

# Compositeness of exotic hadrons with decay and coupled-channel effects



arXiv:2303.07038 [hep-ph]



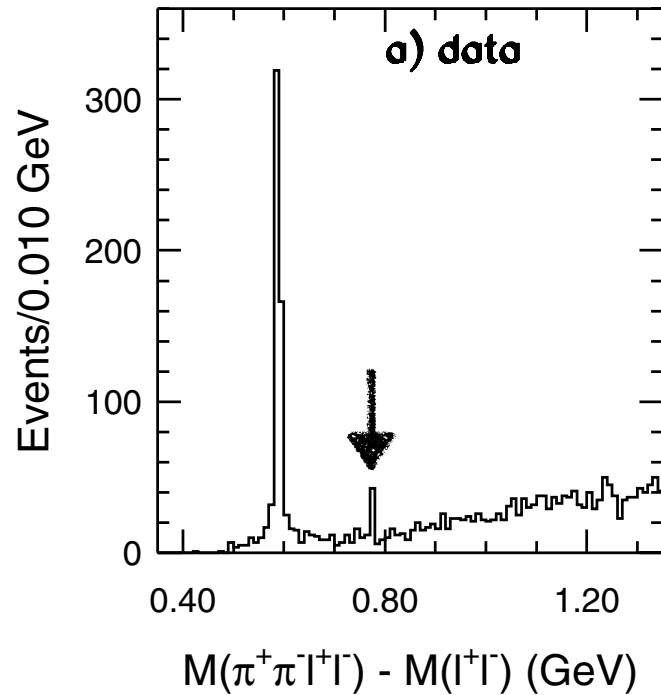
**Tomona Kinugawa**

**Tetsuo Hyodo**

Department of Physics, Tokyo Metropolitan University  
March. 16th, 3rd J-PARC HEF-ex WS

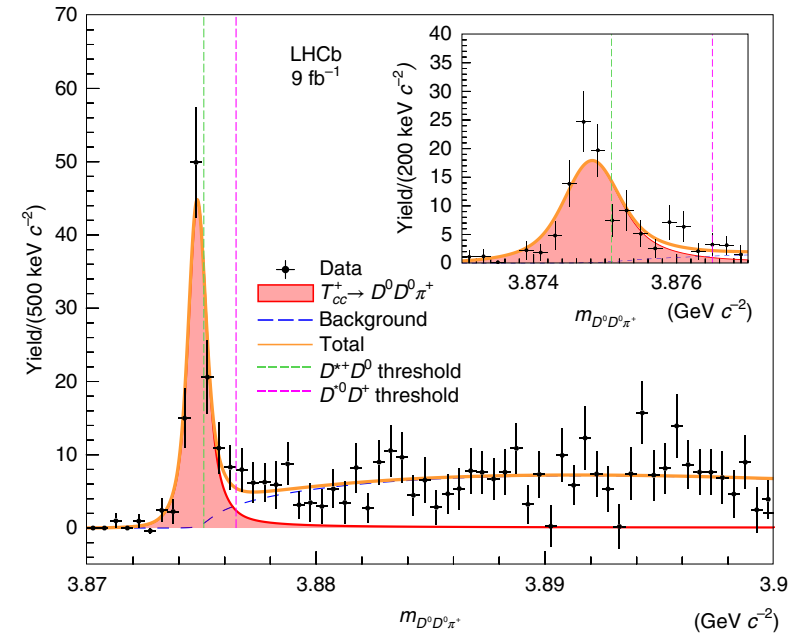
# Near-threshold exotic hadrons

$$X(3872) \rightarrow \pi^+ \pi^- J/\psi$$



S. K. Choi *et al.* (Belle), Phys. Rev. Lett. **91**, 262001 (2003).

$$T_{cc} \rightarrow D^0 D^0 \pi^+ (cc\bar{u}\bar{d})$$



LHCb Collaboration, Nature Phys. **18** (2022) no.7, 751-754;

LHCb Collaboration, Nat. Commun. **13** 3351 (2022).

internal structure?

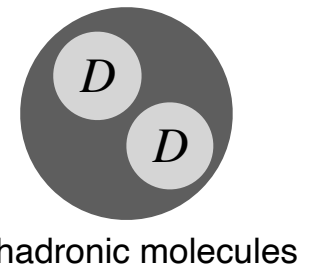
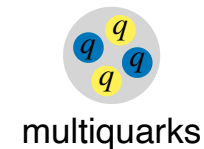
exotic hadron

$\neq qq\bar{q}$  or  $q\bar{q}$



multiquarks

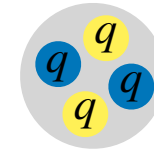
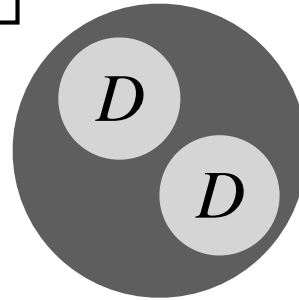
hadronic molecules



# Compositeness

## ● definition

hadron wavefunction



$$|\Psi\rangle = \sqrt{X} |\text{hadronic molecule}\rangle + \sqrt{1-X} |\text{others}\rangle$$

compositeness

elementarity

T. Hyodo, Int. J. Mod. Phys. A 28, 1330045 (2013).

$$\begin{aligned} * 0 \leq X \leq 1 &\longrightarrow X > 0.5 \Leftrightarrow \text{composite dominant} \\ &X < 0.5 \Leftrightarrow \text{elementary dominant} \end{aligned}$$

## ● model calculation

T. Hyodo, D. Jido, and A. Hosaka, Phys. Rev. C **85**, 015201 (2012);

F. Aceti and E. Oset, Phys. Rev. D **86**, 014012 (2012).

compositeness  $X$  ← residue of scattering amplitude

# Low-energy universality

4

- scattering length  $a_0 \gg$  typical length scale of system  $R_{\text{typ}}$

## low-energy universality

E. Braaten and H.-W. Hammer, Phys. Rept. **428**, 259 (2006) ;  
F. P. Naidon and S. Endo, Rept. Prog. Phys. **80**, 056001 (2017).

→ length scales are written only by  $|a_0|$  ( $\rightarrow \infty$ )

for bound states ?

$$a_0 = R \quad R = 1/\sqrt{2\mu B} \quad a_0 \rightarrow \infty \longrightarrow B \rightarrow 0$$

→ universality holds for weakly-bound states!!

- compositeness  $X = 1$  in  $B \rightarrow 0$  limit T. Hyodo, Phys. Rev. C **90**, 055208 (2014) .

→ near threshold poles = composite dominant ?

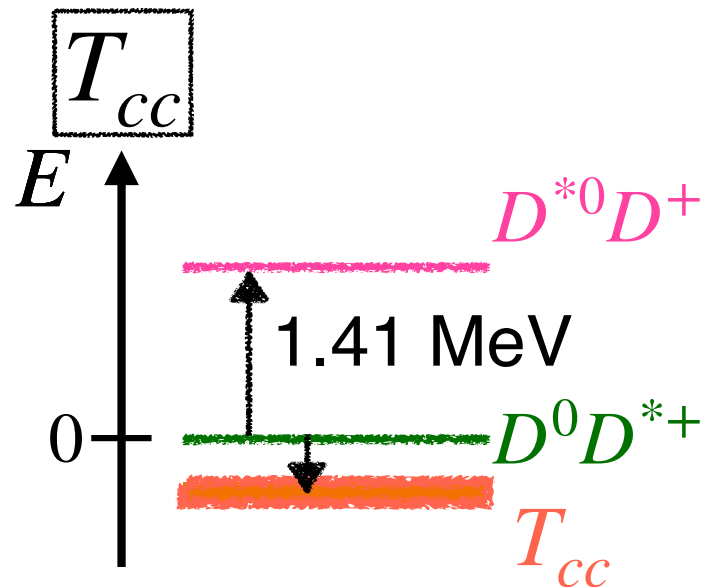
e.g.  ${}^8\text{Be}$ ,  ${}^{12}\text{C}$  →  $\alpha$  cluster? H. Horiuchi, K. Ikeda, and Y. Suzuki,  
Prog. Theor. Phys. Suppl. **52**, 89 (1972) .

# Decay & coupled ch. effects

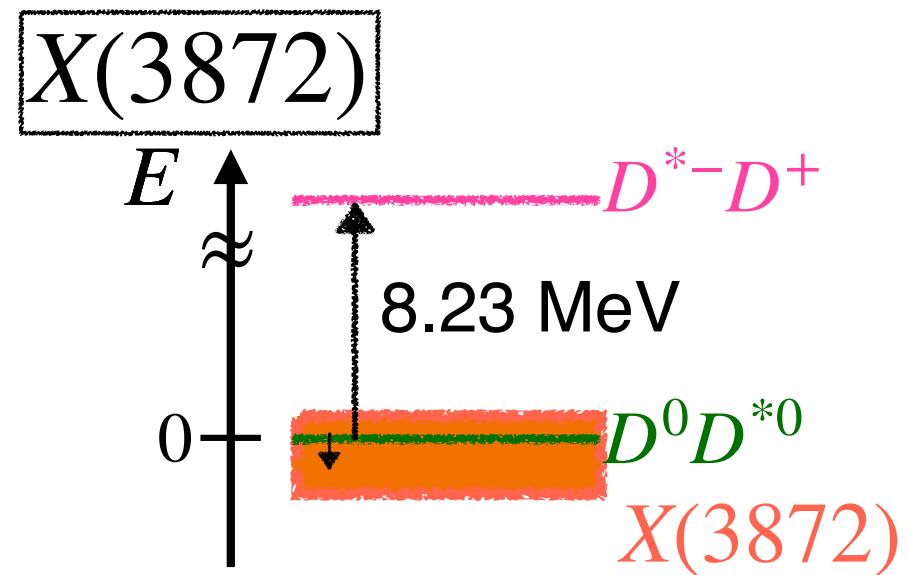
5

However...

actual exotic hadrons  $\longrightarrow$  decay and coupled channel



LHCb Collaboration, Nat. Commun **13** 3351 (2022).



PDG

other ch. than threshold ch. make deviation from  $X = 1$

Y. Kamiya and T. Hyodo, PTEP 2017, 023D02 (2017).

This work...

study those deviations quantitatively!

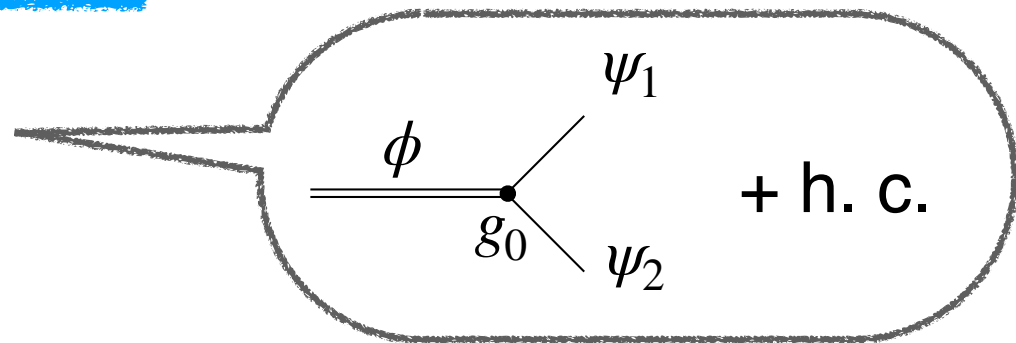
## ● single-channel resonance model

$$\mathcal{H}_{\text{free}} = \frac{1}{2m_1} \nabla \psi_1^\dagger \cdot \nabla \psi_1 + \frac{1}{2m_2} \nabla \psi_2^\dagger \cdot \nabla \psi_2 + \frac{1}{2m_\phi} \nabla \phi^\dagger \cdot \nabla \phi + \nu_0 \phi^\dagger \phi,$$

1.

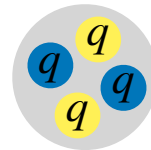
$$\mathcal{H}_{\text{int}} = g_0 (\phi^\dagger \psi_1 \phi_2 + \phi_1^\dagger \psi_2^\dagger \phi).$$

2.



1. single-channel scattering

2. coupling with compact state  $\phi$



## ● scattering amplitude

$$V = \frac{g_0^2}{E - \nu_0}, \quad G = -\frac{\mu}{\pi^2} \left[ \Lambda + ik \arctan\left(\frac{\Lambda}{-ik}\right) \right]. \quad \Lambda : \text{cutoff}$$

$$T = \frac{1}{V^{-1} - G} \longrightarrow f(k) = -\frac{\mu}{2\pi} \left[ \frac{\frac{k^2}{2\mu} - \nu_0}{g_0^2} + \frac{\mu}{\pi^2} \left[ \Lambda + ik \arctan\left(\frac{\Lambda}{-ik}\right) \right] \right]^{-1}.$$

# Model scales and parameters

7

- typical energy scale :  $E_{\text{typ}} = \Lambda^2/(2\mu)$

- three model parameters  $g_0, \nu_0, \Lambda$

1. calculation with given  $B$

$$\text{coupling const. } g_0 : g_0^2(B, \nu_0, \Lambda) = \frac{\pi^2}{\mu} (B + \nu_0) \left[ \Lambda - \kappa \arctan(\Lambda/\kappa) \right]^{-1}$$

$$\because \text{bound state condition } f^{-1} = 0 \quad \kappa = \sqrt{2\mu B}.$$

2. use dimensionless quantities with  $\Lambda$

→ results do not depend on cutoff  $\Lambda$

3. energy of bare quark state  $\nu_0$

varied in the region :  $-B/E_{\text{typ}} \leq \nu_0/E_{\text{typ}} \leq 1$

$\because$  to have  $g_0^2 \geq 0$  & applicable limit of EFT

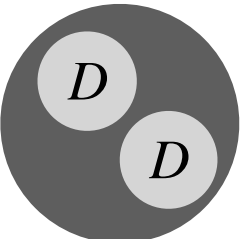
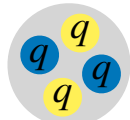
# Calculation

## ● compositeness $X$

scattering amplitude :  $T = \frac{1}{V^{-1} - G}$  Y. Kamiya and T. Hyodo, PTEP 2017, 023D02 (2017).

$$\begin{aligned} \longrightarrow X &= \frac{G'(-B)}{G'(-B) - [V^{-1}(-B)]'}, \quad \alpha'(E) = d\alpha/dE \\ &= \left[ 1 + \frac{\pi^2 \kappa}{g_0^2 \mu^2} \left( \arctan(\Lambda/\kappa) - \frac{\Lambda/\kappa}{1 + (\Lambda/\kappa)^2} \right)^{-1} \right]^{-1}. \end{aligned}$$

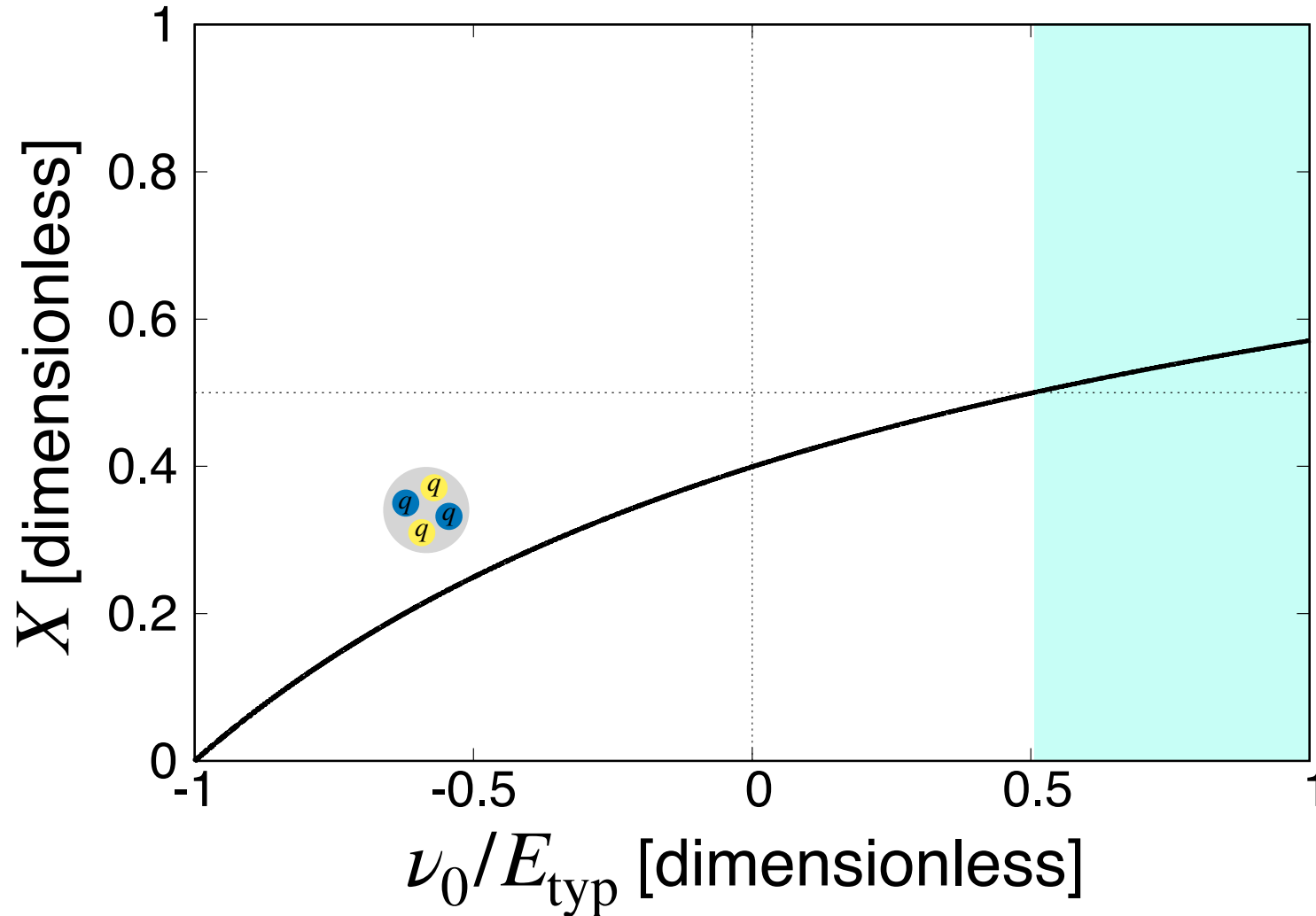
-  $\nu_0$  region :  $-B/E_{\text{typ}} \leq \nu_0/E_{\text{typ}} \leq 1$

compositeness  $X$  as a function of  $\nu_0$   or   $X < 0.5$

$\longrightarrow$  internal structure of bound state?



●  $X$  as a function of  $\nu_0/E_{\text{typ}}$  of bound state  $B = E_{\text{typ}}$



$X > 0.5$  :

$X < 0.5$  :

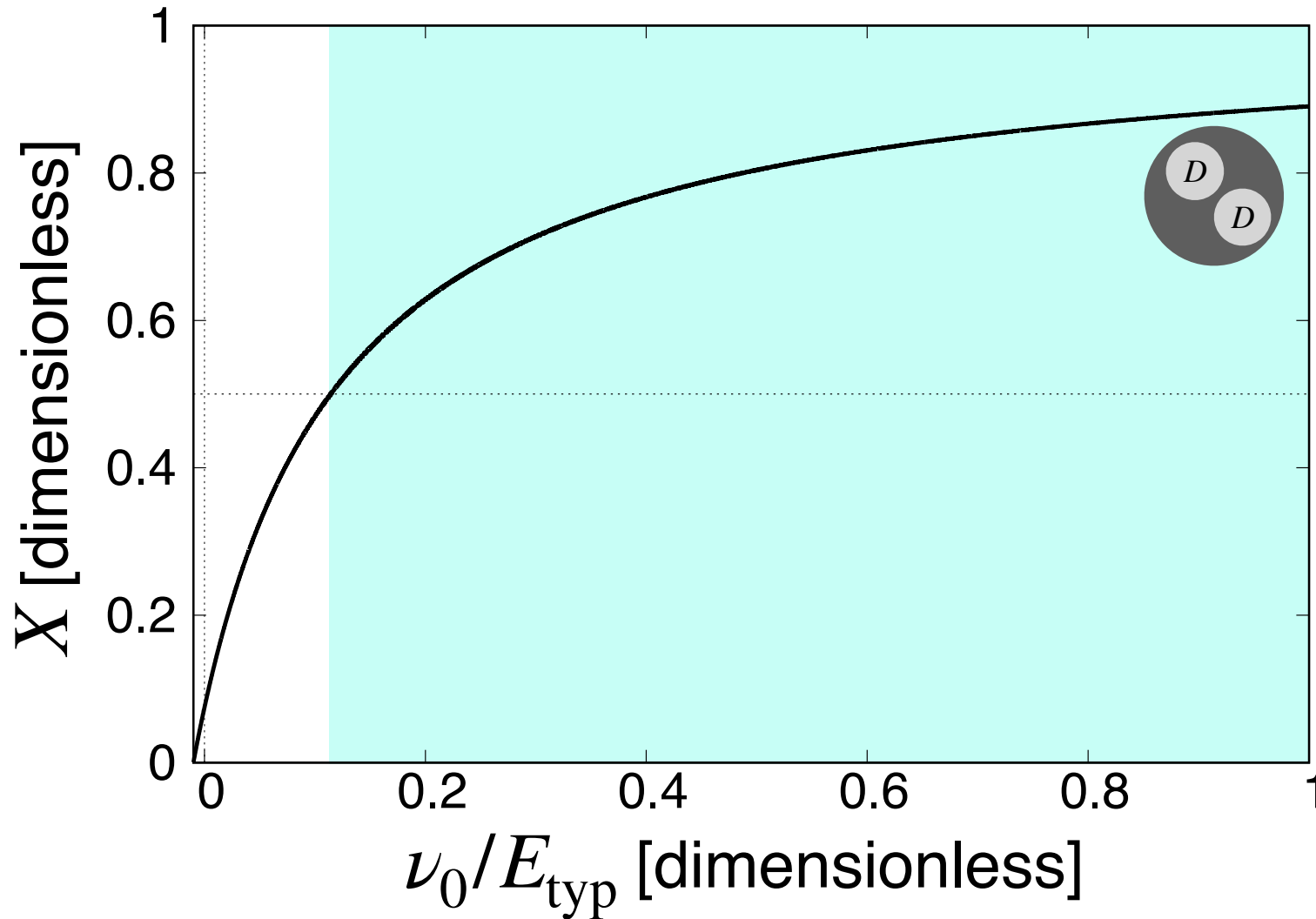
- typical energy scale :  $B = E_{\text{typ}} = \Lambda^2/(2\mu)$

-  $X > 0.5$  only for 25 % of  $\nu_0$  = elementary dominant

$\therefore$  bare state origin

●  $X$  as a function of  $\nu_0/E_{\text{typ}}$  of bound state  $B = 0.01E_{\text{typ}}$

10

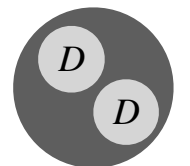


$X > 0.5$  :

$X < 0.5$  :

- weakly-bound state :  $B = 0.01E_{\text{typ}}$

-  $X > 0.5$  for 88 % of  $\nu_0$  = composite dominant



**$\therefore$  low-energy universality !**

# Effect of decay

## ● introducing decay effect

- formally : introducing decay channel in lower energy region than binding energy

→ eigenenergy becomes complex

- effectively : coupling const.  $g_0 \in \mathbb{C}$  ! ← this work

$$E = -B \longrightarrow E = -B - \underline{i\Gamma/2}$$

compositeness

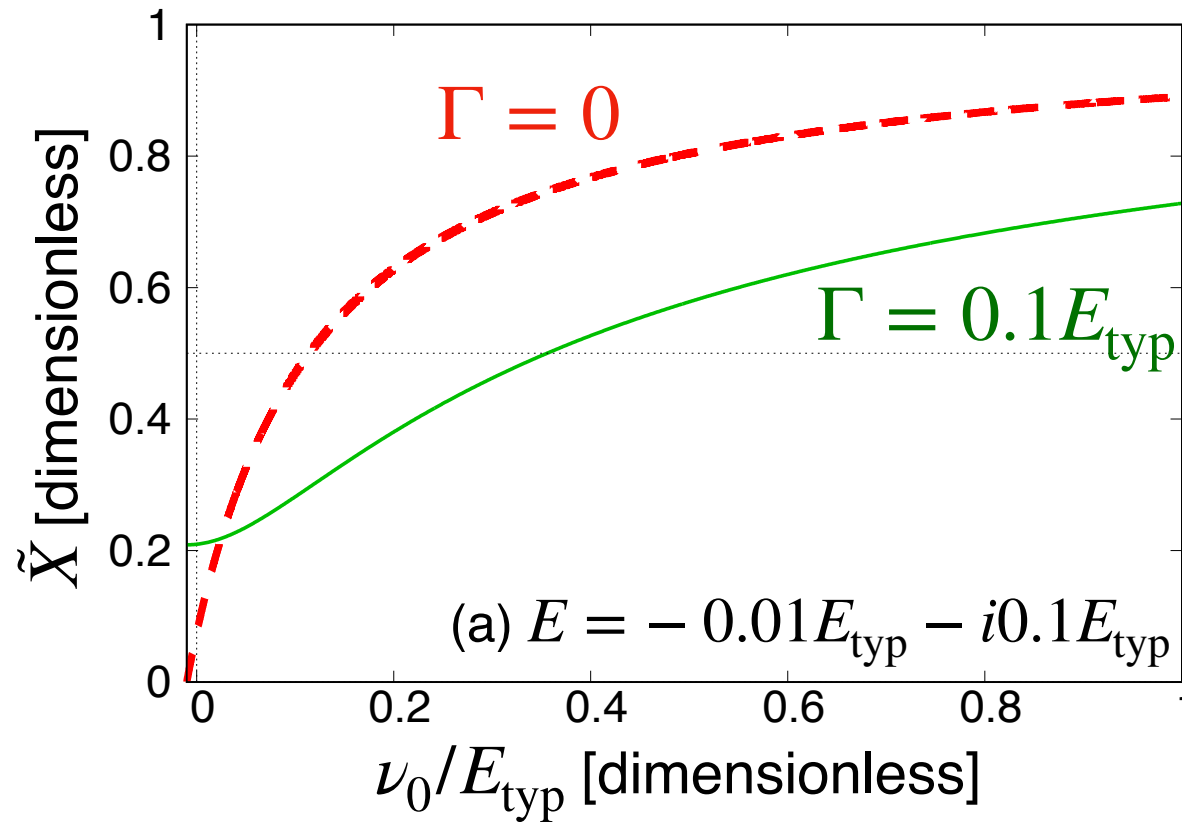
$$X \in \mathbb{R} \longrightarrow X \in \mathbb{C}$$

$$\tilde{X} = \frac{|X|}{|X| + |1 - X|}$$

T. Sekihara, *et. al.*, PRC 93, 035204 (2016).

# Effect of decay

● compositeness  $\tilde{X}$



- $\tilde{X}$  is suppressed by  $\Gamma \neq 0$ 
  - $\because$  threshold ch. component ( $\tilde{X}$ ) decreases with inclusion of decay ch. component ( $1 - \tilde{X}$ )

# Effect of coupled channel

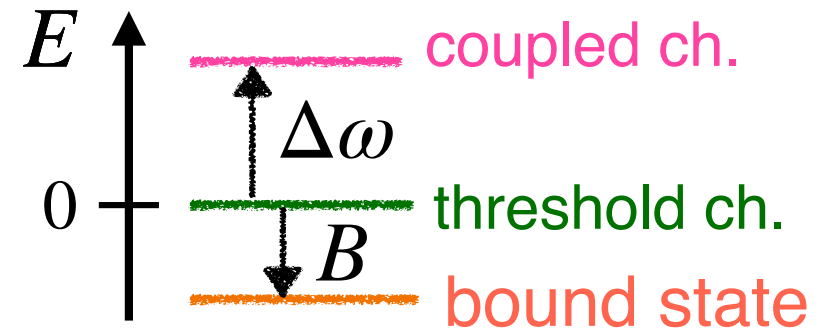
● introducing coupled channel  $\Psi_1 \Psi_2$

$$|\Psi\rangle = \sqrt{X_1} |\text{threshold ch}\rangle + \sqrt{X_2} |\text{coupled ch}\rangle + \sqrt{1 - (X_1 + X_2)} |\text{others}\rangle$$

$X_1$  : threshold ch. compositeness

$X_2$  : coupled ch. compositeness

- threshold energy difference  $\Delta\omega$



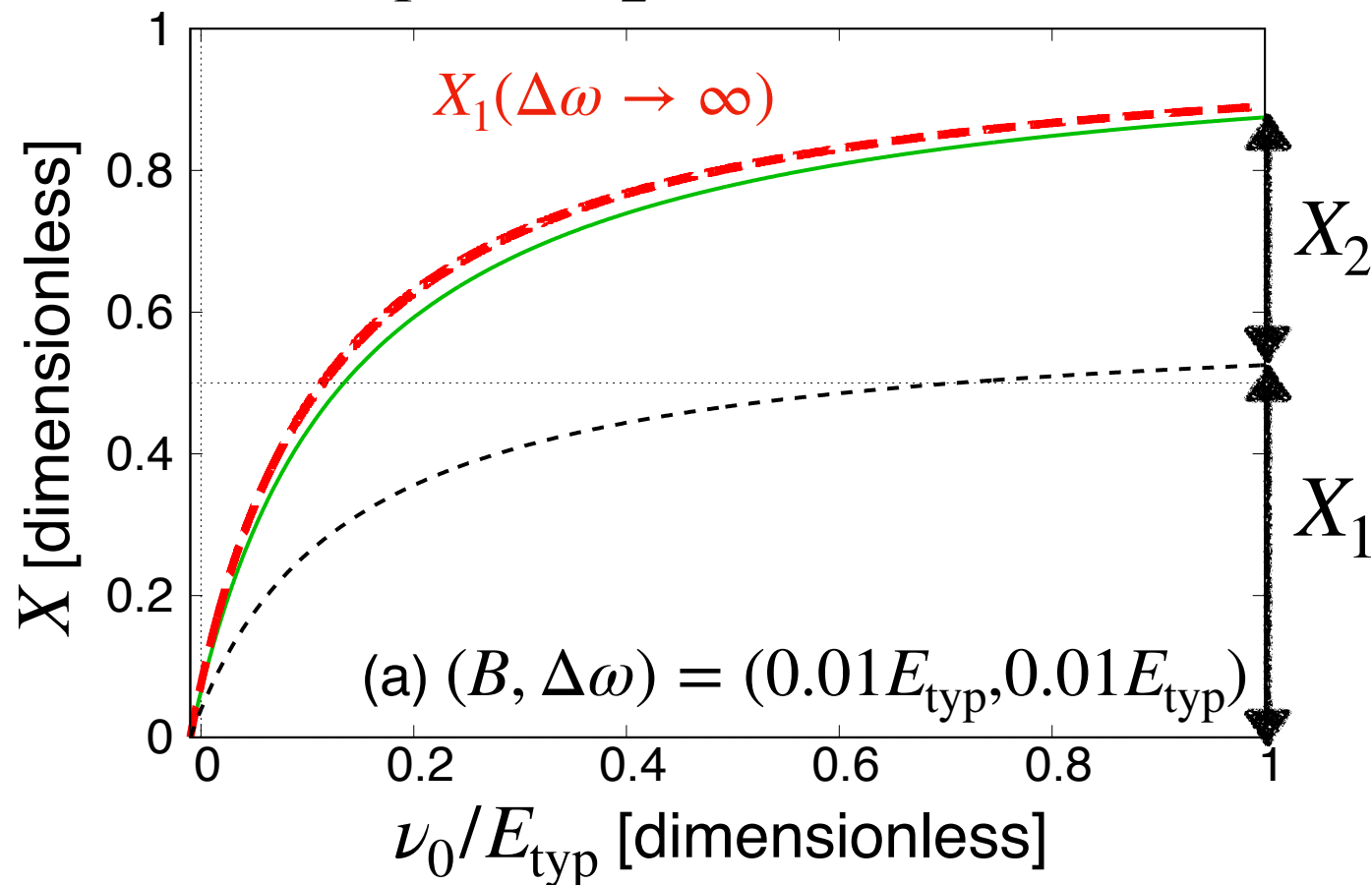
- low-energy universality with coupled-channel effect

$X_1 \sim 1$  (threshold channel)

$X_2 \sim 0$  (other channel)  $Z \sim 0$  (other channel)

# Effect of coupled channel

● compositeness  $X_1$  and  $X_2$

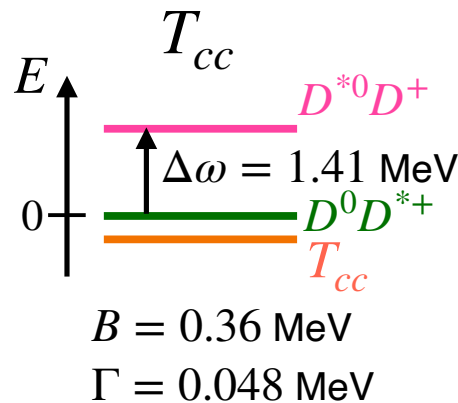


- $X_1$  is suppressed by  $\Delta\omega$ 
  - $\therefore$  threshold ch. component ( $X_1$ ) decreases with inclusion of coupled ch. component ( $X_2$ )

# Application to $T_{cc}$ and $X(3872)$

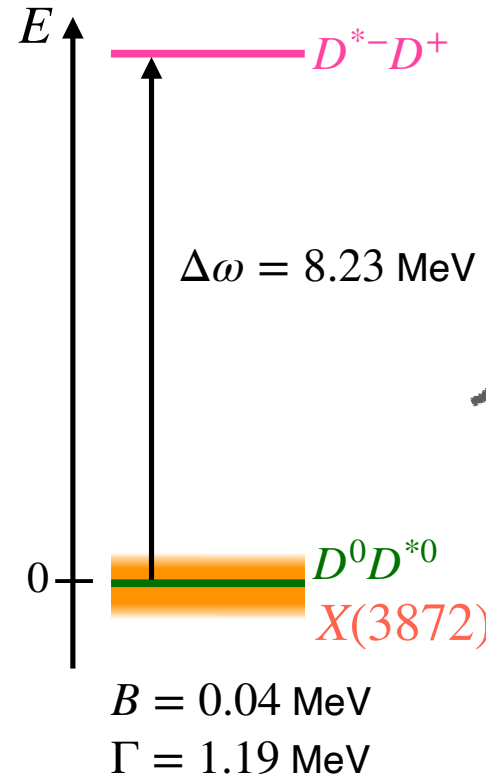
● exotic hadron ← decay and coupled channel  
 $X(3872)$

small  
 $\Gamma$  and  $\Delta\omega$



LHCb Collaboration, Nat. Commun **13** 3351 (2022).

large  
 $\Gamma$  and  $\Delta\omega$



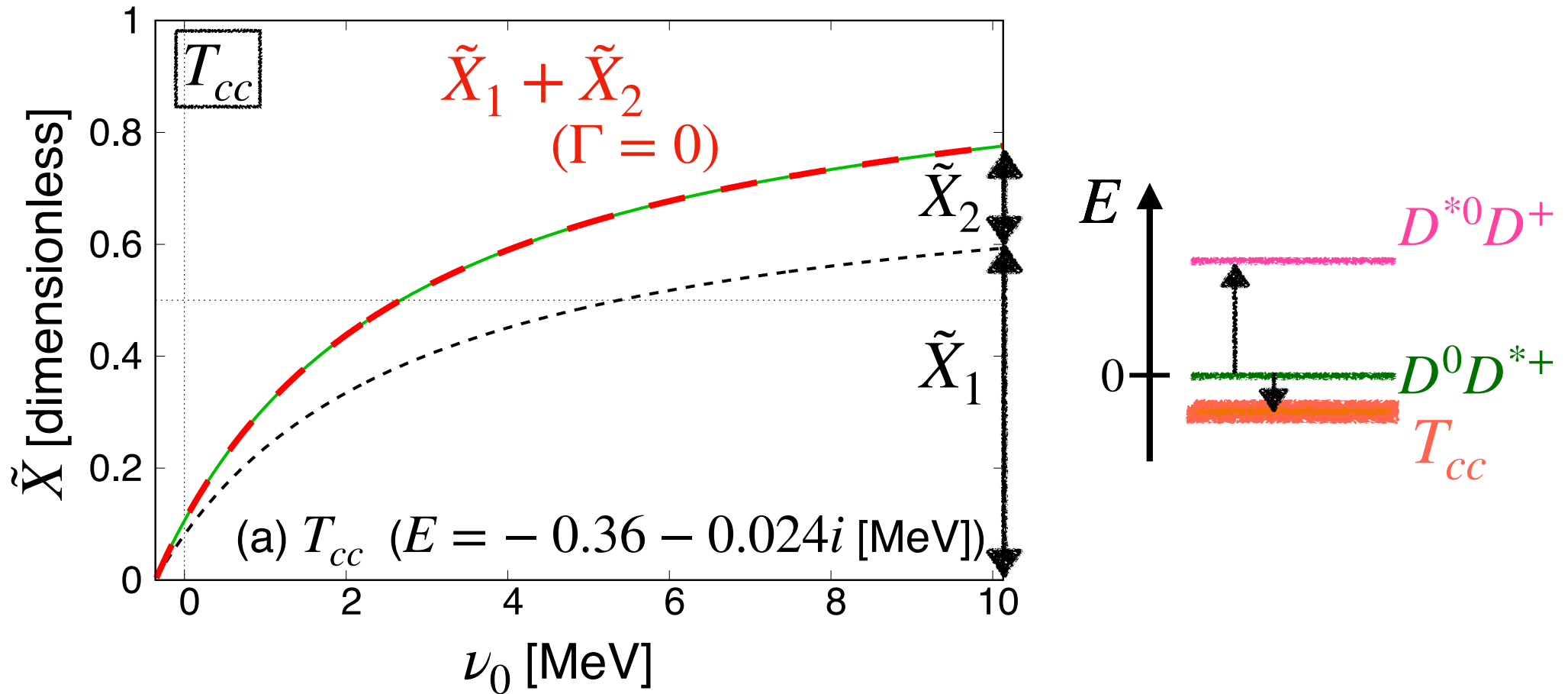
PDG

● compositeness T. Sekihara, *et. al.*, PRC 93, 035204 (2016).

$$\tilde{X}_j = \frac{|X_j|}{\sum_j |X_j| + |Z|}, \quad (j = 1, 2)$$

$\tilde{X}_1$  : threshold ch. compositeness  
 $\tilde{X}_2$  : coupled ch. compositeness

# Application to $T_{cc}$ and $X(3872)$



-  $\tilde{X}_2$  is not negligible

∴ coupled ch. contribution (small  $\Delta\omega$ )

- difference of  $\tilde{X}_1 + \tilde{X}_2(\Gamma = 0)$  and  $\tilde{X}_1 + \tilde{X}_2$  is too small

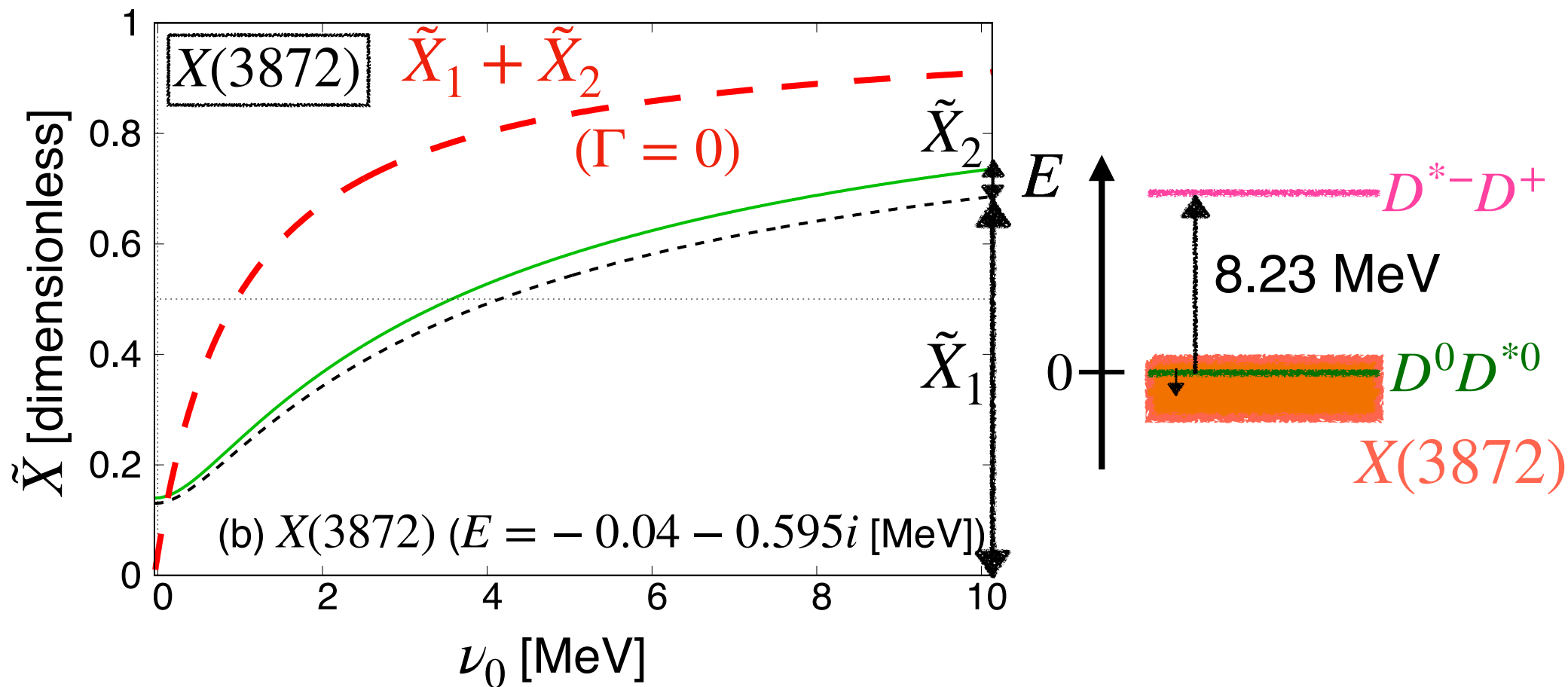
→ We can neglect decay contribution

∴  $\Gamma \ll B$



# Application to $T_{cc}$ and $X(3872)$

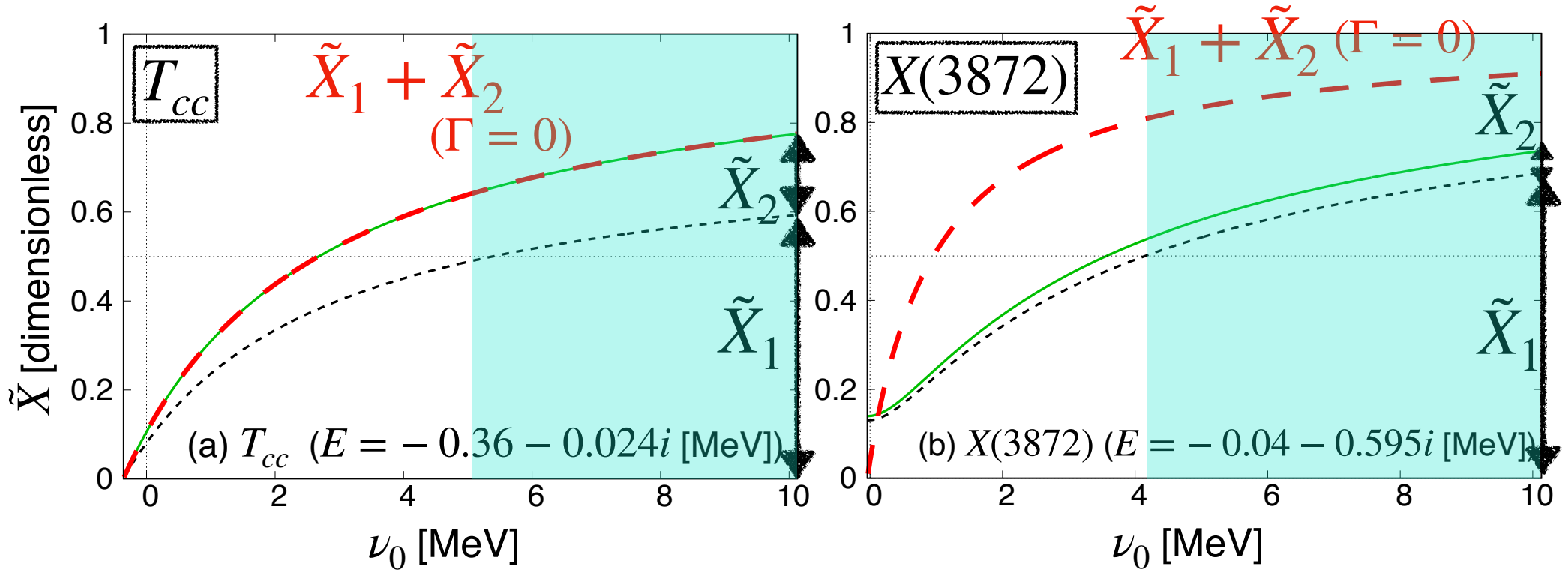
17



- difference of  $\tilde{X}_1 + \tilde{X}_2(\Gamma = 0)$  and  $\tilde{X}_1 + \tilde{X}_2$  is large
  - $\therefore$  large decay width contribution
- $\tilde{X}_2$  is much smaller than  $\tilde{X}_1$ 
  - $\longrightarrow$  coupled ch. effect is small

# Application to $T_{cc}$ and $X(3872)$

18



- $T_{cc}$  :  $\tilde{X}_1 > 0.5$  for 55 % of  $\nu_0$  region
- $X(3872)$  :  $\tilde{X}_1 > 0.5$  for 41 % of  $\nu_0$  region
- coupled ch. effect is more important for  $T_{cc}$  than  $X(3872)$
- decay effect is more important for  $X(3872)$  than  $T_{cc}$

# Summary

arXiv:2303.07038 [hep-ph]

19

- internal structure of exotic hadrons ← EFT & compositeness
- model with bare state coupled to the scattering state
- shallow bound state is composite dominant even from bare state
  - ∴ low-energy universality
- decay and coupled channel effects are introduced
  - both decay and coupled ch. effect suppress compositeness
- $X$  of  $T_{cc}$  and  $X(3872)$  are calculated with decay and coupled ch. effects

$T_{cc}$  : important coupled ch. effect with negligible decay effect

$X(3872)$  : important decay effect with negligible coupled ch. effect

# Compositeness of exotic hadrons with decay and coupled-channel effects



arXiv:2303.07038 [hep-ph]



**Tomona Kinugawa**

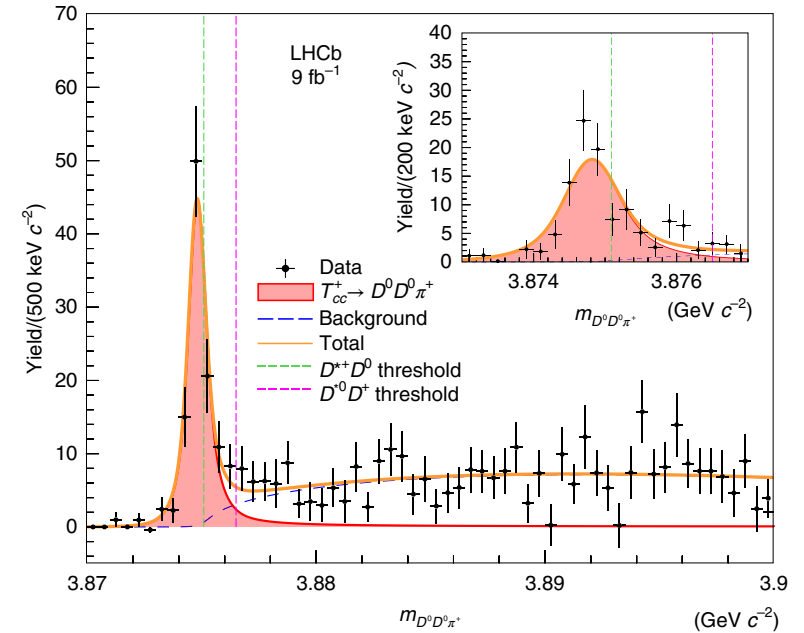
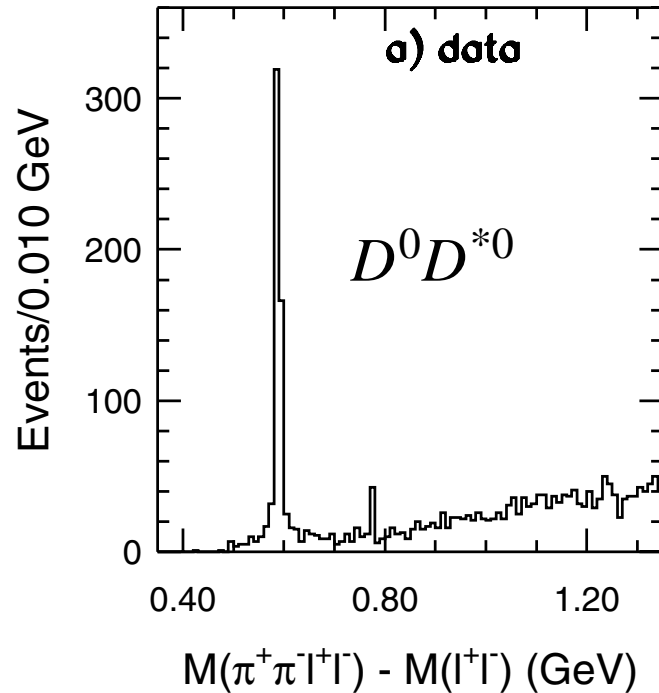
**Tetsuo Hyodo**

Department of Physics, Tokyo Metropolitan University  
March. 16th, 3rd J-PARC HEF-ex WS

# Near-threshold exotic hadrons

$$X(3872) \rightarrow K^\pm \pi^+ \pi^- J/\psi$$

$$T_{cc} \rightarrow D^0 D^0 \pi^+ (c\bar{u}c\bar{u}d\bar{d})$$



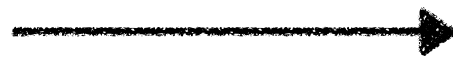
S. K. Choi *et al.* (Belle), Phys. Rev. Lett. **91**, 262001 (2003).

LHCb Collaboration, Nature Phys. **18** (2022) no.7, 751-754;

LHCb Collaboration, Nat. Commun. **13** 3351 (2022).

exotic hadron

$\neq qqq$  or  $q\bar{q}$

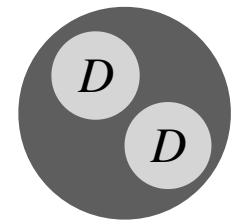


multiquarks

hadronic molecules



multiquarks

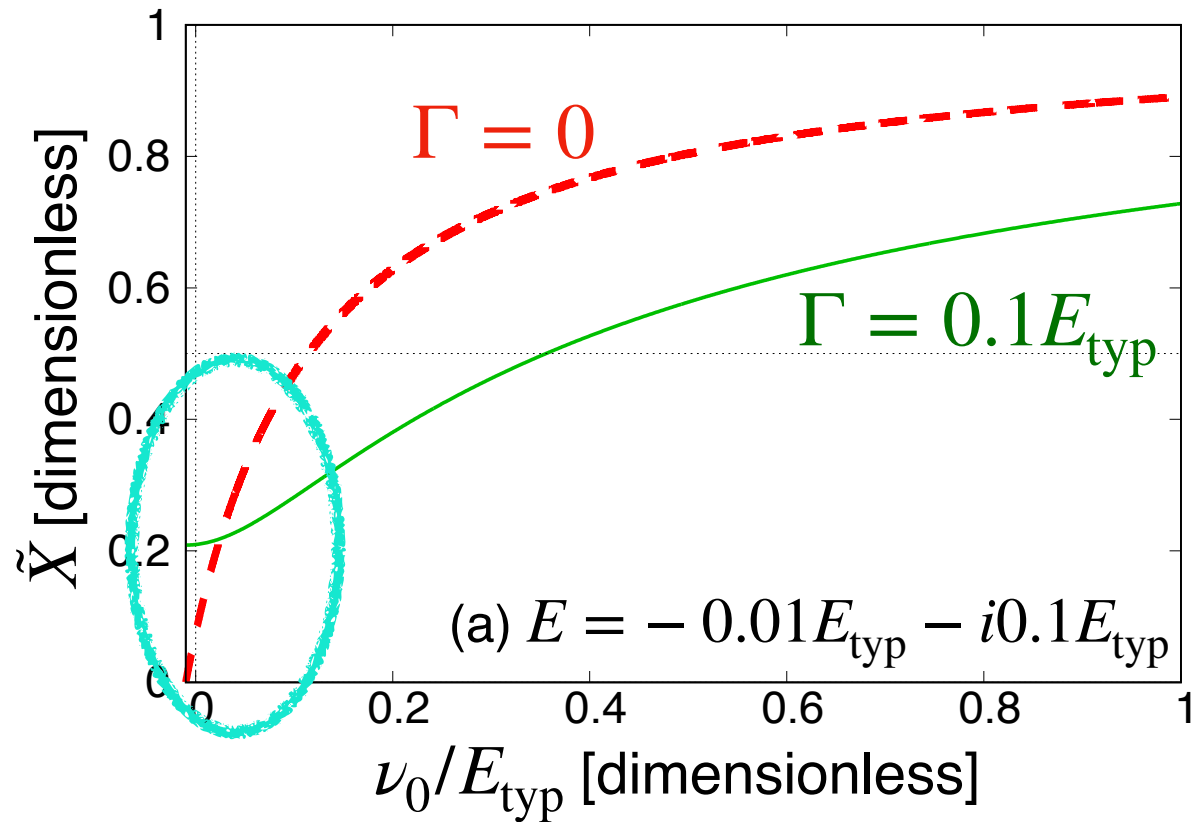


hadronic molecules

internal structure ?



EFT & compositeness



-  $X \neq 0$  with  $\Gamma \neq 0$   
 $\because g_0 \neq 0$  at  $\nu_0 = -B$   
 c.f.  $g_0 = 0$  at  $\nu_0 = -B$   
 with  $\Gamma = 0$

$$g_0^2 \left( -\nu_0 + i\frac{\Gamma}{2}; \nu_0, \Lambda \right) = \frac{\pi^2}{\mu} \left( -i\frac{\Gamma}{2} \right) \left[ \Lambda - \kappa \arctan \left( \frac{\Lambda}{\kappa} \right) \right]^{-1} \neq 0$$

$$X = \left[ 1 + \frac{\pi^2 \kappa}{g_0^2 \mu^2} \left( \arctan(\Lambda/\kappa) - \frac{\Lambda/\kappa}{1 + (\Lambda/\kappa)^2} \right)^{-1} \right]^{-1}$$

# Compositeness for two-channel case

23

$$V(k) = \begin{pmatrix} v(k) & v(k) \\ v(k) & v(k) \end{pmatrix}, \quad v(k) = \frac{g_0^2}{\frac{k^2}{2\mu_1} - \nu_0}.$$

$$G(k) = \begin{pmatrix} G_1(k) & 0 \\ 0 & G_2(k) \end{pmatrix}, \quad G_1(k) = -\frac{\mu_1}{\pi^2} \left[ \Lambda + ik \arctan \left( -\frac{\Lambda}{ik} \right) \right],$$

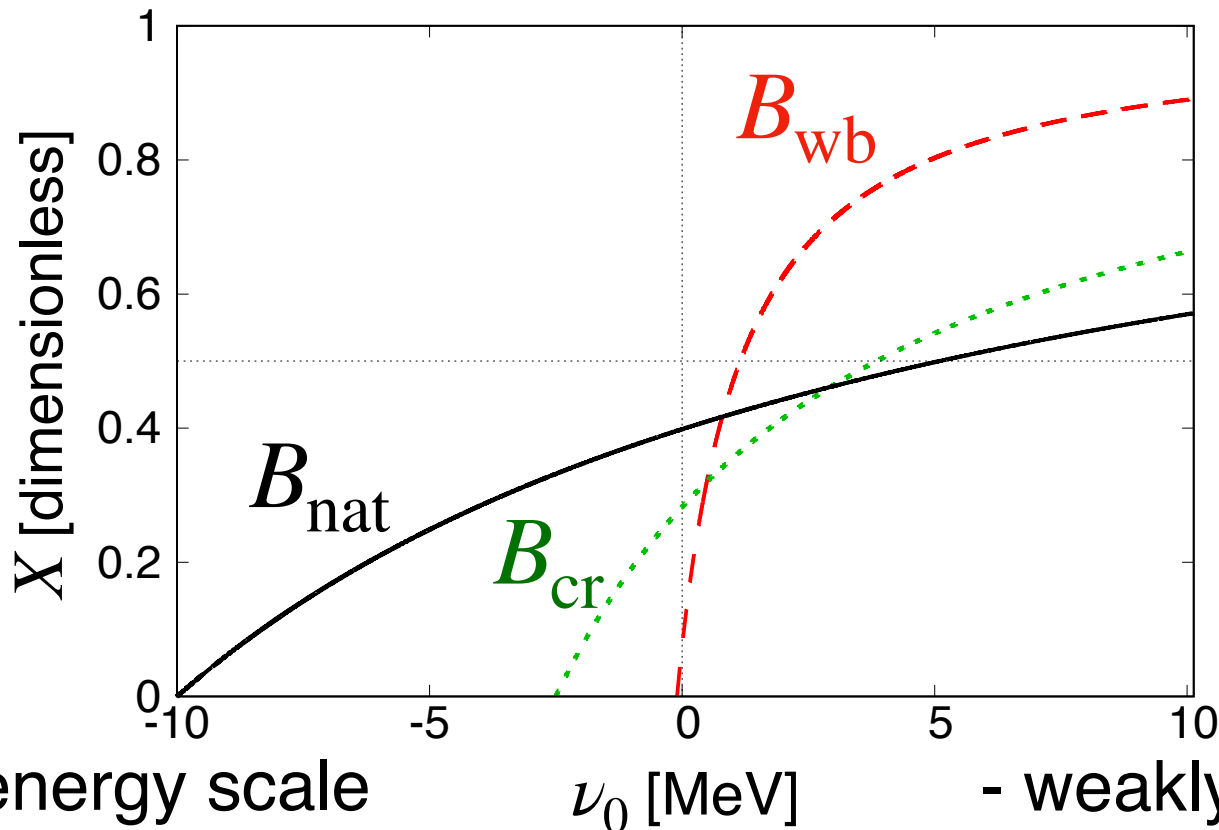
$$G_2(k') = -\frac{\mu_2}{\pi^2} \left[ \Lambda + ik' \arctan \left( -\frac{\Lambda}{ik'} \right) \right].$$

$$k = \sqrt{2\mu_1 E}, \quad k'(k) = \sqrt{2\mu_2(E - \Delta\omega)} = \sqrt{\frac{\mu_2}{\mu_1} k^2 - 2\mu_2 \Delta\omega}.$$

$$X_1 = \frac{G'_1}{(G'_1 + G'_2) - [v^{-1}]'},$$



$$X_2 = \frac{G'_2}{(G'_1 + G'_2) - [v^{-1}]'}.$$



- natural energy scale

- weakly-bound state

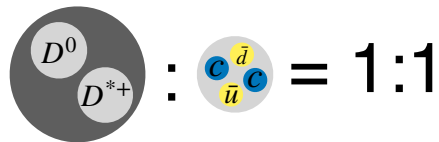
$$B_{\text{nat}} = \Lambda^2 / (2\mu) \sim 10 \text{ MeV}$$

$$B_{\text{cr}} \sim 2.5 \text{ MeV}$$

$$B_{\text{wb}} = 0.1 \text{ MeV}$$

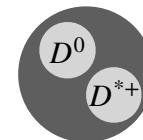
$X > 0.5$  for 25 % of  $\nu_0$   
= elementary dominant

$\therefore$  bare state origin



$X > 0.5$  for 88 % of  $\nu_0$   
= composite dominant

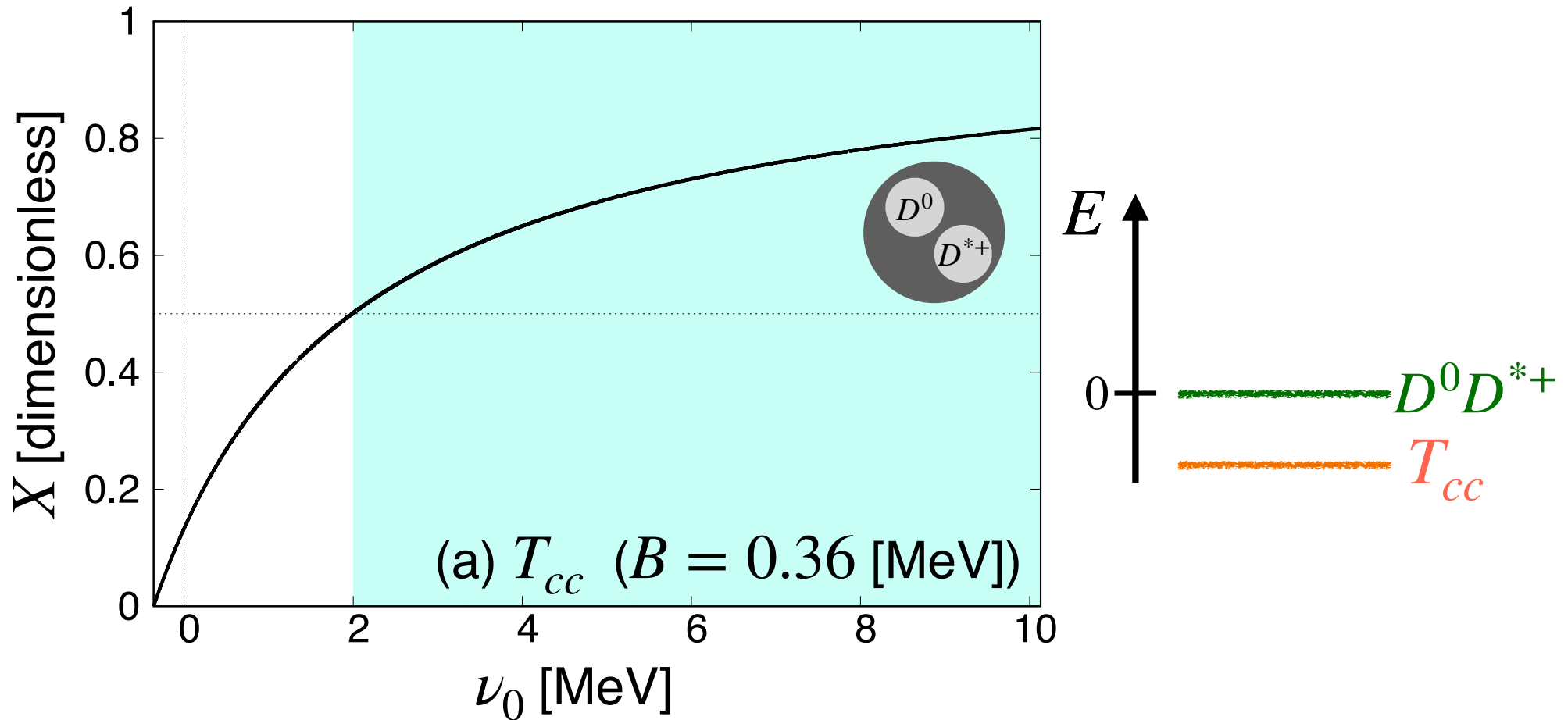
$\therefore$  low-energy universality !





# Application to $T_{cc}$

● single-channel



- $X > 0.5$  for 78 % of  $\nu_0$  = composite dominant
- fine tuning is necessary to realize  $X < 0.5$