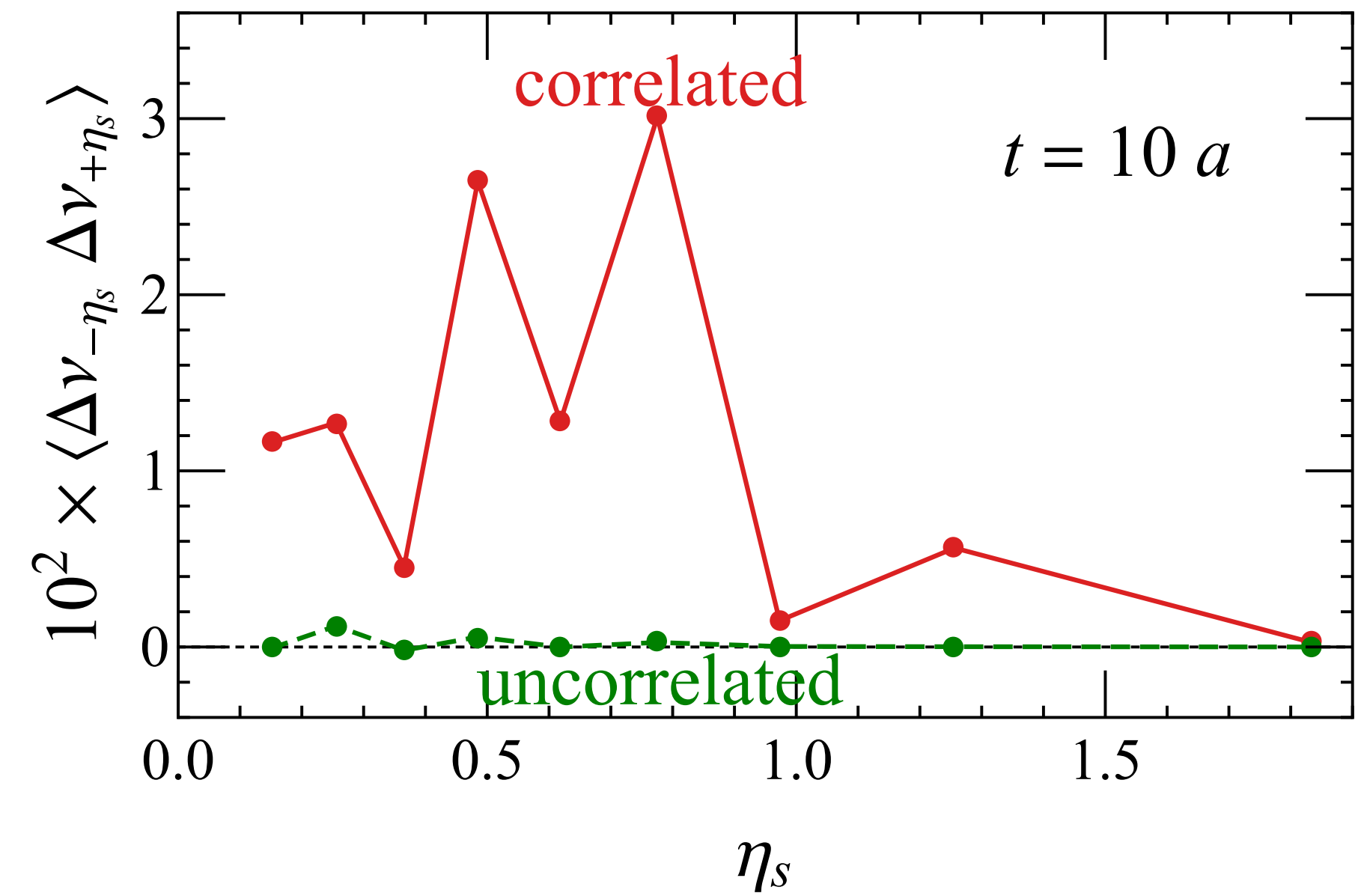
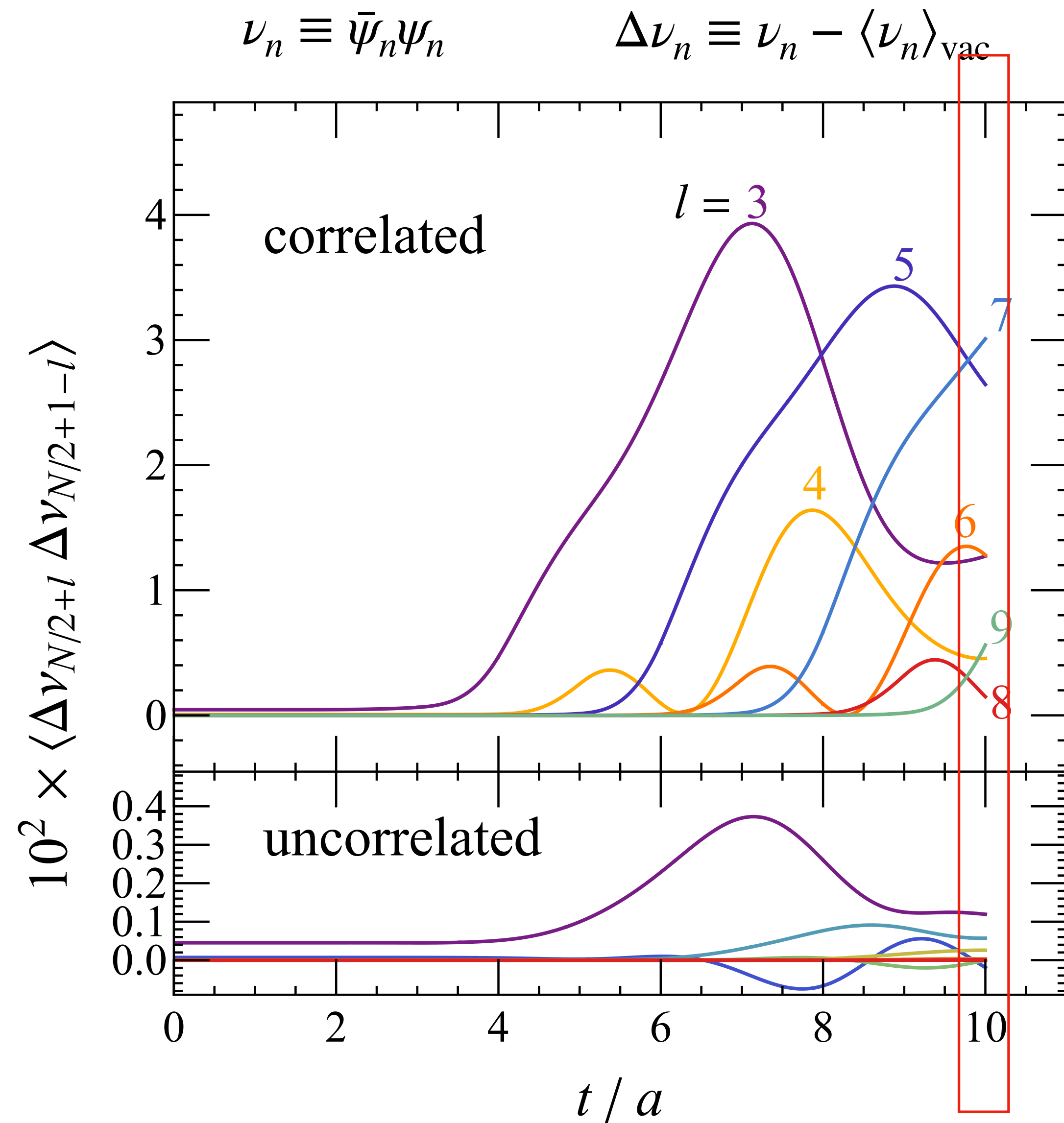


$$|\psi_{\text{uncorr}}\rangle = \frac{1}{\sqrt{2}} |\psi_{\text{left}}\rangle + \frac{e^{i\varphi}}{\sqrt{2}} |\psi_{\text{right}}\rangle$$

$$\begin{aligned} & \langle\langle \psi_{\text{uncorr}} | O | \psi_{\text{uncorr}} \rangle\rangle \\ \equiv & \int \langle \psi_{\text{uncorr}} | O | \psi_{\text{uncorr}} \rangle \frac{d\varphi}{2\pi} \\ = & \frac{\langle \psi_{\text{left}} | O | \psi_{\text{left}} \rangle}{2} + \frac{\langle \psi_{\text{right}} | O | \psi_{\text{right}} \rangle}{2} \end{aligned}$$







(a) correlated:



(b) left:



(c) right:



## *II. Vector & Axial Charge Transport*

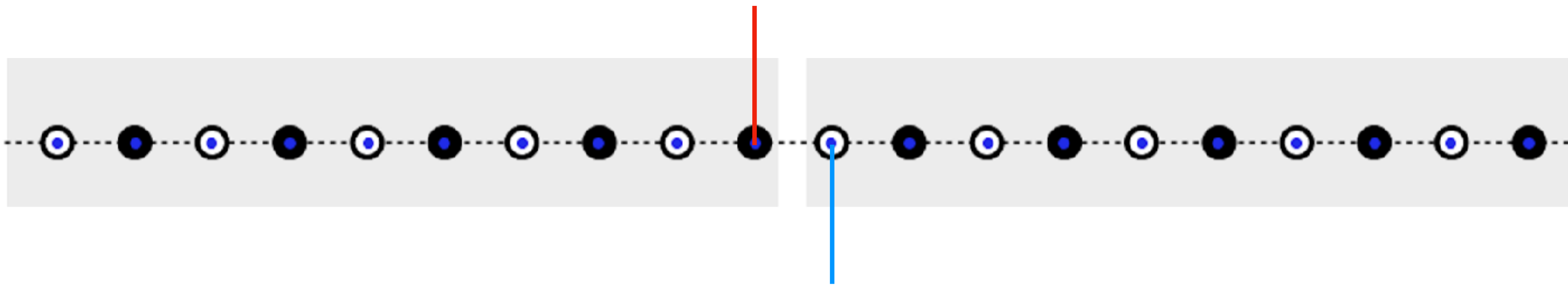
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K. Ikeda, D. Kharzeev, S. Shi,  
Phys. Rev. D in press(arXiv: 2305.05685)

$$\partial_t Q + \partial_z J = 0$$

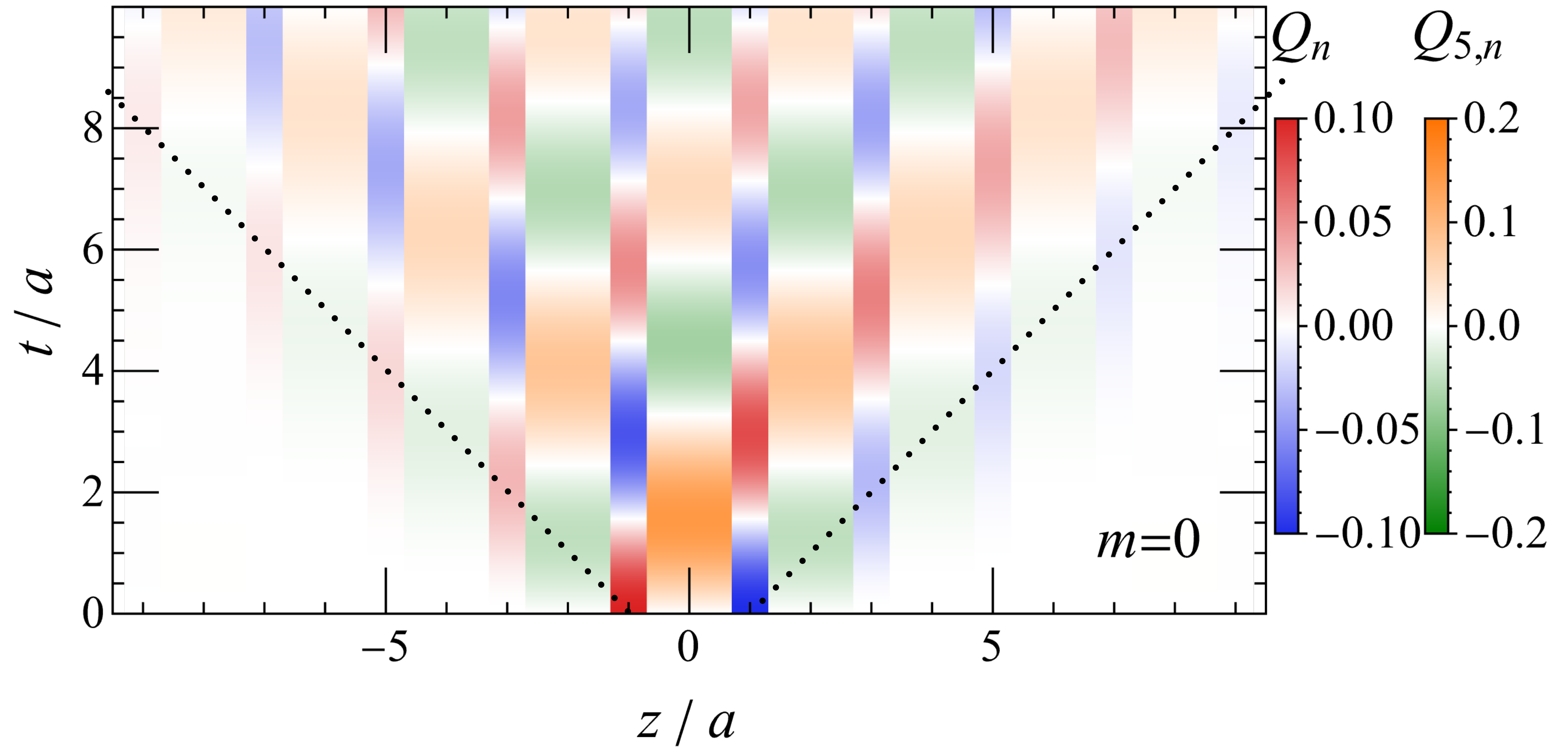
in Schwinger model:

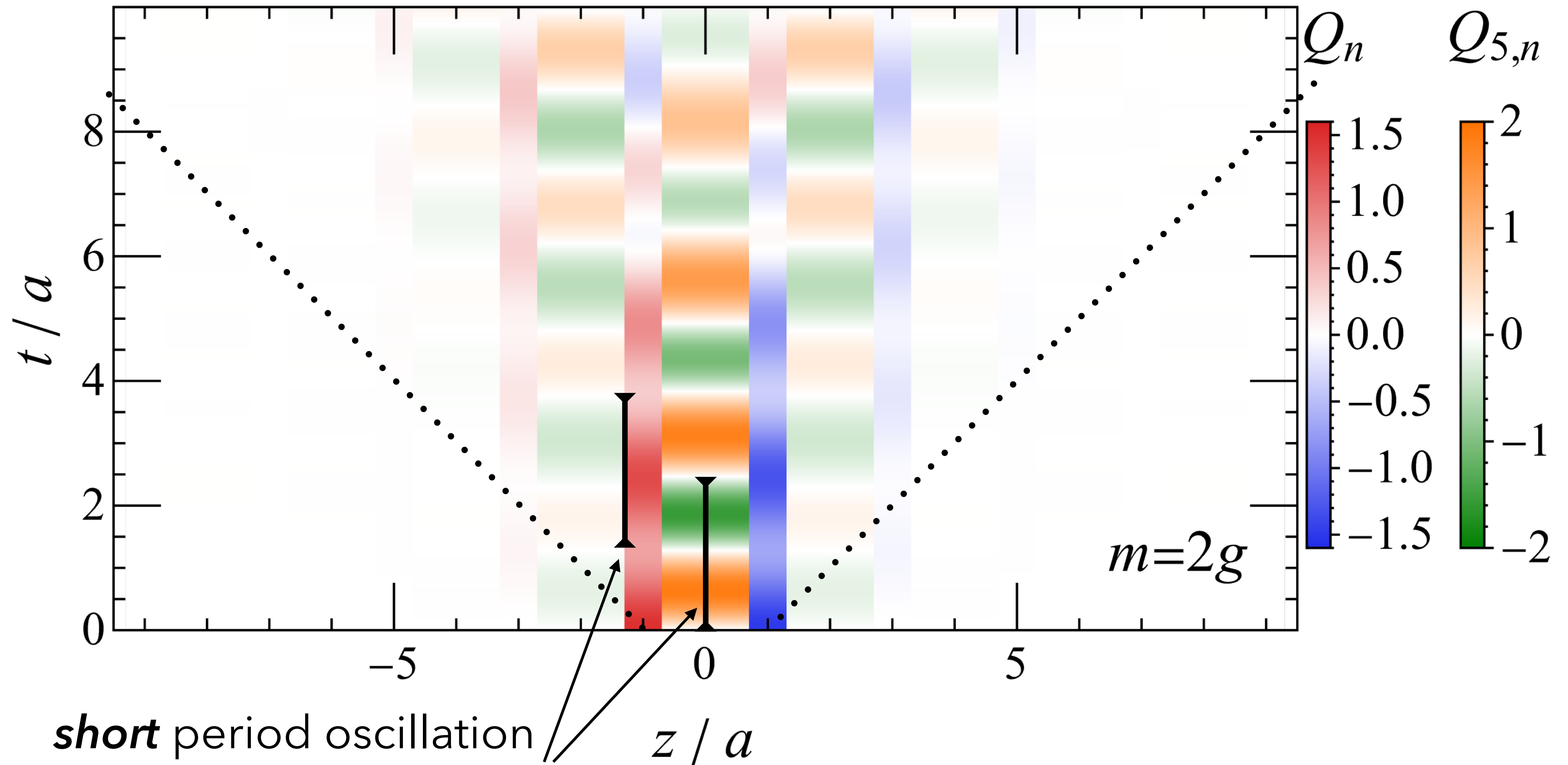
$$Q = J_5, \quad J = -Q_5$$



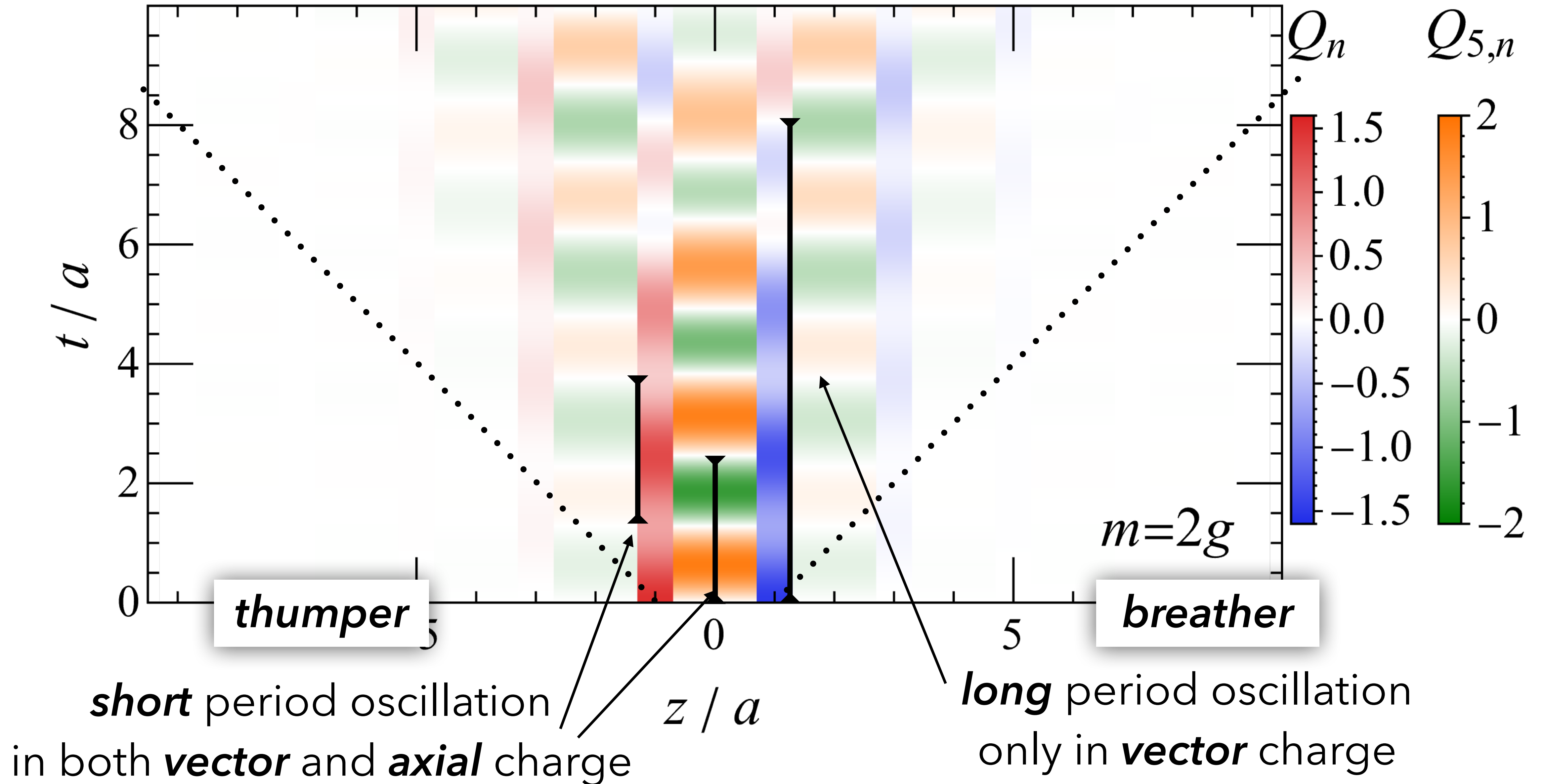
$$Q_n \equiv \langle \bar{\psi}(a n) \gamma^0 \psi(a n) \rangle = \frac{\langle Z_n \rangle + (-1)^n}{2a},$$

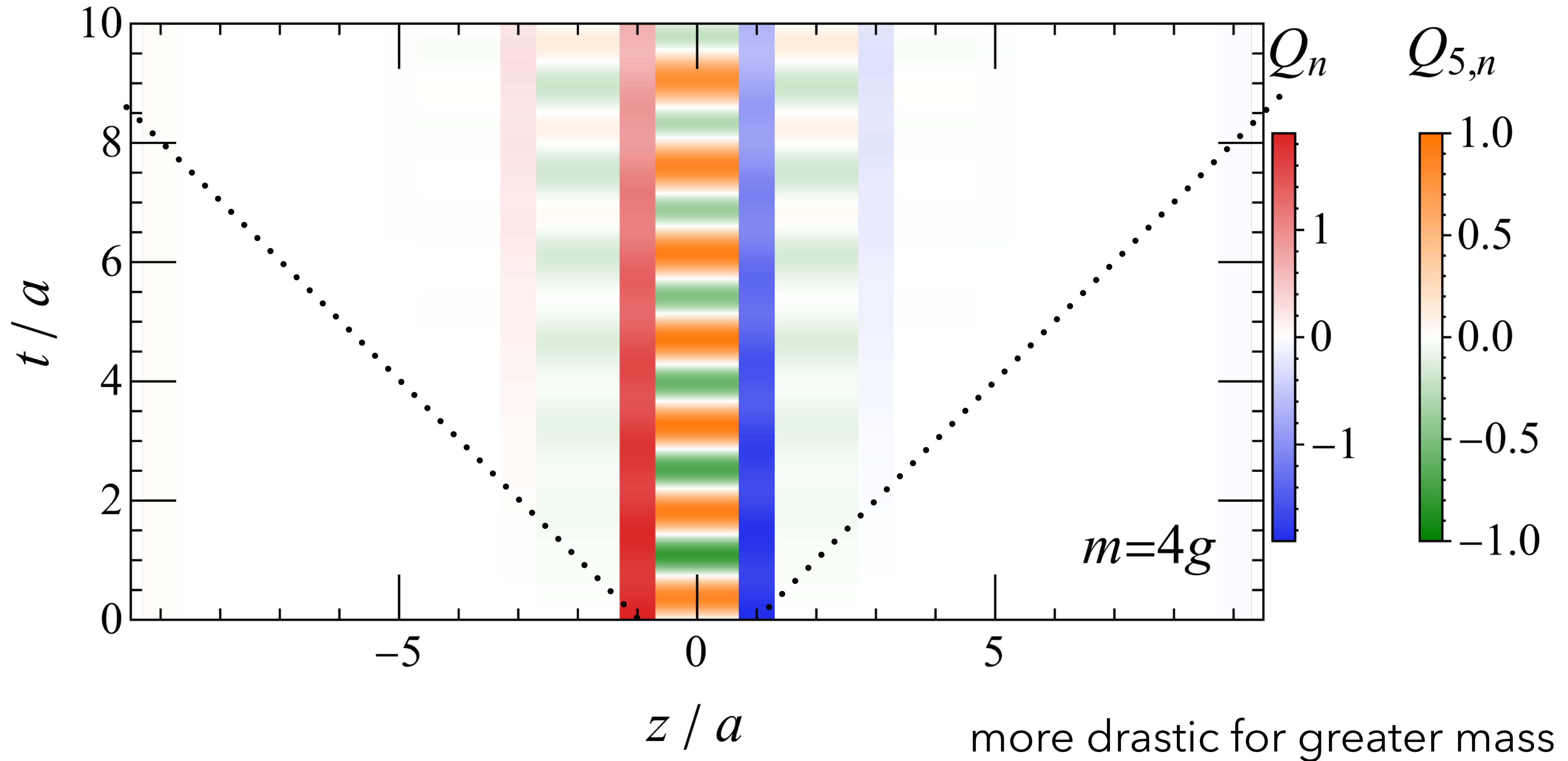
$$Q_{5,n} \equiv \langle \bar{\psi}(a n) \gamma^5 \gamma^0 \psi(a n) \rangle = \frac{\langle X_n Y_{n+1} - Y_n X_{n+1} \rangle}{4a}$$

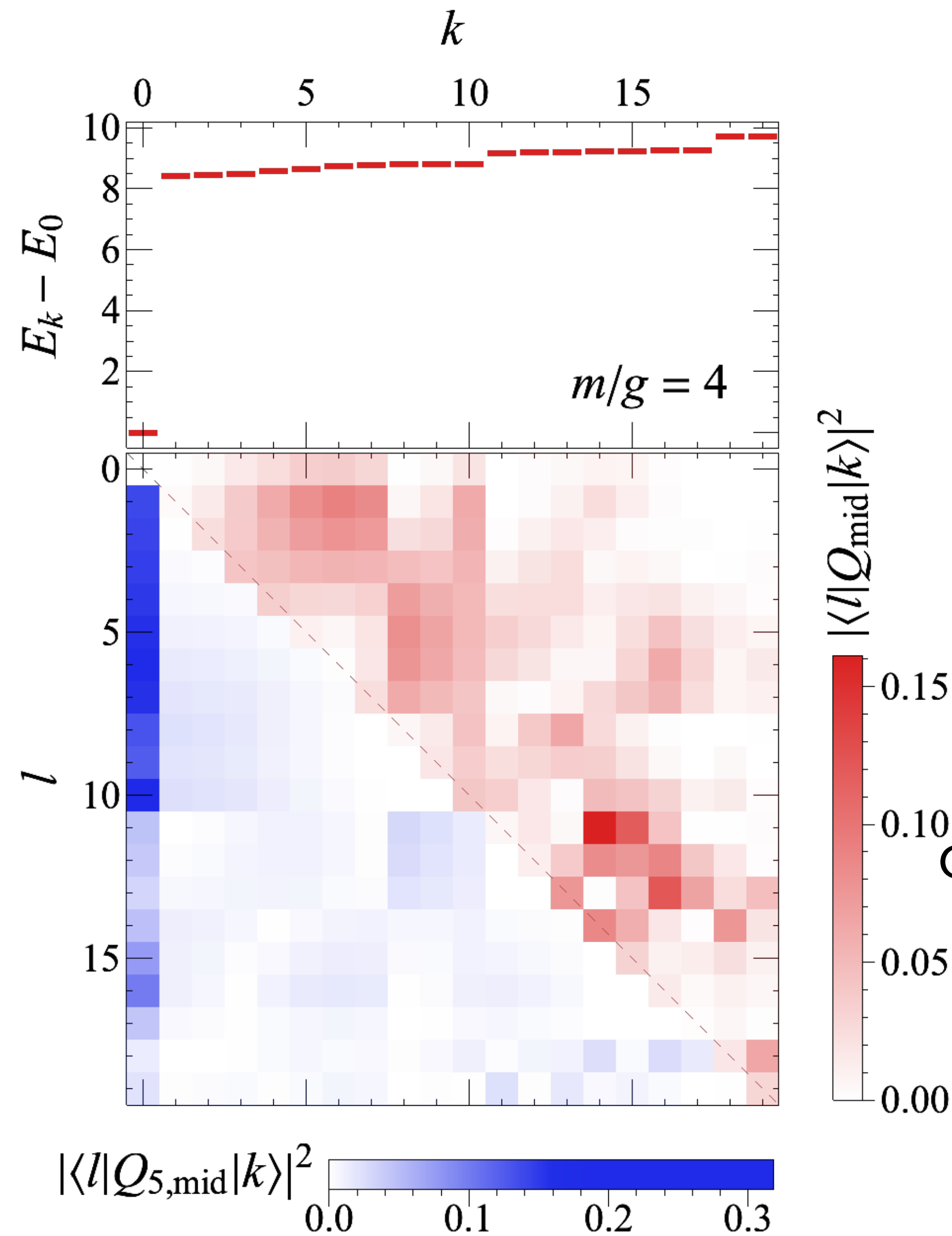




**short** period oscillation  
 in both **vector** and **axial** charge







$$H|k\rangle = E_k|k\rangle$$

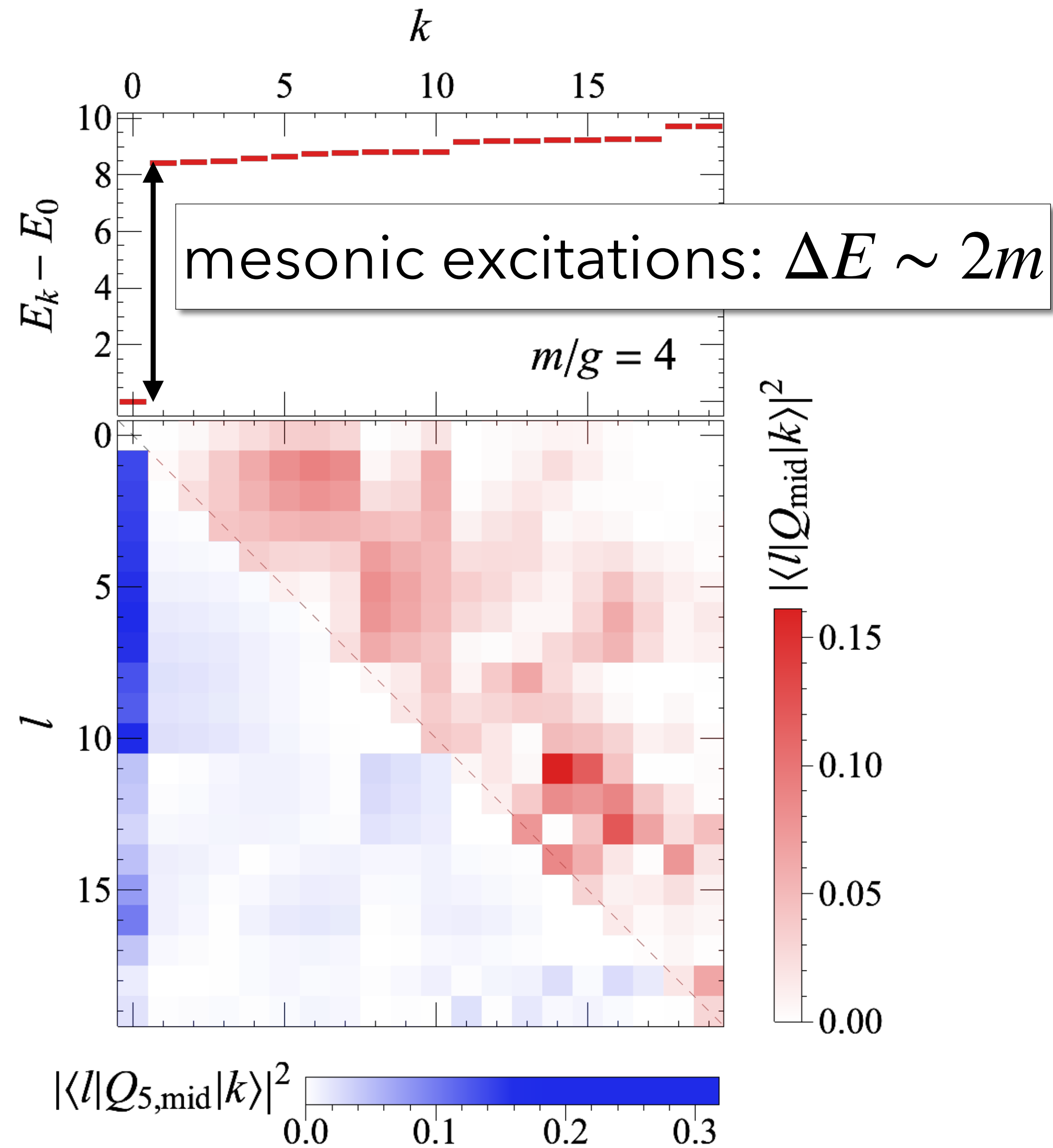
$$|\Psi(t=0)\rangle = \sum_k c_k |k\rangle$$

$$O(t) \equiv \langle \Psi(t) | O | \Psi(t) \rangle = \sum_{k,l} c_k c_l^* e^{i(E_l - E_k)t} \langle l | O | k \rangle$$

oscillation *frequency* ← energy difference

oscillation *strength* ← matrix element

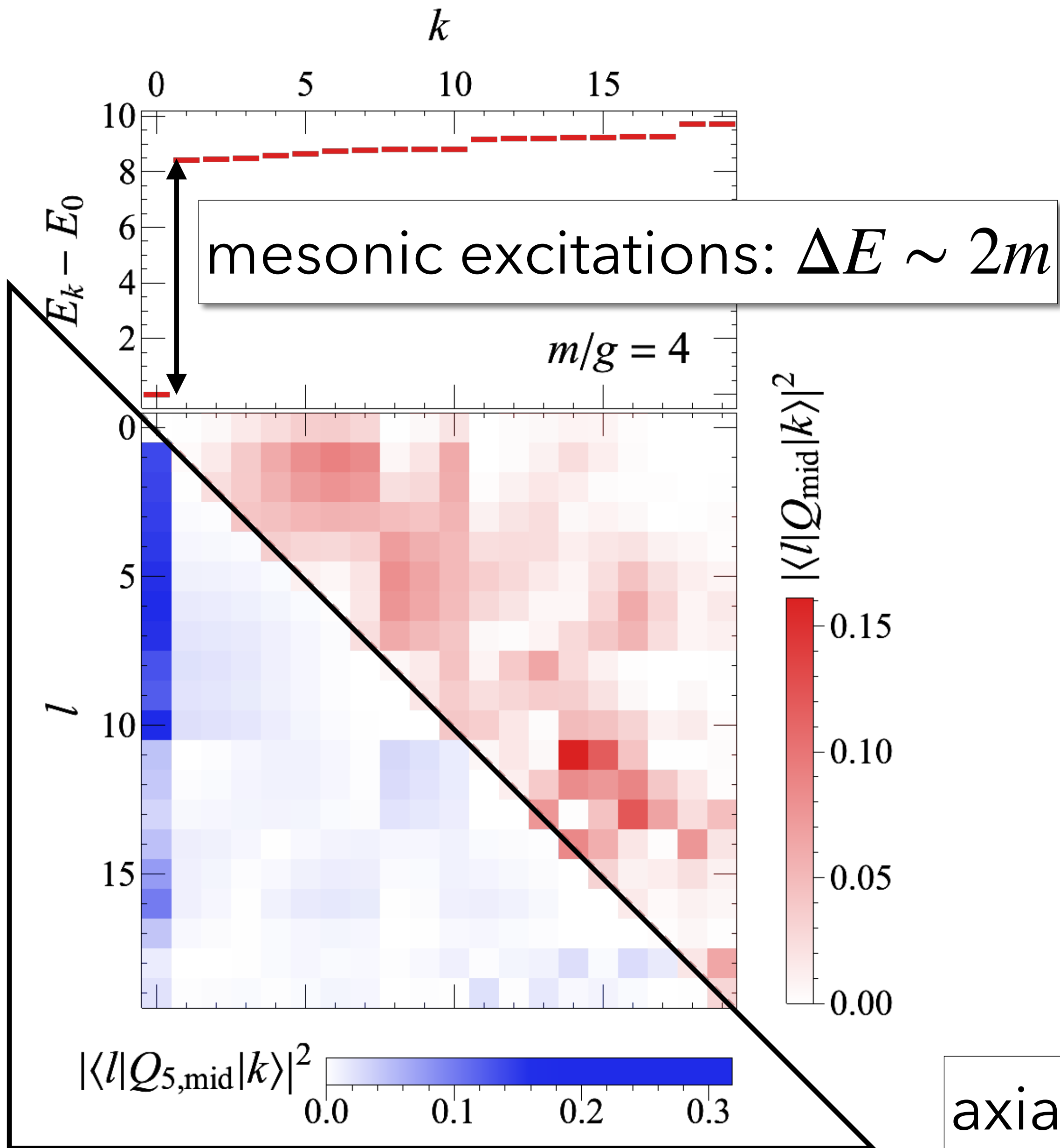




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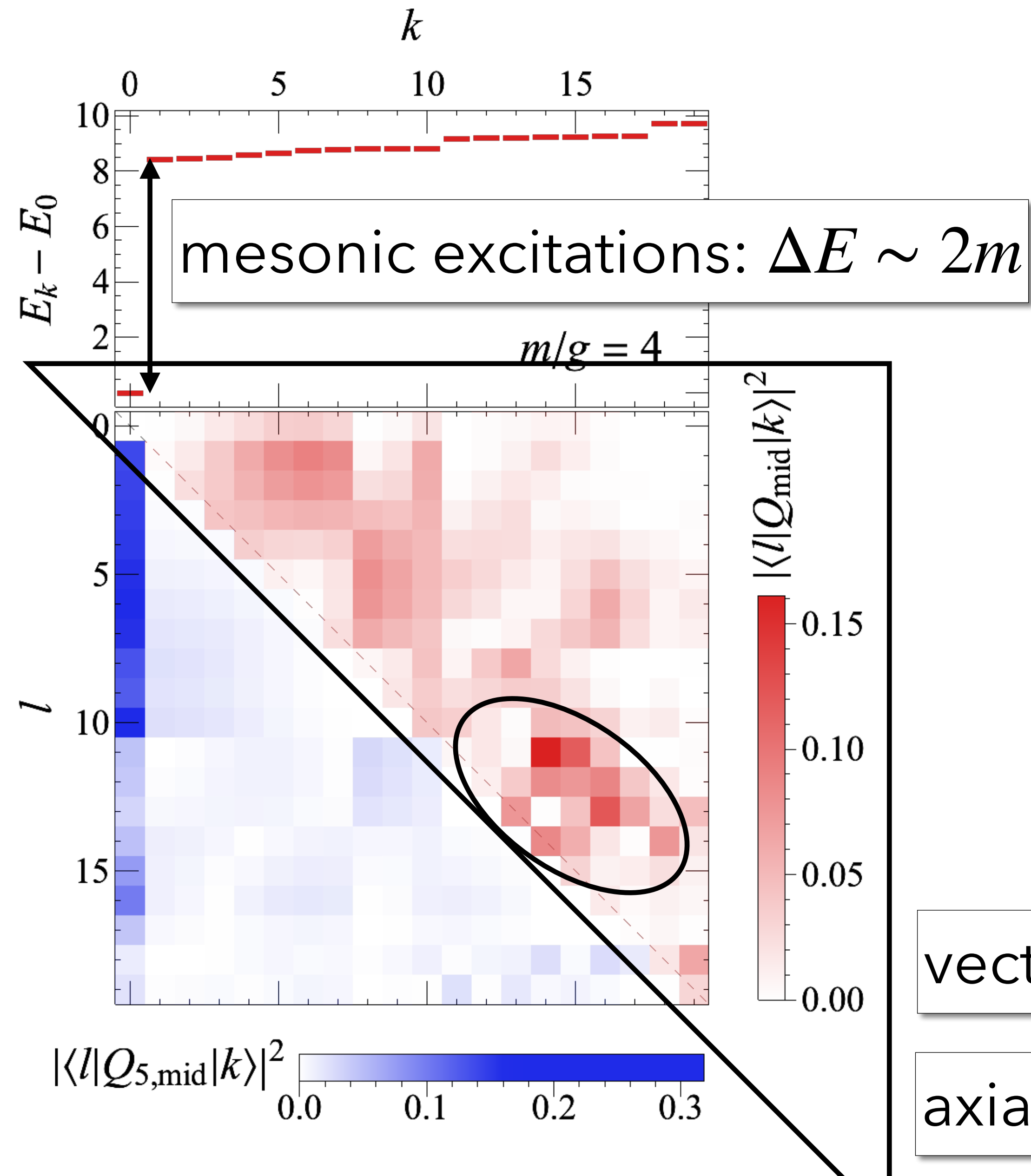


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axial charge: ground state  $\leftrightarrow$  excitation



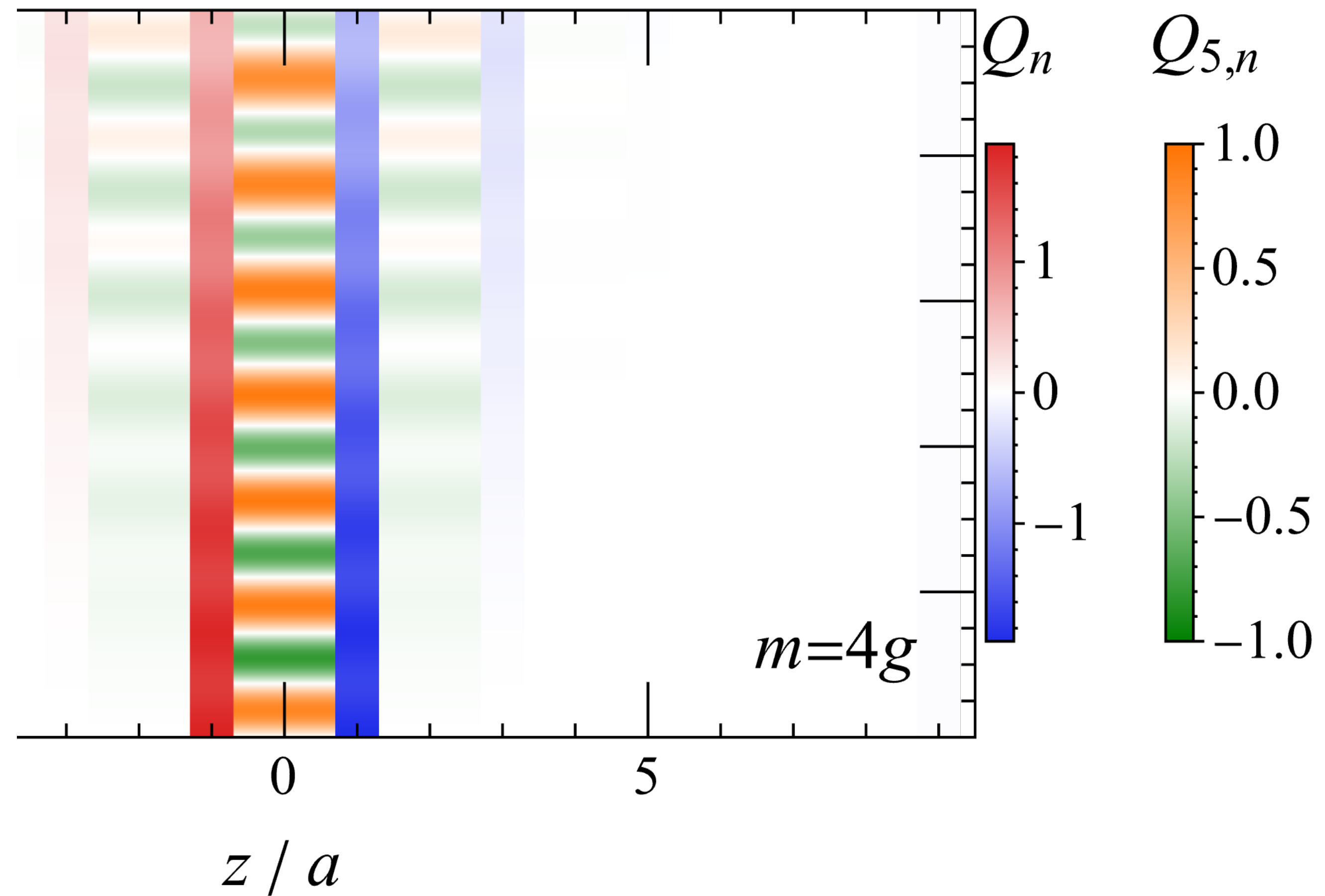
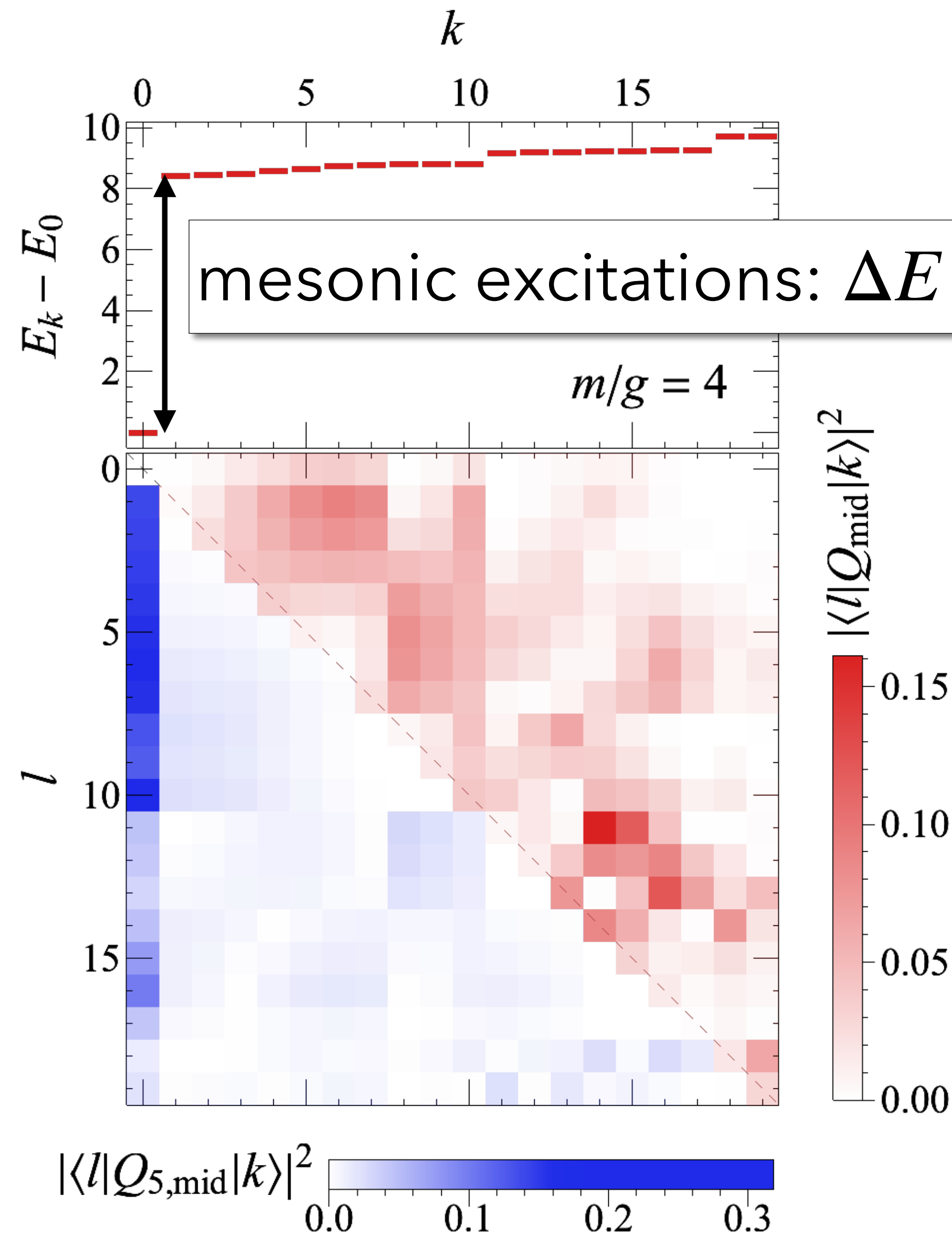
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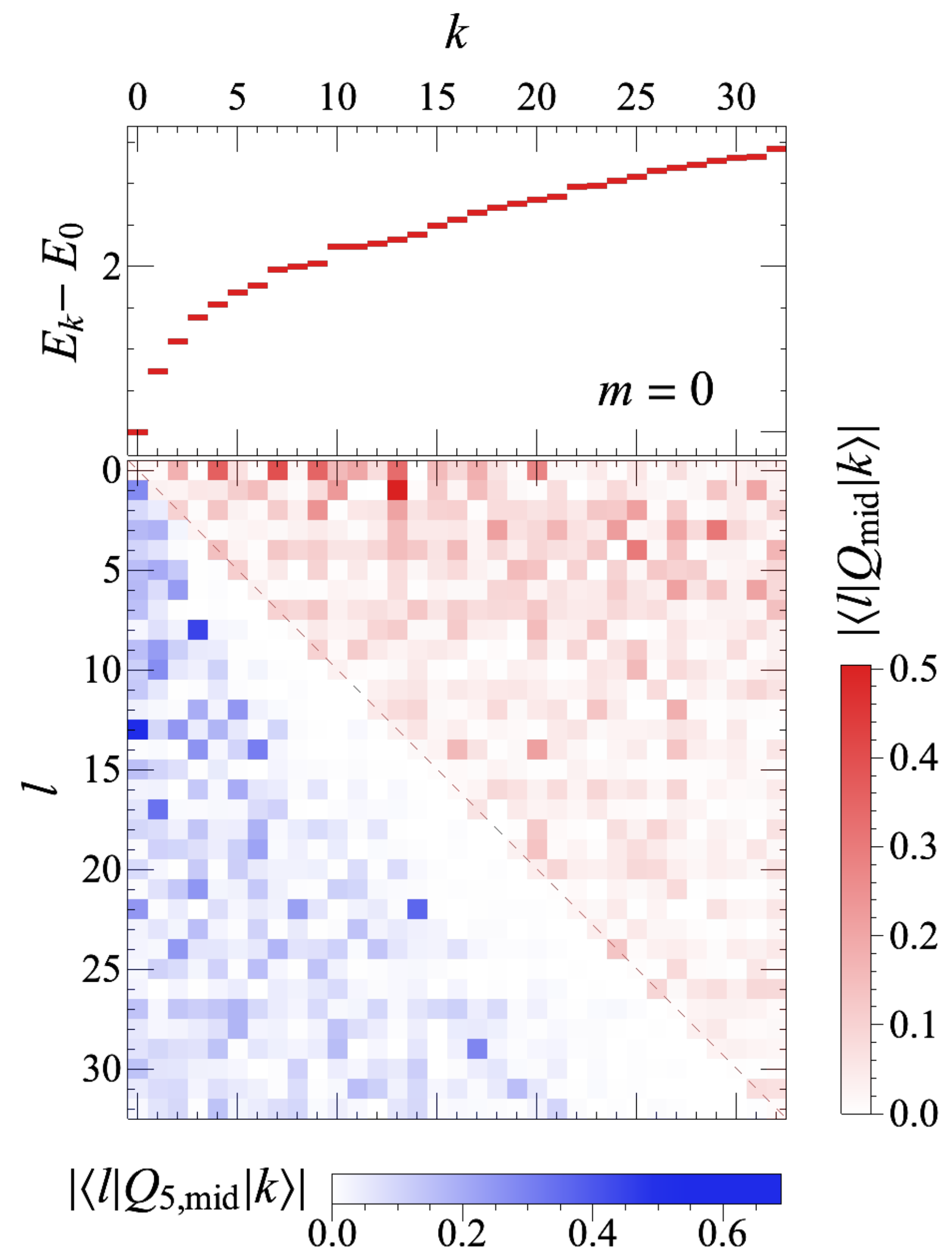
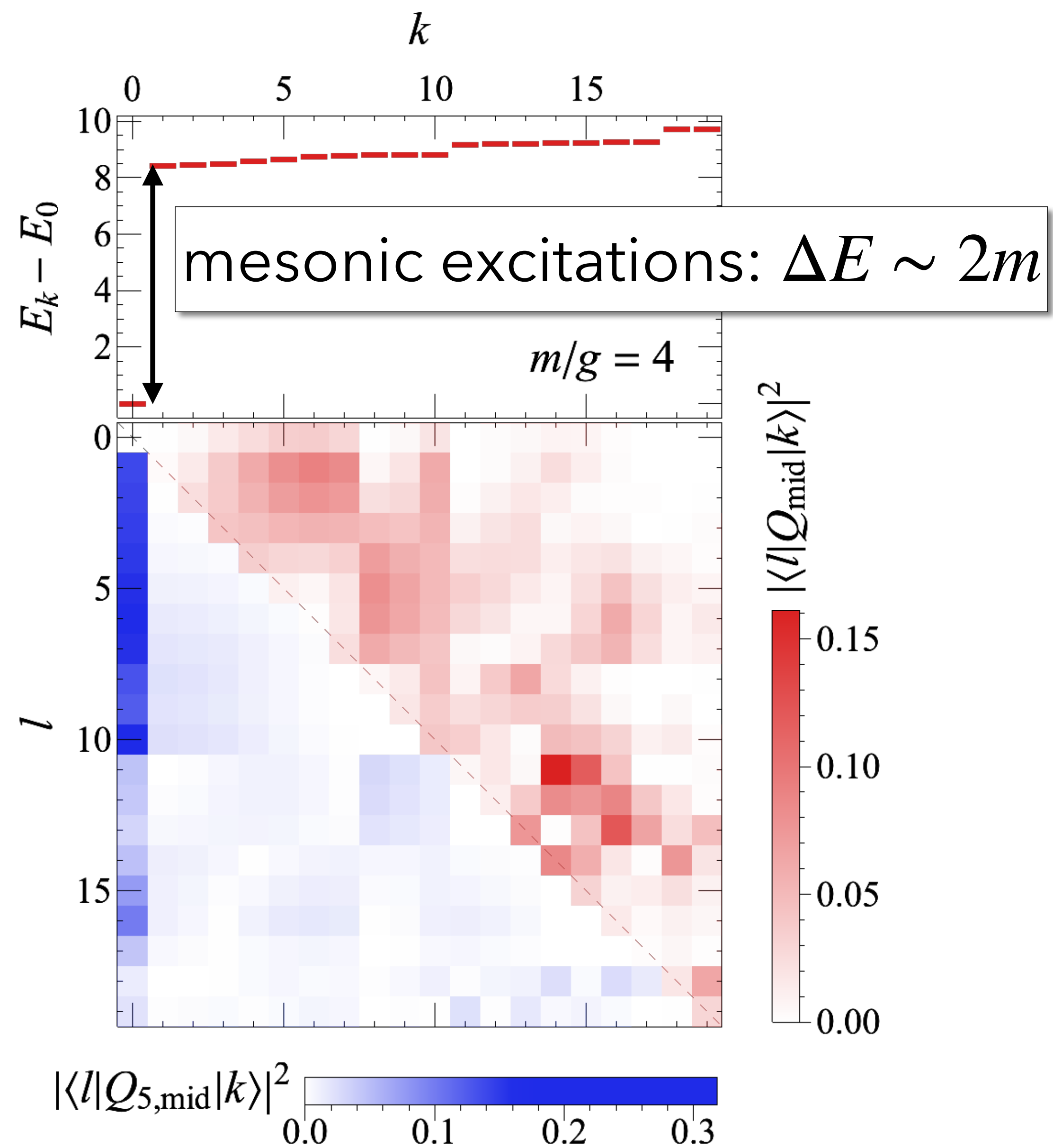
vector charge: excitation  $\leftrightarrow$  excitation

axial charge: ground state  $\leftrightarrow$  excitation



vector charge: excitation  $\leftrightarrow$  excitation

axial charge: ground state  $\leftrightarrow$  excitation



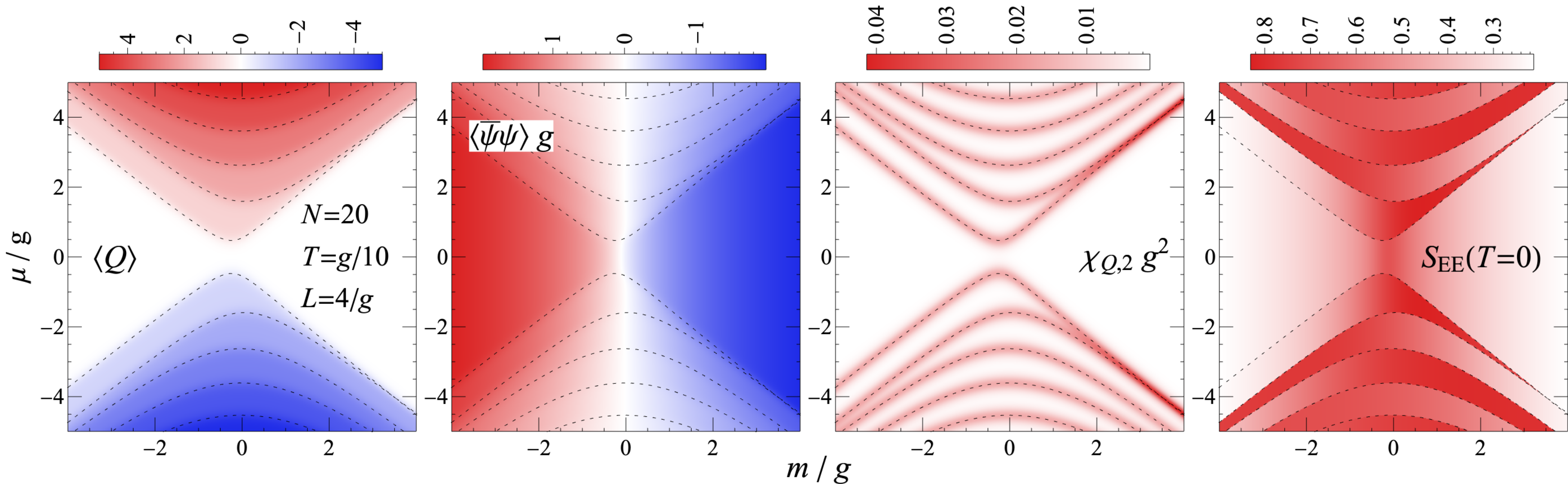
# *III. Phase Structure*

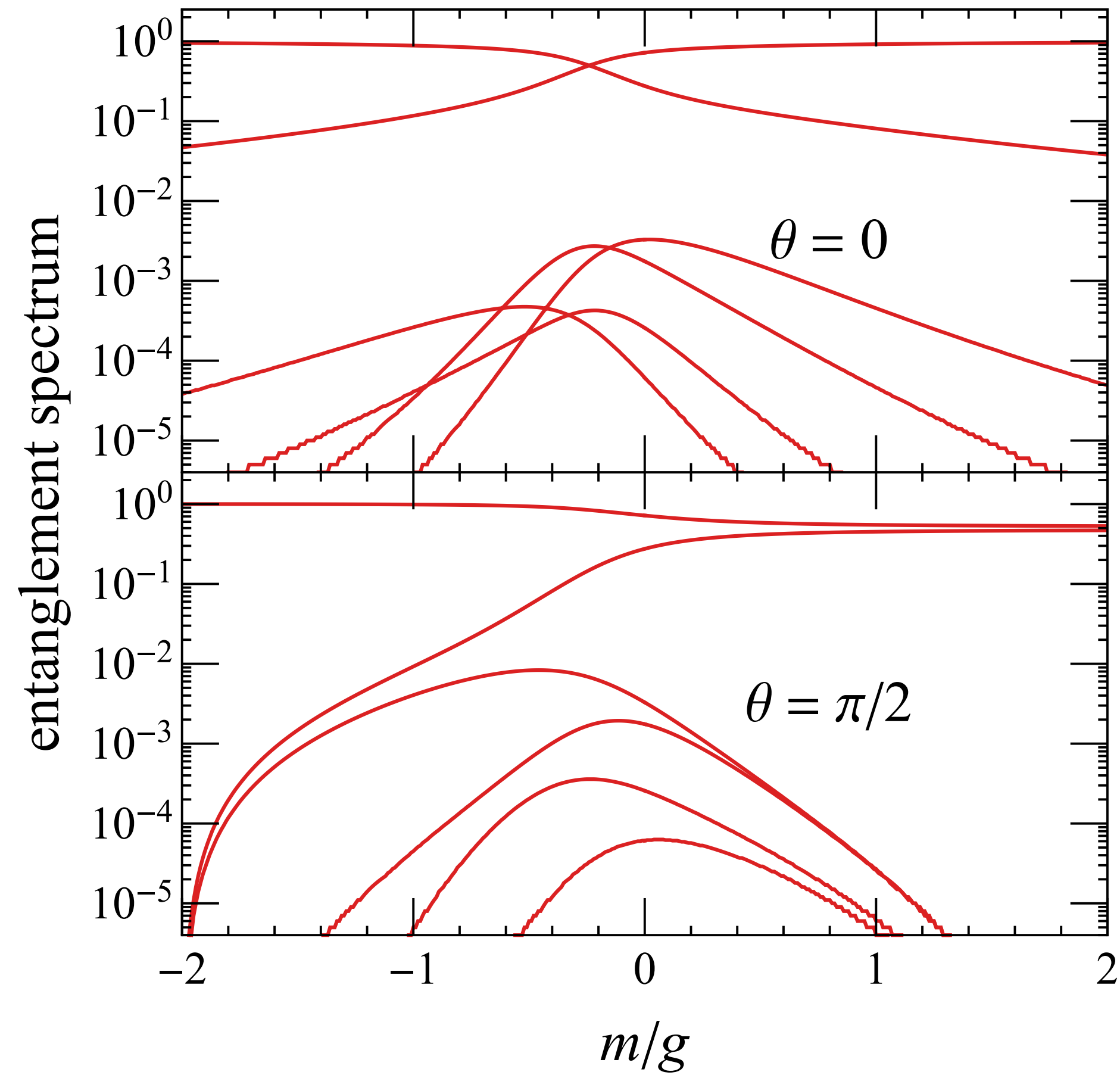
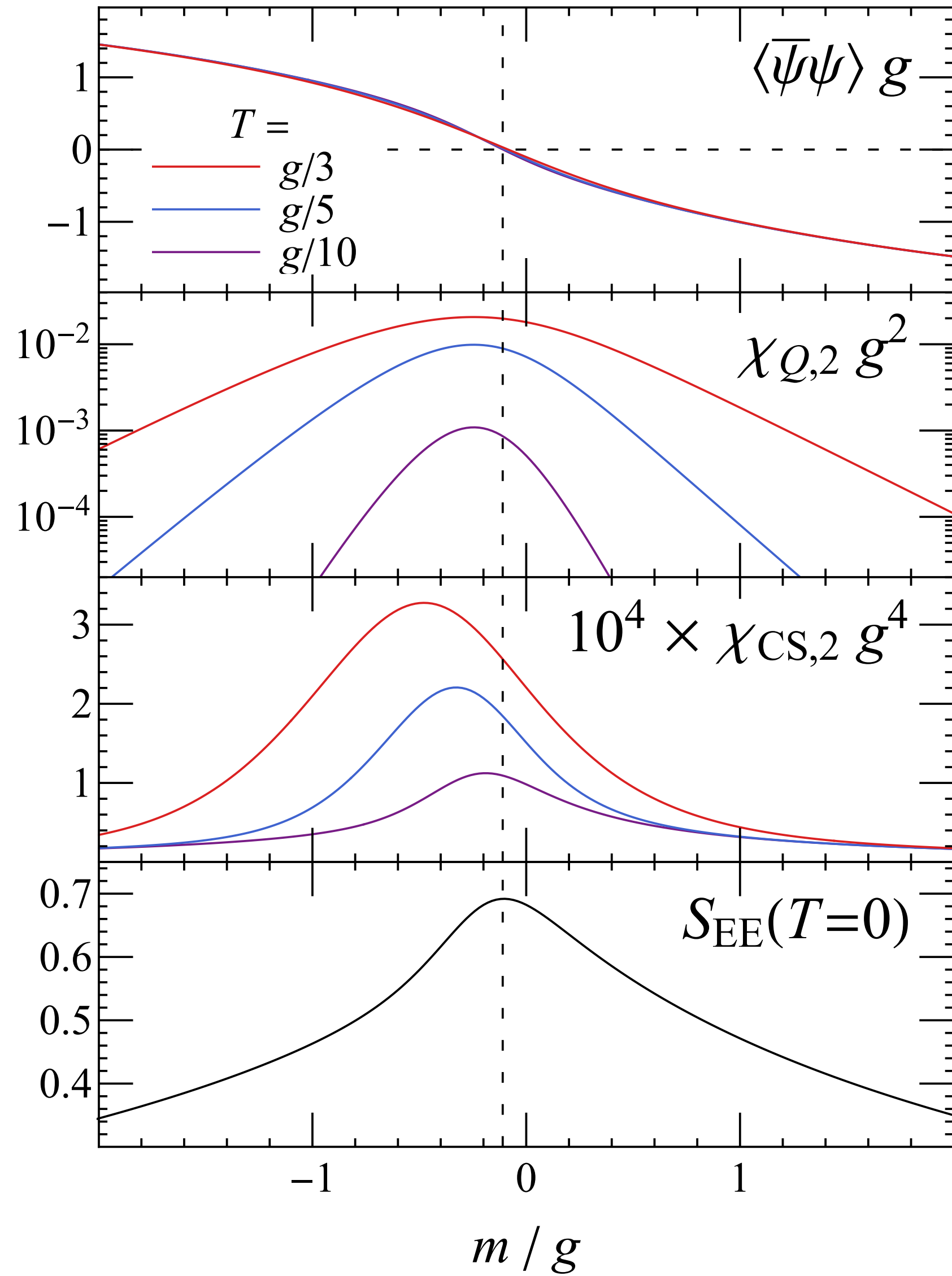
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- finite temperature, finite chemical potential:

$$\langle O \rangle_{\text{th}} \equiv \text{Tr}(\rho_{\text{th}} O) \quad \rho_{\text{th}} \equiv \frac{e^{-(H-\mu Q)/T}}{\text{Tr}(e^{-(H-\mu Q)/T})}$$







## summary

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- real-time dynamics in Schwinger model
  - jet production: spread out of light cone, creation of fermion-antifermion pairs
  - charge transport: thumper and breather modes
  - critical point detected by different signals
- Need **quantum computers** to approach the continuum limit.

Students and Postdocs are welcome!

(so are questions/comments)