

Nonlinear Chiral Kinetic Theory

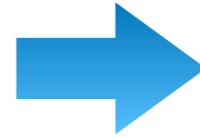
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Transport Phenomena

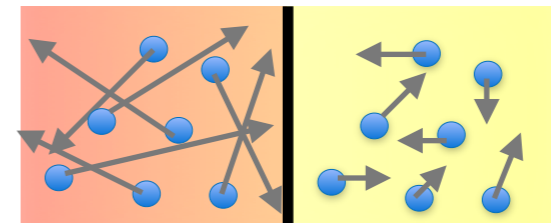
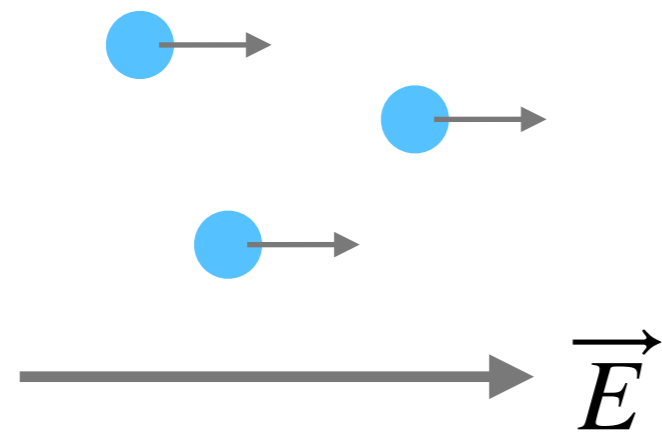
particle



current

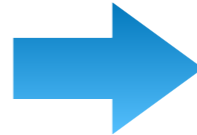


external force



Transport Phenomena

particle

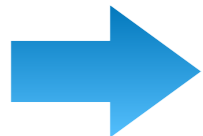


current



external force

Ohm's law

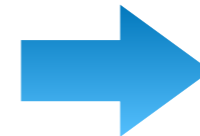


electric current



\vec{E}

Fourier's law



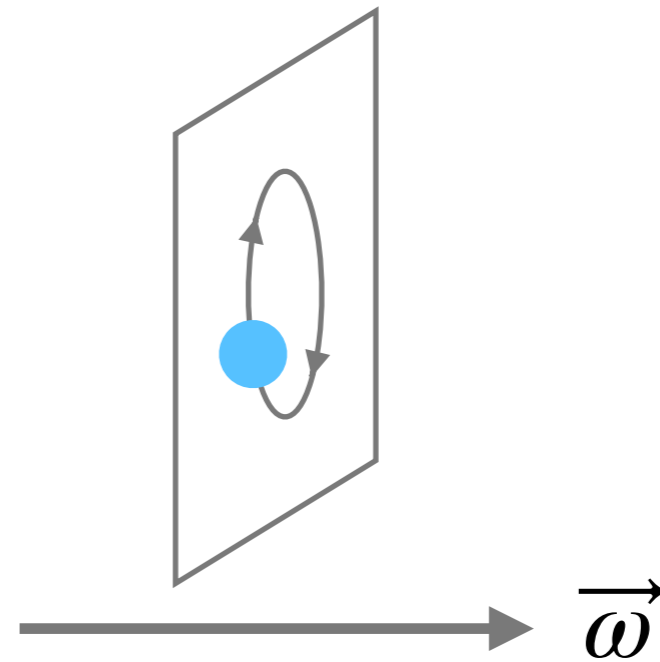
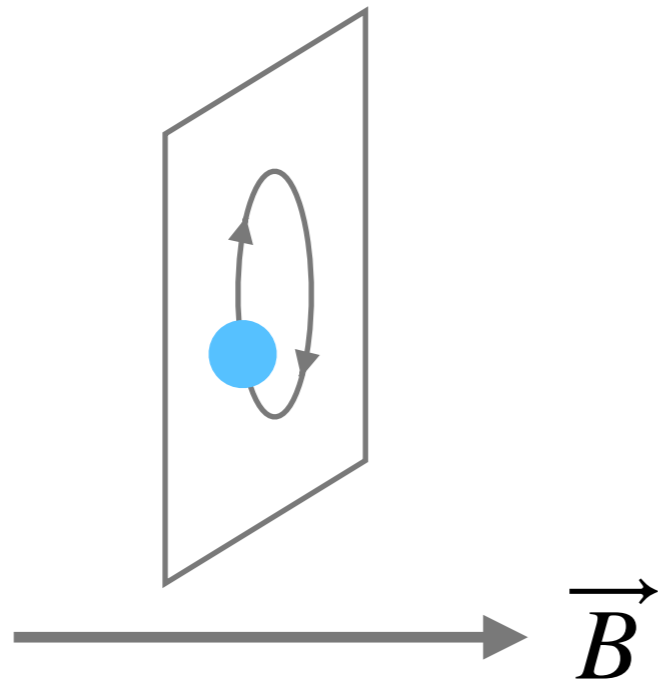
energy current



$-\vec{\nabla} T$

Transport Phenomena

- classical



Transport Phenomena

- quantum and massless

magnetization coupling

momentum



spin

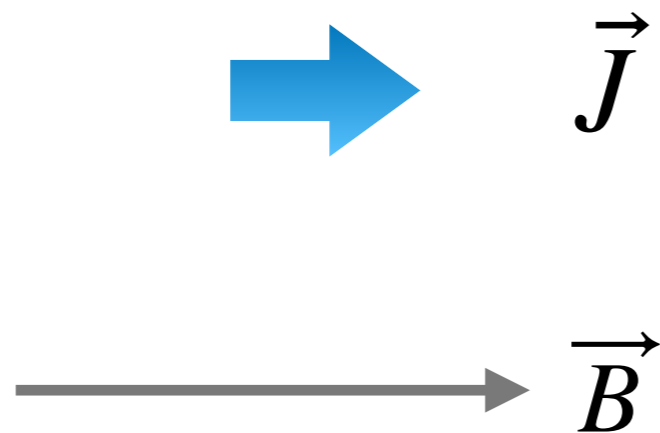


spin-vorticity coupling



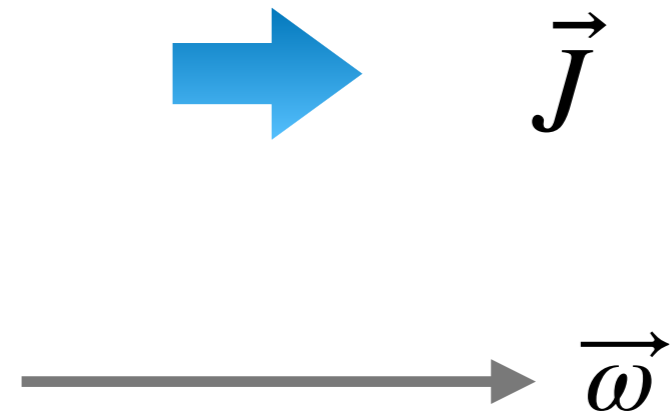
Transport Phenomena

- quantum and massless



chiral magnetic effect

Vilenkin (1980)
Fukushima, Kharzeev, Warringa (2008)



chiral vortical effect

Vilenkin (1979), Son, Surowka (2009)
Landsteiner, Megias, Pena-Benitez (2011)

Chiral Transport Phenomena

chiral magnetic effect

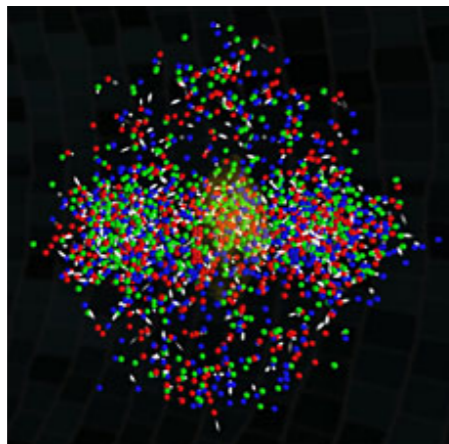
$$\vec{J} = \sigma_{\text{CME}} \vec{B}$$

chiral vortical effect

$$\vec{J} = \sigma_{\text{CVE}} \vec{\omega}$$

from **quantum anomaly** in Dirac theory

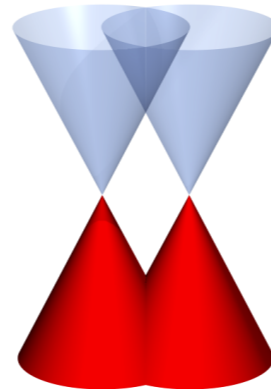
hep



quark-gluon plasma

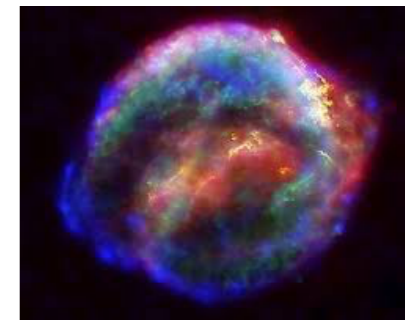
cond-mat

Vazifeh, Franz (2013)



Weyl semimetal

astro



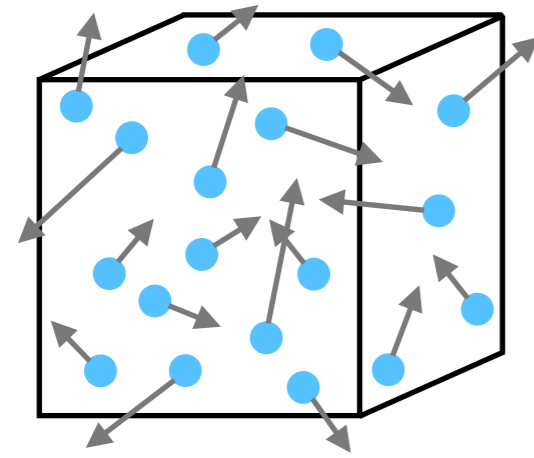
neutrino

for matter in early universe
(2010)



for any chiral matter
(present)

Chiral Kinetic Theory



Boltzmann eq.

+



chirality

=

Chiral Kinetic Theory(CKT)

Stephanov, Yin (2012)
Son, Yamamoto (2012)
Chen et al. (2013)

- Lorentz covariance
- collision
- mass correction
- spin polarization

- strong magnetic field
- gravitational field
- other derivations

Hidaka, Pu, Wang, Yang (2022)

Limitation 1 : magneto-vortical transport

$$\mathcal{L} = \bar{\psi} (i\hbar \partial_{\mu} - eA_{\mu}) \psi$$

$$\longrightarrow O(\partial) \sim O(\hbar)$$

$$O(1)$$

$$O(\hbar)$$

$$O(\hbar^2)$$

Hattori, Yin (2016)

$$J^0 \sim \mu^3$$

Fermi sphere

Boltzmann

$$\vec{J} \sim \mu \vec{B}$$

CME

$$\vec{J} \sim \mu^2 \vec{\omega}$$

CVE

linear CKT

$$J^0 \sim \vec{B} \cdot \vec{\omega}$$

nonlinear CKT

$$B \sim \partial A$$

$$\omega \sim \partial u$$

Limitation 1 : magneto-vortical transport

$$O(\hbar)$$

$$\vec{J} \sim \mu \vec{B} \quad \vec{J} \sim \mu^2 \vec{\omega}$$

anomaly, nondissipative

$$O(\hbar^2)$$

$$J^0 \sim \vec{B} \cdot \vec{\omega}$$

anomaly? nondissipative?

- linear response [Hattori, Yin \(2016\)](#)
- rotating fermions [Ebihara, Fukushima, KM \(2017\)](#)
- near-equilibrium Wigner function [Yang et al. \(2020\), Lin, Yang \(2021\)](#)

Limitation 1 : magneto-vortical transport

$$O(\hbar)$$

$$\vec{J} \sim \mu \vec{B} \quad \vec{J} \sim \mu^2 \vec{\omega}$$

anomaly, nondissipative

$$O(\hbar^2)$$

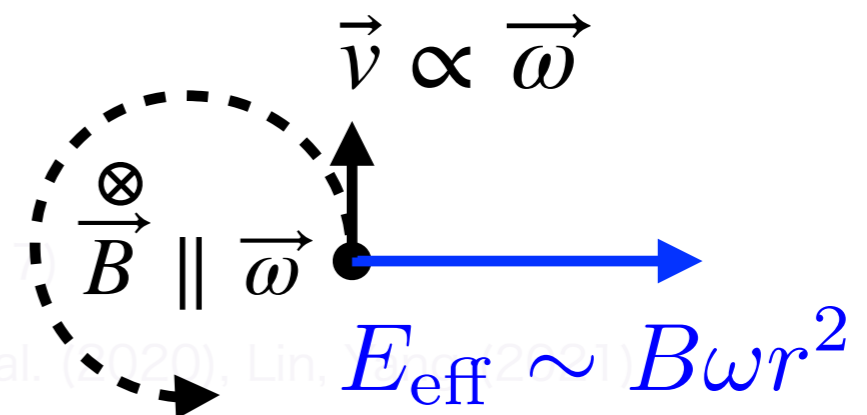
$$J^0 \sim \vec{B} \cdot \vec{\omega}$$

anomaly? nondissipative?

• linear response Hattori, Yin (2016)

• Thermal Field Theory, but...

• near-equilibrium Wigner function Yang et al. (2020), Lin, Yang (2021)



Equilibrated?

Check it from **nonlinear CKT**

Limitation 2 : trace anomaly

$$\mathcal{L} = \bar{\psi}(i\hbar\partial_{\mu} - eA_{\mu})\psi$$

→ $O(\partial) \sim O(\hbar)$

$O(1)$

$\partial_{\mu}J^{\mu} = 0$

$T^{\mu}_{\mu} = 0$

Boltzmann

$O(\hbar)$

$\partial_{\mu}J^{\mu} \sim \vec{E} \cdot \vec{B}$

chiral anomaly

linear CKT

$E, B \sim \partial A$

$O(\hbar^2)$

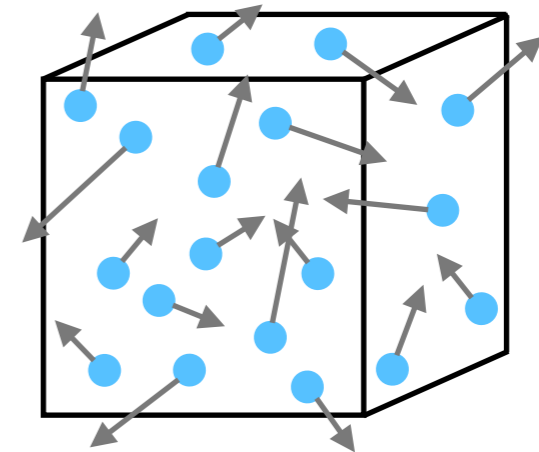
$T^{\mu}_{\mu} \sim \beta(e)F_{\mu\nu}^2$

trace anomaly

nonlinear CKT

Limitation 2 : trace anomaly

$$\mathcal{L}_{\text{QED}} = -\frac{1}{4}F_{\mu\nu}^2 + \bar{\psi} i\not{D}\psi$$



nonlinear CKT

$$O(\hbar)$$

$$\partial_\mu J^\mu \sim \vec{E} \cdot \vec{B}$$



Berry monopole

$$O(\hbar^2)$$

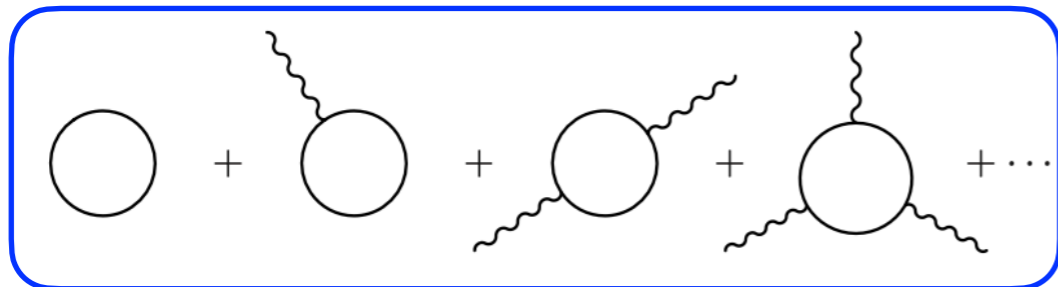
$$T^\mu{}_\mu \sim \beta(e)F_{\mu\nu}^2$$



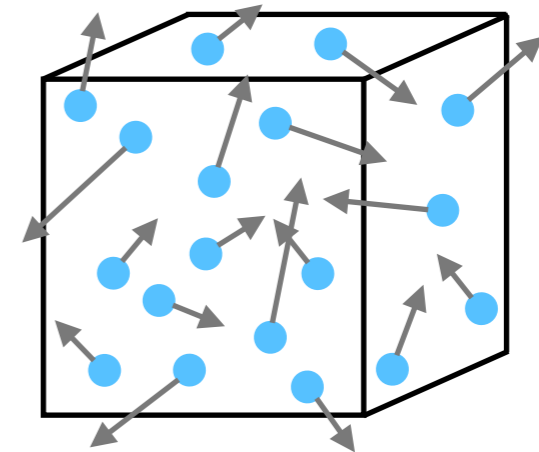
???

Limitation 2 : trace anomaly

$$\mathcal{L}_{\text{EH}} = -\frac{1}{4}F_{\mu\nu}^2 + \mathcal{L}_{\text{fer}} \parallel$$



Euler-Heisenberg theory



nonlinear CKT

$$O(\hbar)$$

$$\partial_\mu J^\mu \sim \vec{E} \cdot \vec{B}$$



Berry monopole

$$O(\hbar^2)$$

$$T^\mu{}_\mu \sim \beta(e)F_{\mu\nu}^2$$



???

Quantum kinetic theory

- classical

$$n_{\mathbf{p}}(t, \mathbf{x}, \mathbf{p}) \quad \text{no uncertainty}$$

- quantum field theory (relativistic)

$$W(x, p) = \int_y \underline{e^{-ip \cdot y / \hbar}} \langle \bar{\psi}(x + \frac{y}{2}) \psi(x - \frac{y}{2}) \rangle$$

uncertainty Green's function

- right-handed current density

$$\mathcal{R}^\mu(x, p) = \text{tr} \left[\gamma^\mu \frac{1 + \gamma^5}{2} W(x, p) \right]$$

Quantum kinetic theory

- classical

$$n_{\mathbf{p}}(t, \mathbf{x}, \mathbf{p}) \quad \text{no uncertainty}$$

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$$W(x, p) = \int_y \underline{e^{-ip \cdot y / \hbar}} \langle \bar{\psi}(x + \frac{y}{2}) \psi(x - \frac{y}{2}) \rangle$$

uncertainty Green's function

- right-handed current density

$$\mathcal{R}^\mu(x, p) \sim (n_{\mathbf{p}}, \mathbf{v}n_{\mathbf{p}}) + O(\hbar)$$

- chiral kinetic theory

$$\text{EoMs for } \mathcal{R}^\mu = \mathcal{R}_{(0)}^\mu + \hbar \mathcal{R}_{(1)}^\mu + \hbar^2 \mathcal{R}_{(2)}^\mu$$

EoMs (zeroth order)

- number conservation (kinetic equation) $(\partial_\mu - F_{\mu\nu}\partial_p^\nu)\mathcal{R}^\mu = 0$
- scale invariance $p_\mu\mathcal{R}^\mu = 0$
- AM conservation $p^\mu\mathcal{R}^\nu - p^\nu\mathcal{R}^\mu = 0$

current $\mathcal{R}^\mu(x, p) \sim (n_{\mathbf{p}}, \mathbf{v}n_{\mathbf{p}}) + O(\hbar)$

EM tensor $\mathcal{R}^\mu p^\nu$

Zeroth-order Wigner function

- general solution

$$\mathcal{R}_{(0)}^\mu = 2\pi \delta(p^2) p^\mu f_{(0)}(x, p)$$

on-shell condition distribution function

- current

$$J_{(0)}^0 = \int_{\mathbf{p}} n_{\mathbf{p}} \quad J_{(0)}^i = \int_{\mathbf{p}} \frac{p^i}{|\mathbf{p}|} n_{\mathbf{p}}$$

$$n_{\mathbf{p}} = f_{(0)}(p_0 = |\mathbf{p}|) - f_{(0)}(p_0 = -|\mathbf{p}|)$$

Linear-order Wigner function

- general solution Hidaka, Pu, Yang (2016)

$$\mathcal{R}_{(1)}^\mu = 2\pi\delta(p^2) \left[p^\mu f_{(1)} + \left(\Sigma_n^{\mu\nu} \Delta_\nu - \frac{1}{p^2} \tilde{F}^{\mu\nu} p_\nu \right) f_{(0)} \right]$$

- spin tensor

$$\Sigma_n^{\mu\nu} = \frac{\varepsilon^{\mu\nu\rho\sigma} p_\rho n_\sigma}{2p \cdot n}$$

e.g. $n^\mu = (1, \mathbf{0})$

$$(\Sigma^{23}, \Sigma^{31}, \Sigma^{12}) = \frac{\mathbf{p}}{2p_0} \sim \frac{\hat{\mathbf{p}}}{2}$$

Chen, Son, Stephanov (2015)

- $\mathcal{R}_{(1)}^\mu$ should be frame-independent

Nonlinear-order Wigner function

- general solution KM (2023)

$$\begin{aligned}
 \mathcal{R}_\mu^{(2)} = & 2\pi\delta(p^2) \left[p_\mu f_{(2)} + \left(\Sigma_{\mu\nu}^u \Delta^\nu - \frac{1}{p^2} \tilde{F}_{\mu\nu} p^\nu \right) f_{(1)} - \Sigma_{\mu\nu}^u \varepsilon^{\nu\rho\sigma\lambda} \Delta_\rho \frac{n_\sigma}{2p \cdot n} \Delta_\lambda f_{(0)} \right] \\
 & + \frac{2\pi}{p^2} \left[-p_\mu Q \cdot p + 2p^\nu Q_{[\mu} p_{\nu]} \right] f_{(0)} \delta(p^2) \\
 & + 2\pi \frac{\delta(p^2)}{p^2} \left(\frac{1}{2} \varepsilon_{\mu\nu\rho\sigma} p^\nu \Delta^\rho + \frac{p_\mu p^\nu}{p^2} \tilde{F}_{\nu\sigma} - \tilde{F}_{\mu\sigma} \right) \left(\Sigma_n^{\sigma\lambda} \Delta_\lambda - \frac{1}{p^2} \tilde{F}^{\sigma\lambda} p_\lambda \right) f_{(0)} \\
 & + 2\pi \frac{\delta(p^2)}{p^2} \Sigma_{\mu\nu}^u \left[\Delta_\alpha \Sigma_n^{\alpha\nu} + \frac{n_\alpha}{p \cdot n} \tilde{F}^{\alpha\nu} + \frac{1}{p^2} \tilde{F}^{\nu\lambda} p_\lambda \right] p \cdot \Delta f_{(0)}
 \end{aligned}$$

- another frame vector? u^μ Hayata, Hidaka, KM (2021)

- $\mathcal{R}_{(2)}^\mu$ should be frame-independent

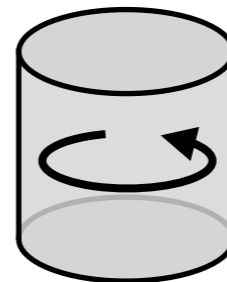
Equilibrium from Boltzmann Theory

$$f^{\text{eq}} = f^{\text{eq}}(g)$$

$$g = \underbrace{-\beta\mu}_{\text{charge}} + \underbrace{\beta \cdot p}_{\text{momentum}} + \underbrace{\frac{\hbar}{2} \Sigma_n^{\mu\nu} \partial_\mu \beta_\nu}_{\text{total angular momentum}}$$

$$\beta^\mu = \beta(1, \mathbf{x} \times \boldsymbol{\omega})$$

e.g. $n^\mu = (1, \mathbf{0})$



$$g = \beta(p_0 - \mu - \boldsymbol{\omega} \cdot (\mathbf{x} \times \mathbf{p}) - \hbar \boldsymbol{\omega} \cdot \hat{\mathbf{p}}/2)$$

Equilibrium from Boltzmann Theory

$$f^{\text{eq}} = f^{\text{eq}}(g)$$

$$g = \underbrace{-\beta\mu}_{\text{charge}} + \underbrace{\beta \cdot p}_{\text{momentum}} + \underbrace{\frac{\hbar}{2} \Sigma_n^{\mu\nu} \partial_\mu \beta_\nu}_{\text{total angular momentum}}$$



$\mathcal{R}_{\text{eq}}^\mu$ frame independent



kinetic eq. satisfied

$$(\partial_\mu - F_{\mu\nu} \partial_p^\nu) \mathcal{R}_{\text{eq}}^\mu = 0$$

Linear-order equilibrium Wigner function

$$O(\partial) \sim O(\hbar)$$

- one derivative Hidaka, Pu, Yang (2016)

$$\mathcal{R}_\mu^{(1)} = \mathcal{R}_\mu^{(F)} + \mathcal{R}_\mu^{(\omega)}$$

- momentum integral at equilibrium

$$\text{e.g. } J_{(1)}^\mu = \underbrace{\frac{\mu}{4\pi^2} B^\mu}_{\text{CME}} + \underbrace{\left(\frac{\mu^2}{4\pi^2} + \frac{T^2}{12} \right)}_{\text{CVE}} \omega^\mu$$

- Nondissipative due to quantum anomaly

Equilibrium at **nonlinear** order

linear

$$f^{\text{eq}} = f^{\text{eq}}(g)$$

$$g = -\beta\mu + \beta \cdot p + \frac{\hbar}{2} \Sigma_n^{\mu\nu} \partial_\mu \beta_\nu$$



nonlinear

$$f^{\text{eq}} = f^{\text{eq}}(g)$$



minimal conditions



$\mathcal{R}_{\text{eq}}^\mu$ frame independent



kinetic eq. satisfied

Equilibrium at **nonlinear** order

nonlinear-order f^{eq} exists only when

$$\partial_{\mu}(\beta_{\mu}) + F_{\mu\nu}\beta^{\nu} = 0, \quad \partial_{\mu}\beta_{\nu} + \partial_{\nu}\beta_{\mu} = 0$$

with

$$1) \quad \beta \cdot \partial F_{\mu\nu} = 0, \quad \beta \cdot \partial \beta_{\mu} = 0$$

or

$$2) \quad \partial_{\lambda} F_{\mu\nu} = 0 \quad \text{Yang et al. (2020)}$$

Nonlinear-order equilibrium Wigner function

$$O(\partial) \sim O(\hbar)$$

- two derivatives Yang et al. (2020) KM (2023)

$$\mathcal{R}_\mu^{(2)} = \mathcal{R}_\mu^{(\partial F)} + \mathcal{R}_\mu^{(FF)} + \mathcal{R}_\mu^{(F\omega)} + \mathcal{R}_\mu^{(\omega\omega)}$$

- momentum integral at equilibrium

e.g.
$$J_{(F\omega)}^0 = -\frac{1}{8\pi^2} B \cdot \omega$$

- Nondissipative due to quantum anomaly

conformed from quantum field theory

UV divergence

- $f_{(0)}(x, p)$ contains **vacuum** contribution

$$\mathcal{R}_{(0)}^\mu = 2\pi\delta(p^2)p^\mu f_{(0)}(x, p)$$

- **vacuum** contribution to current

$$J_{(\partial F)} \underset{+3}{\sim} \underset{+3}{\partial F} \times \int_{\mathbf{p}} \underset{+0}{|\mathbf{p}|^{-3}} \left[n_F(|\mathbf{p}| - \mu) + n_F(|\mathbf{p}| + \mu) \underset{-1}{-1} \right]$$

- IR safe

$$n_F(|\mathbf{p}| - \mu) + n_F(|\mathbf{p}| + \mu) \longrightarrow 1$$

- UV **logarithmic** divergence

$$n_F(|\mathbf{p}| - \mu) + n_F(|\mathbf{p}| + \mu) \longrightarrow 0$$

Regularization problem

- Pauli-Villars $\frac{1}{k^2 - m^2} \rightarrow \frac{1}{k^2 - m^2} - \frac{1}{k^2 - \Lambda^2}$

inapplicable: no mass parameter in CKT

Yang et al. (2020)

- 't Hooft-Veltman (dimensional regularization) $\int d^4k \rightarrow \int d^{4-\epsilon}k$

inapplicable: $\mathcal{R}^\mu = \text{tr} \left[\gamma^\mu \frac{1 + \gamma^5}{2} W \right]$ incorrect in $d \neq 4$

$\gamma^\mu = (\gamma^0, \dots, \gamma^3, \gamma^4, \dots, \gamma^{d-1})$ commute with γ^5

Point-Splitting regularization in CKT

- charge current

$$J^\mu(x, \mathbf{y}) = \text{tr} \left\langle \bar{\psi}(x + \frac{\mathbf{y}}{2}) \gamma^\mu \frac{1 + \gamma^5}{2} \psi(x - \frac{\mathbf{y}}{2}) \right\rangle = \int_p e^{ip \cdot \mathbf{y} / \hbar} \mathcal{R}^\mu(x, p)$$

- UV divergence $\mathcal{J} = \int_0^{\mathbf{y}^{-1}} \frac{dp}{p}$

$$p \sim \infty \longleftrightarrow \mathbf{y} \sim 0$$

- regularized vacuum part

$$\begin{aligned} J_{(\partial F) \text{ vac}}^\mu &= \frac{\mathcal{J}}{12\pi^2} \partial_\lambda F^{\mu\lambda} \\ +3 &= +0 +3 \end{aligned}$$

$$\begin{aligned} T_{(FF) \text{ vac}}^{\mu\nu} &= -\frac{\mathcal{J}}{12\pi^2} \left[F^\mu{}_\sigma F^{\nu\sigma} - \frac{1}{4} g^{\mu\nu} F_{\alpha\beta}^2 \right] \\ +4 &= +0 +4 \end{aligned}$$

Compensation of PS regularization

- gauge invariance unbroken

$$\partial_{\mu} J_{(2)}^{\mu} = 0$$



- translational invariance **broken**

$$\partial_{\mu} T_{(2)}^{\mu\nu} + F^{\mu\nu} J_{\mu}^{(2)} \neq 0$$



- conformal invariance **unbroken**

$$T^{\mu}_{\mu(2)} = 0$$



because of ~~classical EM fields?~~ **regularization**

Yang et al. (2020)

Euler-Heisenberg theory

- Effective Lagrangian

$$\mathcal{L}_{\text{EH}} = -\mathcal{F} - \frac{e^2}{8\pi^2} \int_{s_0}^{\infty} \frac{ds}{s} e^{-sm^2} \frac{\text{Re} \cosh \left[\hbar e s \sqrt{2(\mathcal{F} + i\mathcal{G})} \right]}{\text{Im} \cosh \left[\hbar e s \sqrt{2(\mathcal{F} + i\mathcal{G})} \right]} \mathcal{G}$$

$$\mathcal{F} := F_{\alpha\beta}^2/4, \quad \mathcal{G} := F^{\alpha\beta} \tilde{F}_{\alpha\beta}/4$$

Heisenberg, Euler (1936)
Schwinger (1950)

IR regulator

$$\sim n_F(p - \mu) + n_F(p + \mu)$$

Euler-Heisenberg theory

- Effective Lagrangian

$$\mathcal{L}_{\text{EH}} = -\mathcal{F} - \frac{e^2}{8\pi^2} \int_{s_0}^{\infty} \frac{ds}{s} e^{-sm^2} \frac{\text{Re} \cosh \left[\hbar e s \sqrt{2(\mathcal{F} + i\mathcal{G})} \right]}{\text{Im} \cosh \left[\hbar e s \sqrt{2(\mathcal{F} + i\mathcal{G})} \right]} \mathcal{G}$$

$$\mathcal{F} := F_{\alpha\beta}^2/4, \quad \mathcal{G} := F^{\alpha\beta} \tilde{F}_{\alpha\beta}/4$$

Heisenberg, Euler (1936)
Schwinger (1950)

UV regulator $\sim y^2$

- current & EM tensor

$$J_{\text{EH}(2)}^{\mu} \Big|_{m \rightarrow 0} = \frac{\mathcal{J}}{6\pi^2} \partial_{\lambda} F^{\mu\lambda} = \underline{2J_{(\partial F)}^{\mu}} \text{ vac} \quad \text{chirality}$$

$$T_{\text{EH}(2)}^{\mu\nu} \Big|_{m \rightarrow 0} = -\frac{\mathcal{J}}{6\pi^2} \left[F^{\mu}_{\sigma} F^{\nu\sigma} - \frac{1}{4} g^{\mu\nu} F_{\alpha\beta}^2 \right] = \underline{2T_{(FF)}^{\mu\nu}} \text{ vac}$$

nonlinear CKT is consistent with EH theory

Physics in logarithm

- effective charge

$$\mathcal{L}_{\text{EH}} = -\frac{1}{4} F_{\mu\nu}^2 \left(1 + \frac{e^2}{12\pi^2} \log \frac{s_0^{-1}}{m^2} \right) + \text{const.} + O(F^4)$$
$$= e_{\text{eff}}^2(s_0)/e^2$$

Schwartz (2014)

EH theory

nonlinear CKT

$$\log s_0^{-1/2}$$



$$\log y^{-1}$$

charge renormalization

$$\beta(e_{\text{eff}}) = M \frac{de_{\text{eff}}(M)}{dM} = \frac{e^3}{12\pi^2}$$

trace anomaly

$$T^\mu{}_\mu = \frac{\beta}{2e^2} F_{\mu\nu}^2 = \frac{e^2}{24\pi^2} F_{\mu\nu}^2$$

Summary

- Formulation of **nonlinear** CKT

- Verification for the nondissipativeness of $J^0 \sim \vec{B} \cdot \vec{\omega}$

- Regularization problem

PV : **no**

DR : **no**

PS : relatively **OK**

- Consistency with Euler-Heisenberg theory

$\log y^{-1}$ \longleftrightarrow charge renormalization, trace anomaly

- Potential developments

meaning of another frame vector?

nonlinear transport by Berry dipole Sodemann, Fu (2015)

collisional effects Kadanoff, Baym (1962)

Chiral plasma instability from dynamical gauge fields

Akamatsu, Yamamoto (2013)