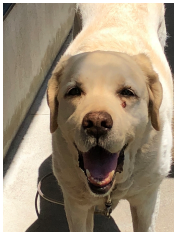


# Range correction in the weak-binding relation for unstable states



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# Background

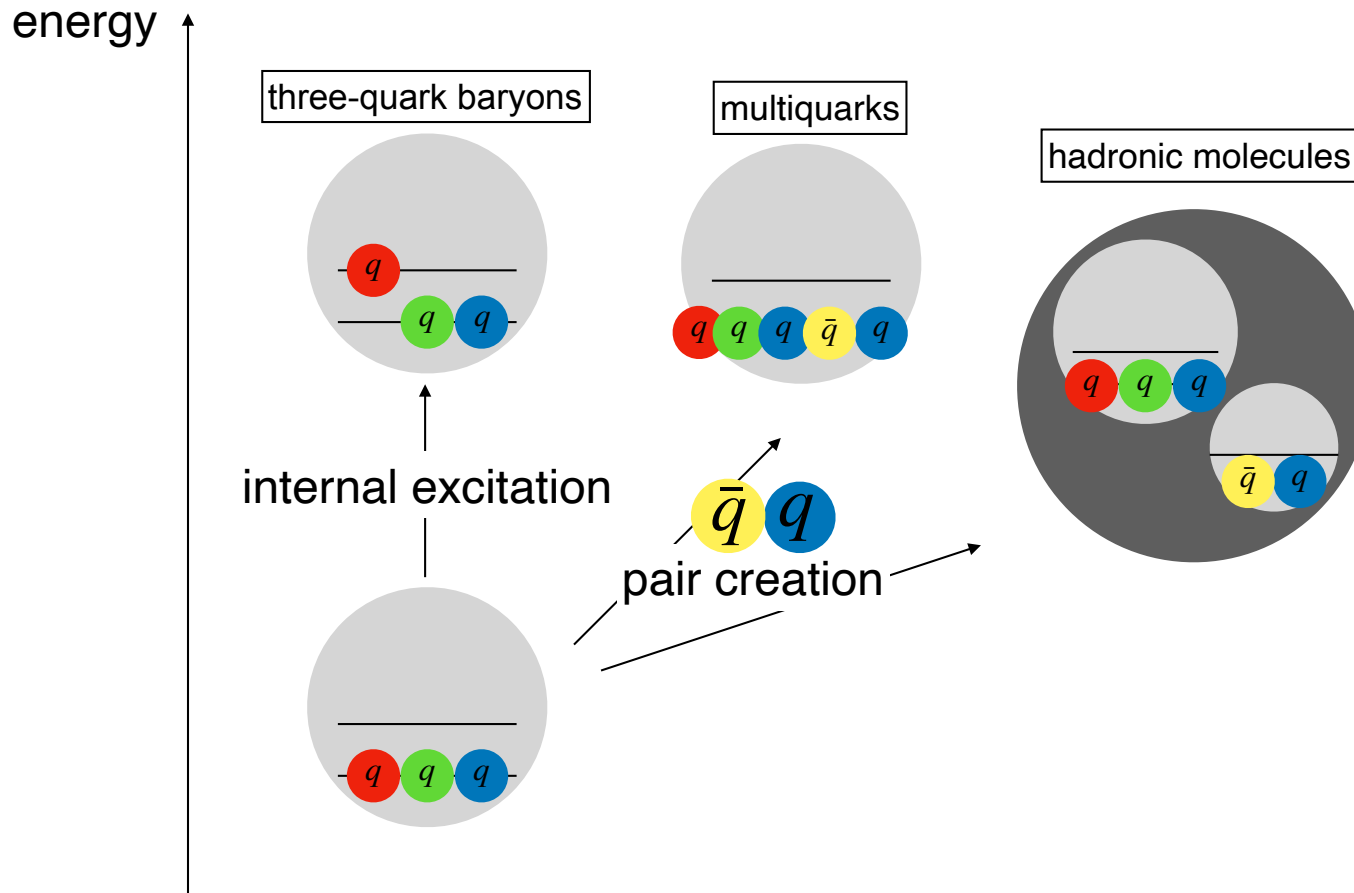
candidates for exotic hadrons

$\Lambda(1405)$ ,  $XYZ$  meson etc...



multiquarks

hadronic molecules



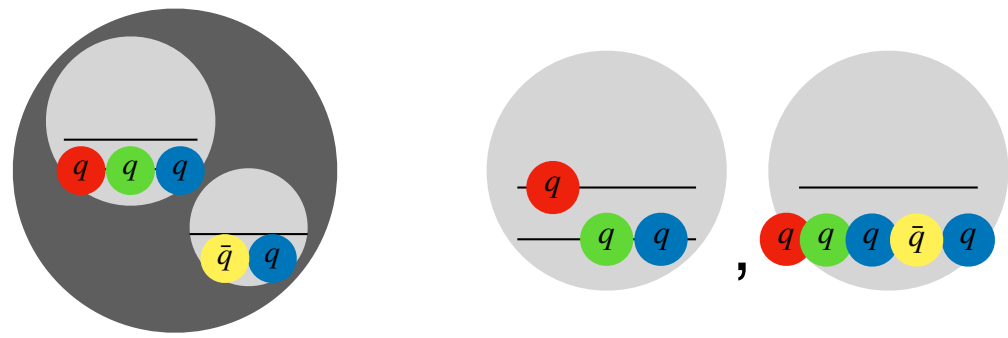
structure of hadrons



model independent

observable

# Previous work



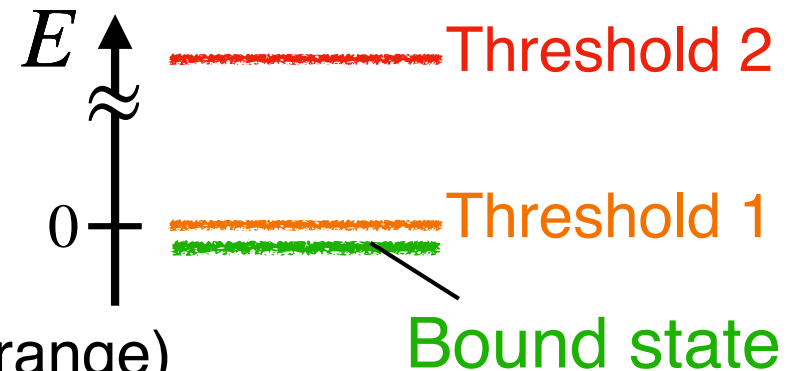
Hadron wave function

$$|\Psi\rangle = \sqrt{X} |\text{hadronic molecule}\rangle + \sqrt{1-X} |\text{others}\rangle$$

Compositeness (weight of hadronic molecule)

## Weak-binding relation for bound state

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}$$



$a_0$  (scattering length)  $R_{\text{typ}}$  (interaction range)

$R \equiv (2\mu B)^{-1/2}$ ,  $B$  (binding energy)

When  $R \gg R_{\text{typ}}$  : observable( $a_0, B$ )  $\longrightarrow$  compositeness( $X$ )

S. Weinberg, Phys. Rev. 137, B672 (1965); Y. Kamiya and T. Hyodo, PTEP 2017, 023D02 (2017).

# Motivation

Weak-binding relation  $a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}$

Low-energy universality  $\rightarrow a_0 = R (R \rightarrow \infty)$

- Deviation by contributions from other channels  $\leftarrow X \neq 1$
- Deviation by interaction range  $\leftarrow R_{\text{typ}} \neq 0$

$\rightarrow$  We study the **range correction** in the weak-binding relation by introducing the effective range  $r_e$ .

# Effective range model

E. Braaten, M. Kusunoki, and D. Zhang, Annals Phys. 323, 1770 (2008), 0709.0499.

Single channel scattering of identical bosons with mass  $m$  :

$$\mathcal{H}_{\text{int}} = \frac{1}{4}\lambda_0(\psi^\dagger\psi)^2 + \frac{1}{4}\rho_0\nabla(\psi^\dagger\psi) \cdot \nabla(\psi^\dagger\psi)$$

Off-shell T-matrix:  $T(E, k, k') = T_1(E) + T_2(E)(k^2 + k'^2) + T_3(E)k^2k'^2$ ,

$$\begin{pmatrix} T_1 & T_2 \\ T_2 & T_3 \end{pmatrix} = -i \begin{pmatrix} \lambda_0 & \rho_0 \\ \rho_0 & 0 \end{pmatrix} - i \begin{pmatrix} \lambda_0 & \rho_0 \\ \rho_0 & 0 \end{pmatrix} \begin{pmatrix} \frac{1}{2} \int \frac{d^3q}{(2\pi)^2} \frac{i}{E - q^2/m + i0^+} & \frac{1}{2} \int \frac{d^3q}{(2\pi)^2} \frac{iq^2}{E - q^2/m + i0^+} \\ \frac{1}{2} \int \frac{d^3q}{(2\pi)^2} \frac{iq^2}{E - q^2/m + i0^+} & \frac{1}{2} \int \frac{d^3q}{(2\pi)^2} \frac{iq^4}{E - q^2/m + i0^+} \end{pmatrix} \begin{pmatrix} T_1 & T_2 \\ T_2 & T_3 \end{pmatrix}.$$

cut off at  $\Lambda$ , typical range  $R_{\text{typ}} \sim 1/\Lambda$

On-shell scattering amplitude:

$$f(k) = \left[ -\frac{8\pi}{m} \frac{\left(1 + \frac{m}{12\pi^2}\Lambda^3\rho_0\right)^2}{N(k)} - \frac{2}{\pi}\Lambda - ik \right]^{-1} N(k) = \left[ \lambda_0 - \frac{m}{20\pi^2}\Lambda^5\rho_0^2 \right] + 2\rho_0 \left( \frac{m}{24\pi^2}\Lambda^3\rho_0 + 1 \right) k^2$$

$$= \left[ -\frac{1}{a_0} + \frac{r_e}{2}k^2 + \mathcal{O}\left(\frac{1}{\Lambda}\right) - ik \right]^{-1} \quad \leftarrow \text{Renormalization}$$

$$\rightarrow \left[ -\frac{1}{a_0} + \frac{r_e}{2}k^2 - ik \right]^{-1} \quad (\Lambda \rightarrow \infty) \quad \leftarrow \text{Zero range limit}$$

# Effective range model

Properties of the effective range model:

-Single channel: | hadronic molecule  $\rangle$  only  $\Leftrightarrow X = 1$

-Zero range limit:  $\Lambda \rightarrow \infty \Leftrightarrow R_{\text{typ}} = 1/\Lambda \rightarrow 0$

$$\Rightarrow a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\} \rightarrow R ?$$

Renormalized scattering amplitude ( $\Lambda \rightarrow \infty$ ):

$$1/f(k = i/R) = 0$$

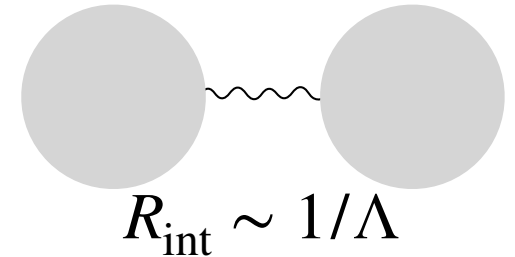
$$a_0 = R \frac{2r_e/R}{1 - (r_e/R - 1)^2} = R \left[ 1 + \mathcal{O}\left(\left|\frac{r_e}{R}\right|\right) \right] \Rightarrow a_0 \neq R$$

→ range correction in the weak-binding relation form  $r_e$

# Improved weak-binding relation

Weak-binding relation  $a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}$

interaction range:  $R_{\text{typ}} \longrightarrow R_{\text{int}} \sim 1/\Lambda$



Redefinition of  $R_{\text{typ}}$ :

$$R_{\text{typ}} = \max \left\{ R_{\text{int}}, R_{\text{eff}} \right\},$$

$$R_{\text{eff}} = \max \left\{ |r_e|, \frac{|P_s|}{R^2}, \dots \right\}.$$

$$f(k) = \left[ -\frac{1}{a_0} + \frac{r_e}{2}k^2 - \frac{P_s}{4}k^4 + \dots - ik \right]^{-1}$$

Length scale in the effective range expansion except for  $a_0$

# Numerical calculation

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}$$

When dose the weak-binding relation work?

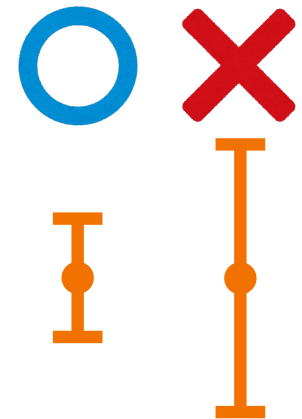
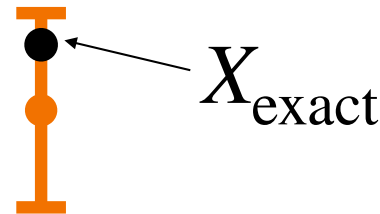
Estimation with correction terms ( $\xi \equiv R_{\text{typ}}/R$ ): Y. Kamiya and T. Hyodo, PTEP 2017, 023D02 (2017).

Central value: 
$$X_c = \frac{a_0/R}{2 - a_0/R}$$

$$X_{\text{upper}}(\xi) = \frac{a_0/R}{2 - a_0/R} + \xi, \quad X_{\text{lower}}(\xi) = \frac{a_0/R}{2 - a_0/R} - \xi.$$

Weak-binding relation works when...

$$\left\{ \begin{array}{l} X_{\text{lower}} < X_{\text{exact}} < X_{\text{upper}} \\ \longrightarrow \text{Validity condition} \\ (X_{\text{upper}} - X_c)/X_c < 0.1 \text{ and } (X_c - X_{\text{lower}})/X_c < 0.1 \\ \longrightarrow \text{Precision condition} \end{array} \right.$$





# Numerical calculation

Effective range model ( $\Lambda < \infty$ )

$$f(k; \lambda_0, \rho_0, \Lambda) = \left[ -\frac{1}{a_0} + \frac{r_e}{2}k^2 + \mathcal{O}(R_{\text{int}}) - ik \right]^{-1} \text{ (two length scales } r_e \text{ and } R_{\text{int}})$$

$$1/f(k = i/R) = 0$$

-  $r_e \neq 0$  (range correction):  $\xi_{r_e} = |r_e/R| \longrightarrow$  Uncertainty from  $r_e$

$r_e < 0$  (effective range model)

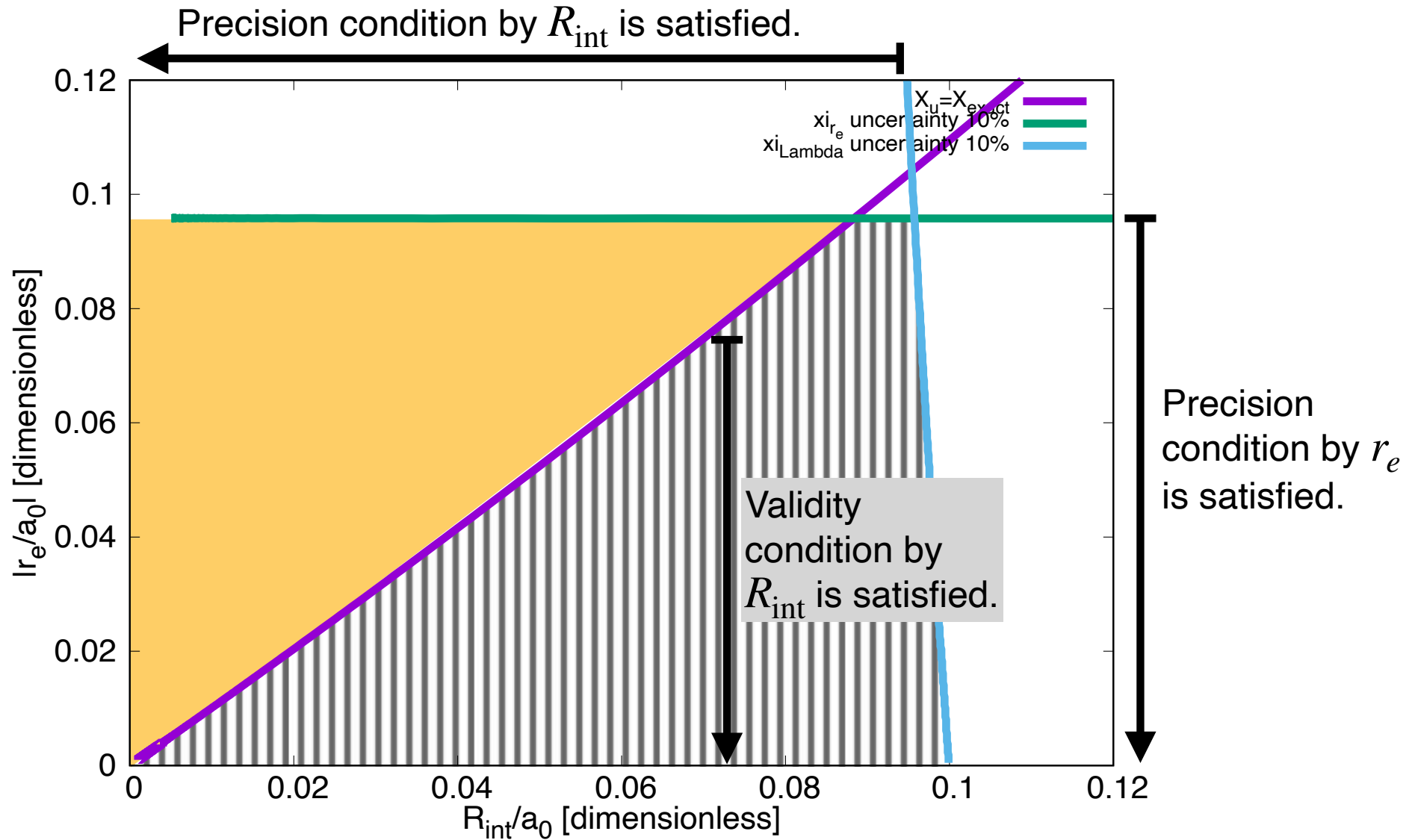
-  $R_{\text{int}} = 1/\Lambda \neq 0$ :  $\xi_{\text{int}} = R_{\text{int}}/R. \longrightarrow$  Uncertainty from  $R_{\text{int}}$

-  $X_{\text{exact}} = 1$

$\longrightarrow$  We search for the regions of  $r_e$  and  $R_{\text{int}}$  in which validity and precision conditions are satisfied.

# Numerical calculation

Validity and precision conditions in  $R_{\text{int}}/a_0 - |r_e/a_0|$  plane



Only the improved weak-binding relation can be applied.

# Extension to unstable states

Candidates for exotic hadrons are unstable states!

Unstable systems  $\longrightarrow E, a_0, r_e \in \mathbb{C}$

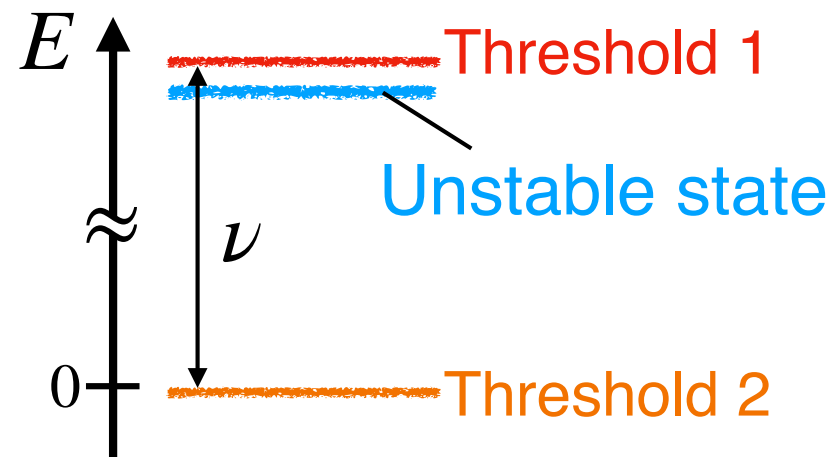
## Weak-binding relation for **unstable states**

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) + \mathcal{O}\left(\left|\frac{l}{R}\right|^3\right) \right\}$$

$$l \equiv (2\mu\nu)^{-1/2}$$

$\nu$  (difference between the threshold energies)

$$R \equiv (2\mu E)^{-1/2}, E \text{ (eigenenergy)}$$



When  $|R| \gg R_{\text{typ}}$  and  $|R| \gg l$ : observables  $(a_0, E)$

$\longrightarrow$  compositeness( $X$ )

Interpretation of  $X \longrightarrow \tilde{X} \equiv \frac{1 - |1 - X| + |X|}{2} \in \mathbb{R} \quad \because X \in \mathbb{C}$

# Extension to unstable states

What are the features of weak-binding relation in the effective range model with zero range limit ( $\Lambda \rightarrow \infty$ )?

$$f(k) = \left[ -\frac{1}{a_0} + \frac{r_e}{2}k^2 - ik \right]^{-1} \quad \text{One length scale } r_e \text{ except for } a_0$$

$\longrightarrow R_{\text{typ}} = |r_e|$

- description of unstable states in the **single channel** case

$a_0, r_e \in \mathbb{C} \longrightarrow \text{pole of } f(k) \in \mathbb{C} \longrightarrow \text{eigenenergy } E \in \mathbb{C}$

**unstable states**

$$a_0 = |a_0| e^{i\theta_{a_0}}, r_e = -|r_e| e^{i\theta_{r_e}}$$

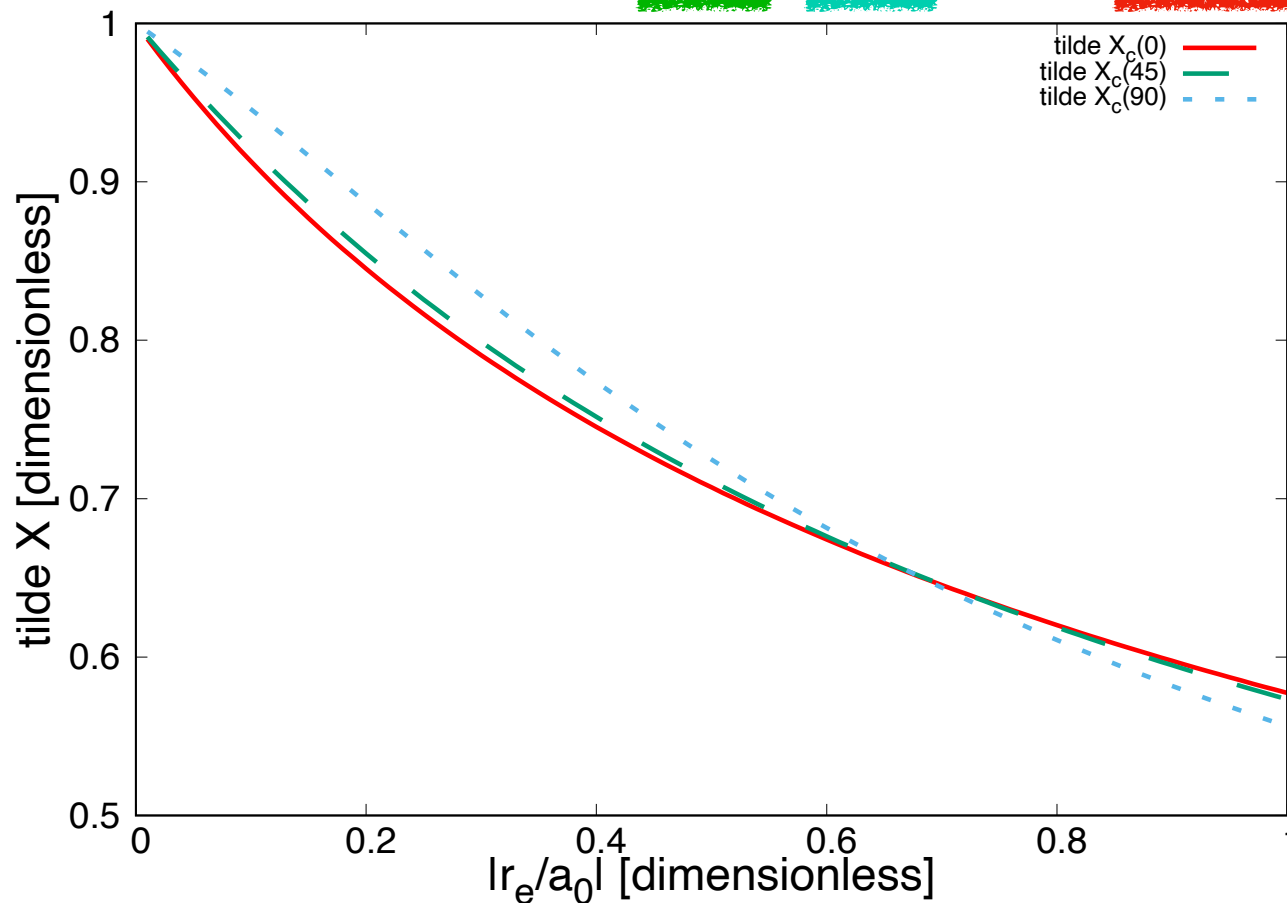
-  $X(a_0, r_e) = X(a_0/r_e)$  in the zero range limit

$\longrightarrow X$  depends only on  $\theta \equiv \theta_{a_0} - \theta_{r_e}$ .

$$a_0 = |a_0| e^{i\theta}, r_e = -|r_e|$$

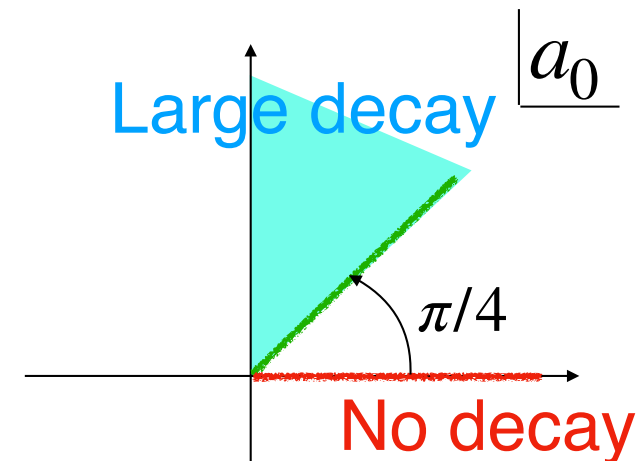
# Extension to unstable states

$\tilde{X}$  for unstable ( $\theta = \pi/4, \pi/2$ ) and stable ( $\theta = 0$ ) states



$\theta = \pi/4, \pi/2$

➔ Far from the stable state



There are almost no differences between  $\tilde{X}$  for unstable and stable states.

➔ Is this a feature of effective range model in the zero range limit?

# Extension to unstable states

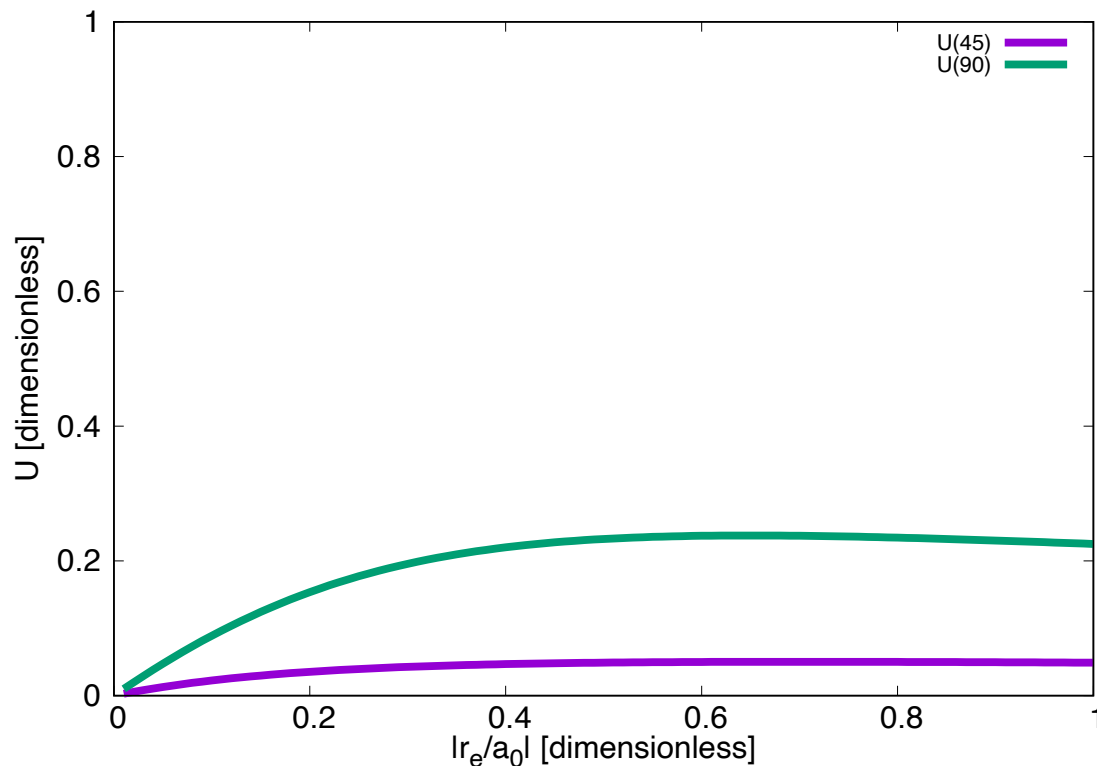
Estimation of uncertainty of interpretation ( $U$ ):

$$U \equiv |1 - X| + |X| - 1. \quad \text{Y. Kamiya and T. Hyodo, PTEP 2017, 023D02 (2017).}$$

-When  $U$  is small, we can interpret  $\tilde{X}$  as compositeness and the system is similar to the stable state.

-When  $X \in \mathbb{R}$ ,  $U = 0$ .

$$U (\theta = \pi/2 \text{ and } \pi/4)$$



$$U \lesssim 0.25 \text{ in } \theta = \pi/2 \text{ case}$$

Far from the stable state

$\because a_0 = |a_0| e^{i\theta}$  is pure imaginary

This system has similar features to the stable state.

$\because X$  depends only on  $a_0/r_e$ .

# Conclusion and future prospect

- Weak-binding relation : observable  compositeness (X)

$$a_0 = R \left\{ \frac{2X}{1+X} + \mathcal{O}\left(\frac{R_{\text{typ}}}{R}\right) \right\}$$

- We study the range correction in weak-binding relation from  $r_e$ .
- Improved weak-binding relation by redefinition of  $R_{\text{typ}}$  :

$$R_{\text{typ}} = \max \left\{ R_{\text{int}}, |r_e|, \dots \right\}$$

- We find the region where only the improved weak-binding relation can be applied.
- In effective range model in the zero range limit, the unstable state is similar to the stable state.
- Future prospect: Apply the improved relation to hadron systems.