



HPCI project field 5
“The origin of matter and the universe”

2014/06/05 ARIS2014
at Ito Hall, Tokyo

Frontier of Nuclear Shell- Model Calculations and High- Performance Computing



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Y. Tsunoda (Tokyo), and T. Yoshida (CNS)

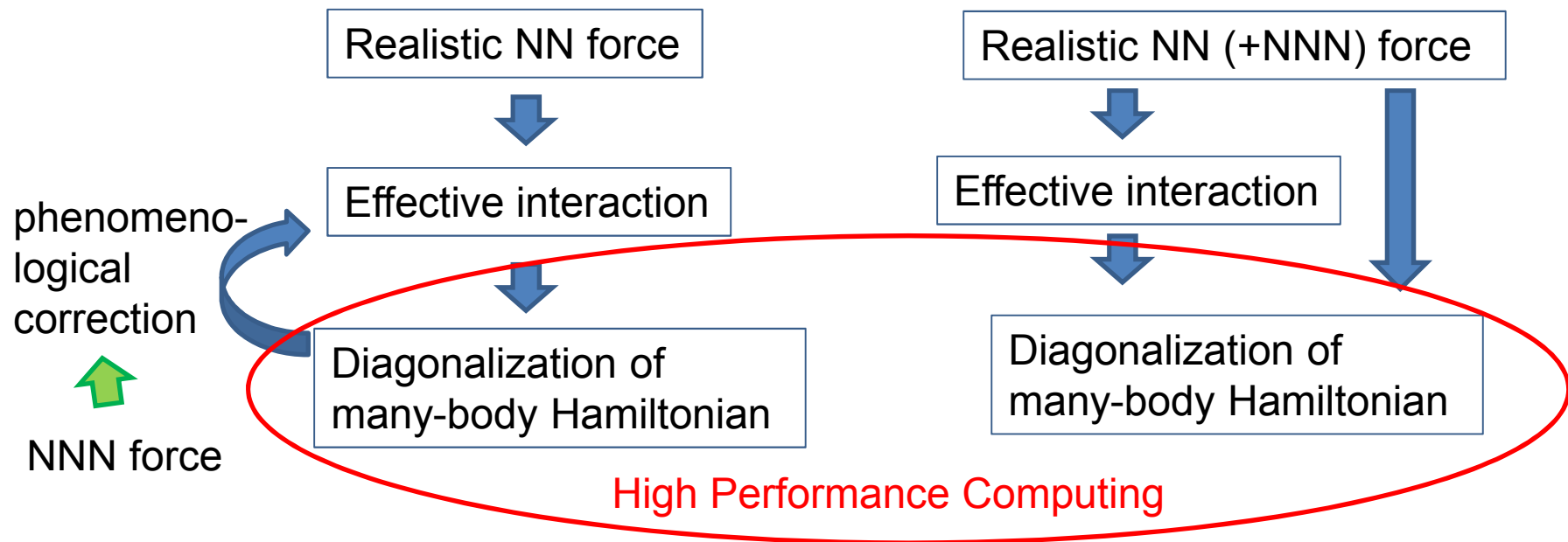
Recipe for shell-model calculations

large-scale shell-model calc.

model space: valence shell (+ extension)

no-core shell model

excitation energy ($N_{\max} \hbar \omega$)



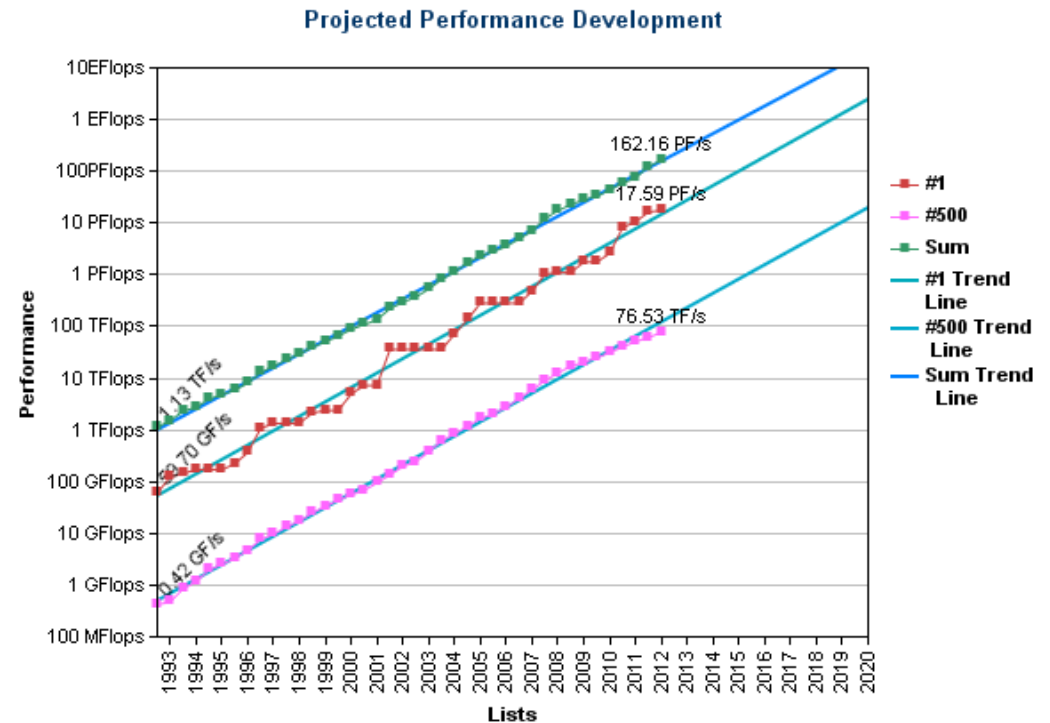
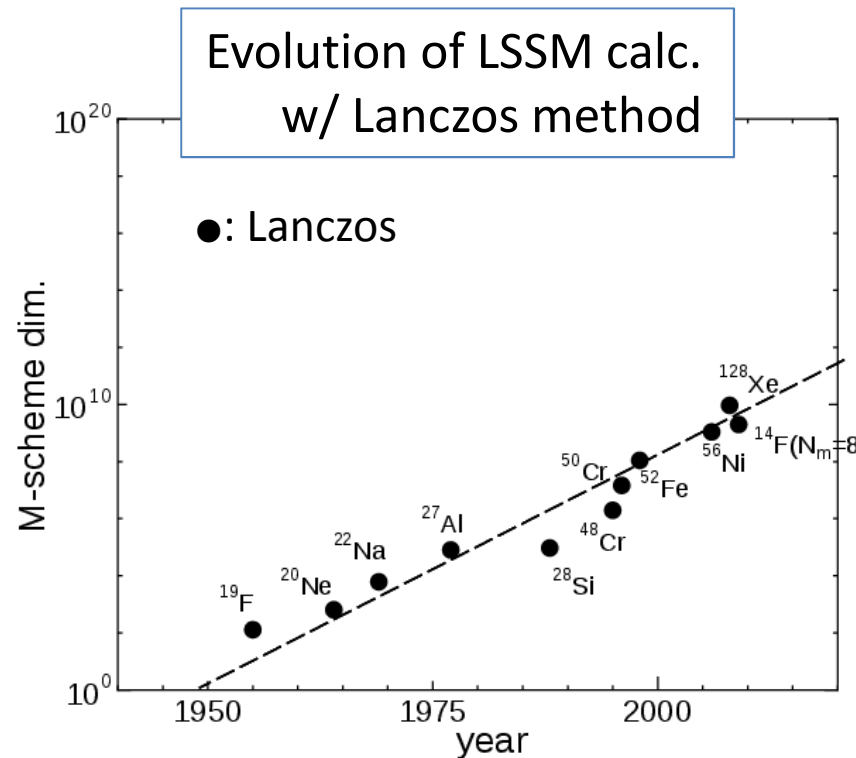
$$|\Psi\rangle = \sum_m v_m |m\rangle \quad \sum_{m'} \langle m|H|m'\rangle v_{m'} = E|m\rangle$$

$|m\rangle$: configuration, “spherical” Slater determinant of harmonic oscillator

The number of $|m\rangle \approx \frac{1}{f} \binom{N_p}{Z} \binom{N_p}{Z}$

Hamiltonian matrix is sparse
 $\Rightarrow$ Lanczos algorithm

Frontier of large-scale shell-model calculations



- The maximum dimension of the M -scheme Hamiltonian matrix of the LSSM evolve exponentially as the performance development of supercomputers
- Present limitation $\approx 10^{10-11}$ (as expected by the extrapolation in 2000)
e.g. ^{68}Ni in pf+g9/2+d5/2 5×10^{15}

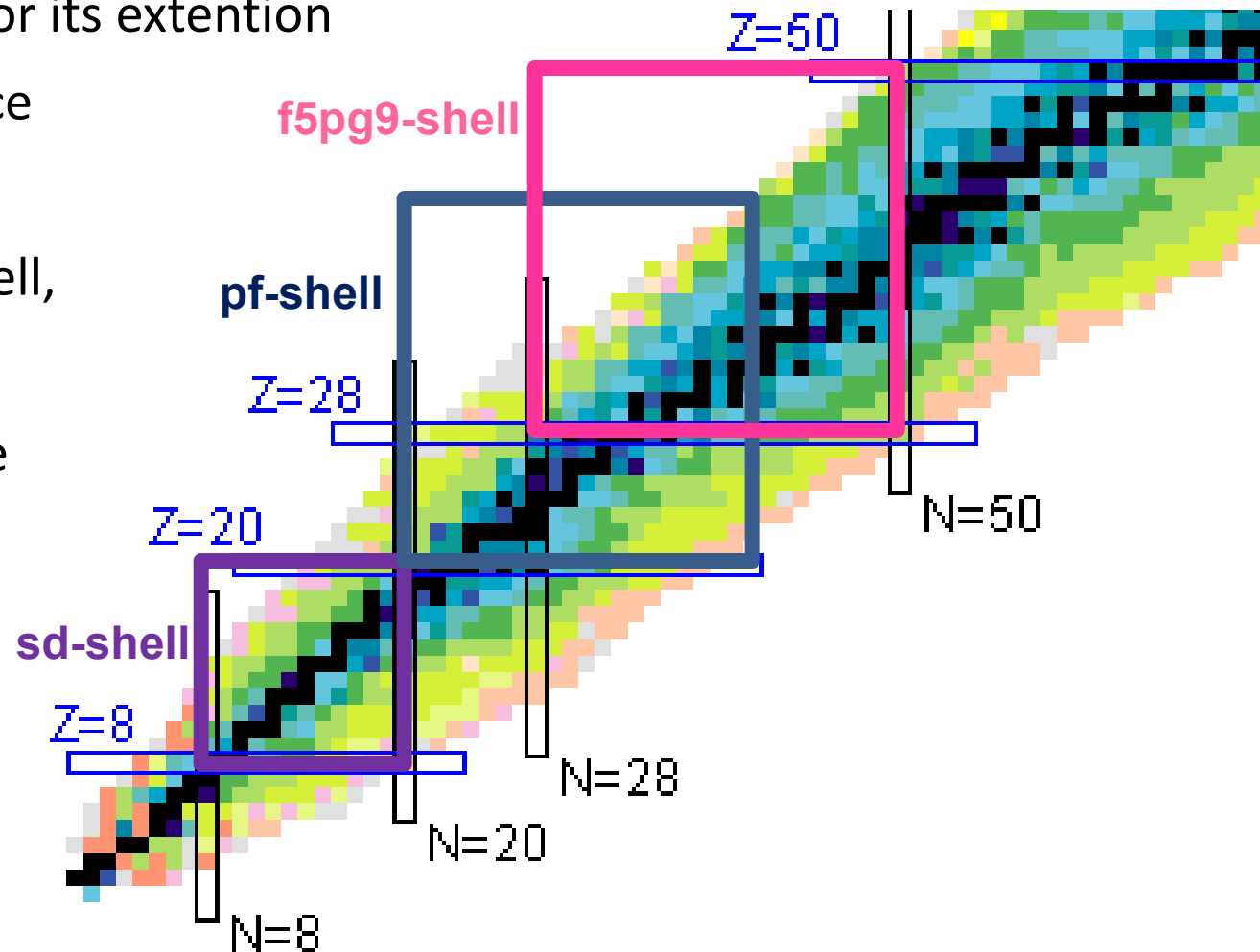


2040 for 10^{15} dim.?

Large-scale shell model calculations

valence shell or its extension
as model space

sd shell, *pf* shell,
f5pg9 shell
diagonalizable



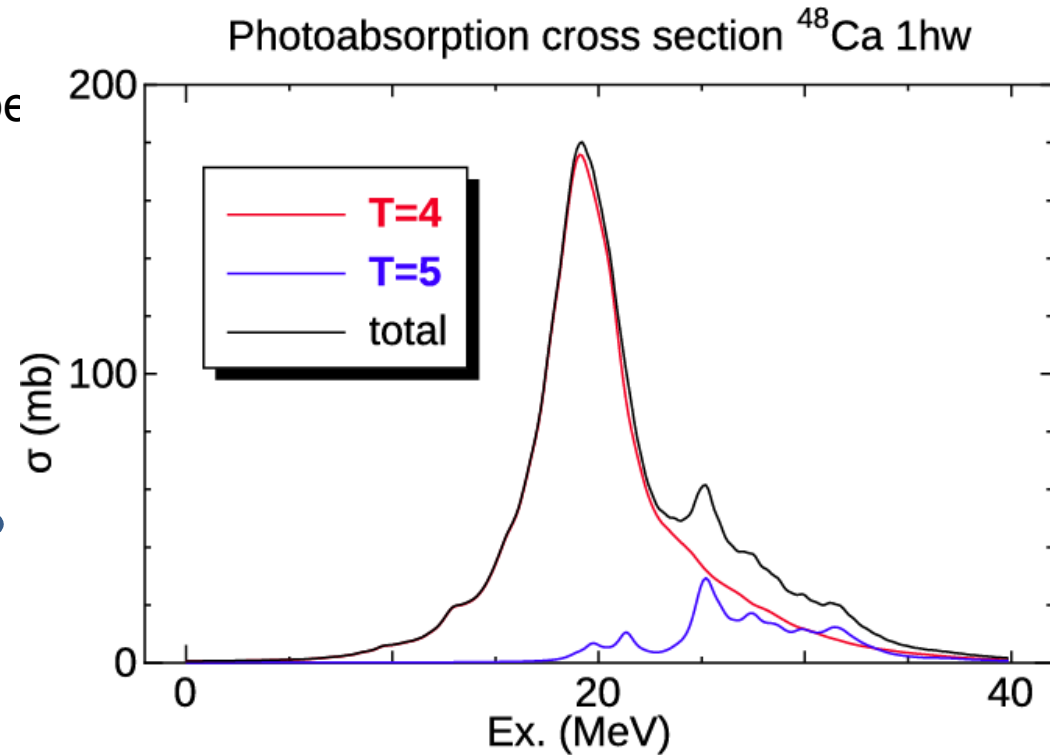
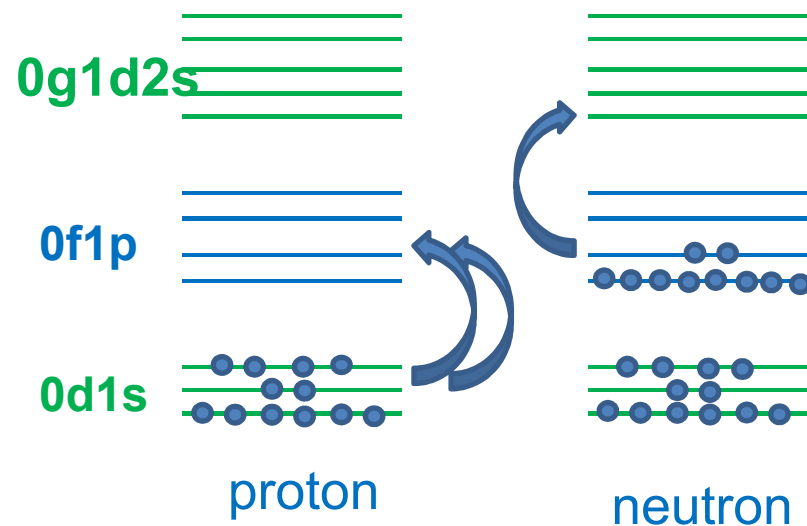
^{60}Zn in *pf*-shell : 2×10^9

^{78}Y in *f5pg9*-shell : 1.3×10^{10}

Example of the large-scale shell-model calculations

NS, Y. Utsuno, S. Ebata, T. Otsuka, M. Honma and T. Mizusaki, in preparation

1hw/3hw sd-pf-sdg shell calculations for negative-parity states of Ca isotope to describe E1 excitations



1hw : upto 1hw excitation in sd-pf-sdg shell

4.1×10^6 M-scheme dim.

at PC

We develop *M*-scheme shell-model code with Lanczos method for massive parallel computation “KSHELL”.

(1+3)hw: up to 3hw excitation in sd-pf-sdg shell

1.2×10^{10} M-scheme dim.

at supercomputer



HPCI Strategic Program for Innovative Research

by MEXT, Japan FY2011 – FY2015

to facilitate efficient use of K computer and HPC

- Strategic 5 Fields:

- Field1: “Supercomputational Life Science”



- Field2: “Computational Materials Science Initiative”



- Field 3: “Projection of Planet Earth Variations for Mitigating Natural Disasters”



- Field 4: “Industrial Innovation”



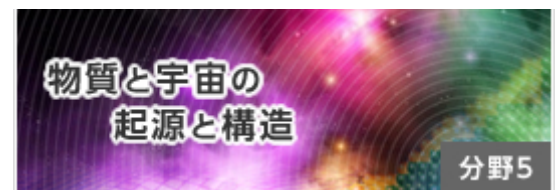
- Field 5: “The origin of matter and the universe”

- lattice QCD

- Nucleus ← Our group

- Supernova Explosion Head: Takaharu Otsuka (Tokyo)

- Early Star Formation

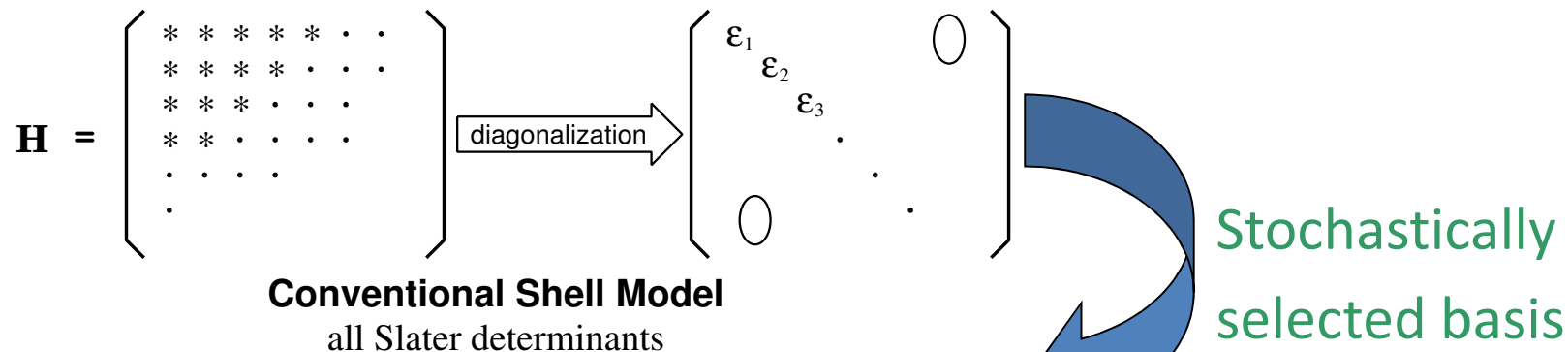


nuclear shell-model calculations using Monte Carlo shell model

Monte Carlo Shell Model

Otsuka, Honma, Mizusaki, Shimizu,
Utsuno, PPNP47, 319 (2001)

a tool to go beyond the conventional Lanczos diagonalization method
for Large-Scale Shell Model calculations



$$\mathbf{H} \approx \begin{pmatrix} * & * & * & \cdot \\ * & * & * & \cdot \\ * & * & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} \xrightarrow{\text{diagonalization}} \begin{pmatrix} \epsilon'_1 & & 0 \\ & \epsilon'_2 & \\ 0 & & \cdot \end{pmatrix}$$

Monte Carlo Shell Model

bases important for a specific eigenstate

$$|\phi(\sigma)\rangle = \prod e^{\Delta\beta \cdot h(\sigma)} \cdot |\phi^{(0)}\rangle$$

$\hat{h}(\sigma)$ one-body Hamiltonian

σ ... auxiliary field

random numbers

$$|\Psi\rangle = \sum_{k=1}^{N_{MCSM}} f_k P^{J,\pi} |\phi_k\rangle$$

↑ MCSM basis, deformed Slater det.

$$|\phi_k\rangle = \prod_{\alpha=1}^N \left(\sum_{i=1}^{N_{sp}} c_i^{\dagger} D_{i\alpha}^{(k)} \right) | - \rangle$$

Wave function is described as a linear combination of J -projected deformed Slater determinants.
Increase MCSM dimension, or number of deformed Slater det. till energy converges.

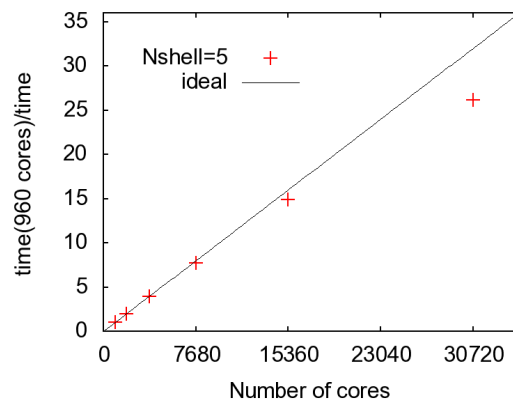
“Advanced” Monte Carlo Shell Model towards K computer

T. Otsuka, M. Honma, T. Mizusaki, NS, and Y. Utsuno, Prog. Part. Nucl. Phys. **47**, 319 (2001).

NS, Y. Utsuno, T. Mizusaki, T. Otsuka, T. Abe, and M. Honma, Phys. Rev. C **82**, 061305(R) (2010).

PC cluster
100CPU

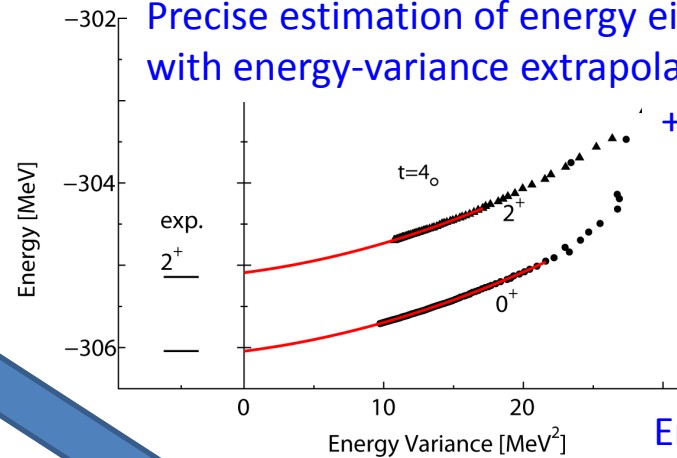
Parallel efficiency



OpenMP+MPI hybrid parallel
100,000 cores

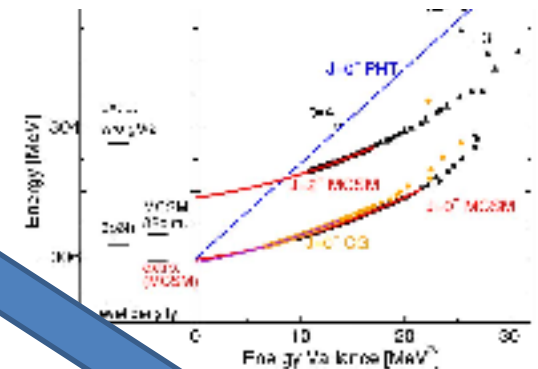
NS, T. Abe, Y. Tsunoda, Y. Utsuno, T. Yoshida, T. Mizusaki, M. Honma and T. Otsuka, Prog. Theor. Exp. Phys. **2012** 01A205 (2012).

Precise estimation of energy eigenvalues
with energy-variance extrapolation
+ basis reordering technique

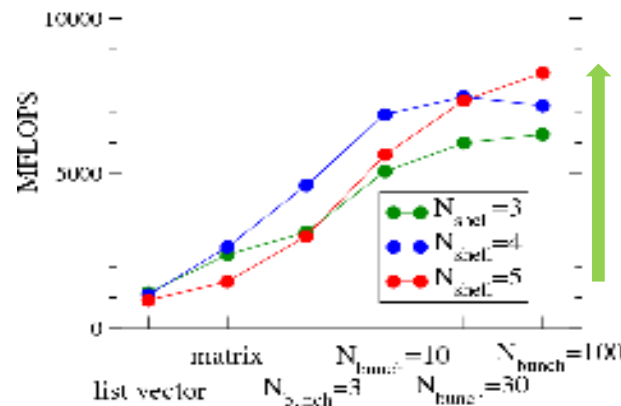


NS et al., Phys. Rev. C **85** 054301(2012).

Energy gradient for
optimization of wave function



Computation algorithm



8 times
acceleration

Y. Utsuno, NS, T. Otsuka, and T. Abe,
Comp. Phys. Comm. **184** 102 (2013).

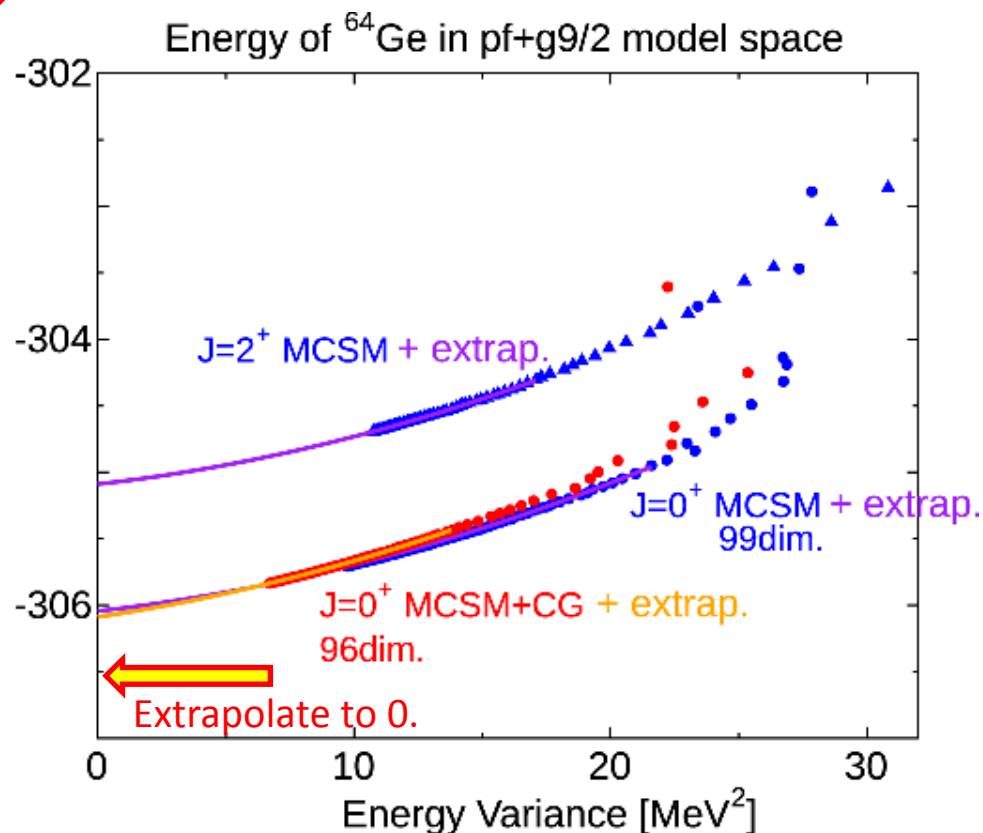
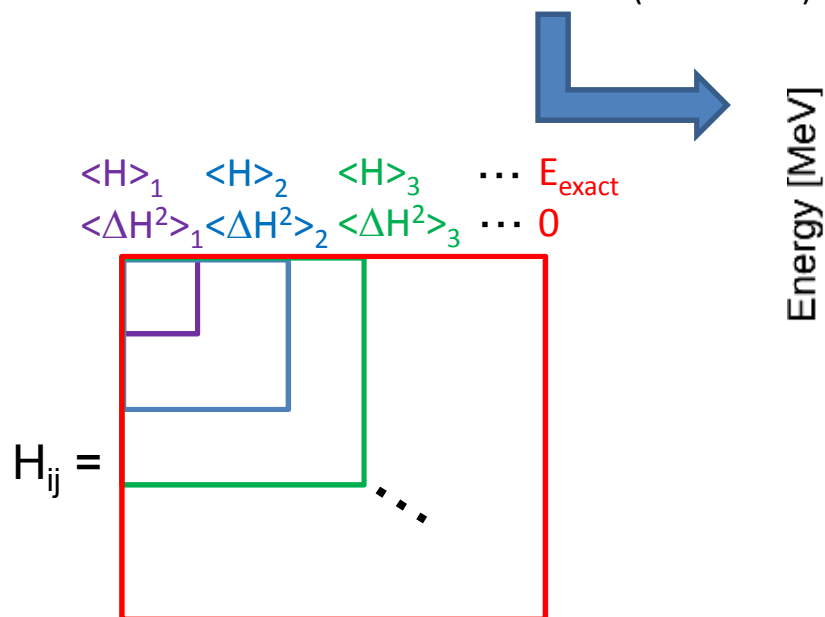
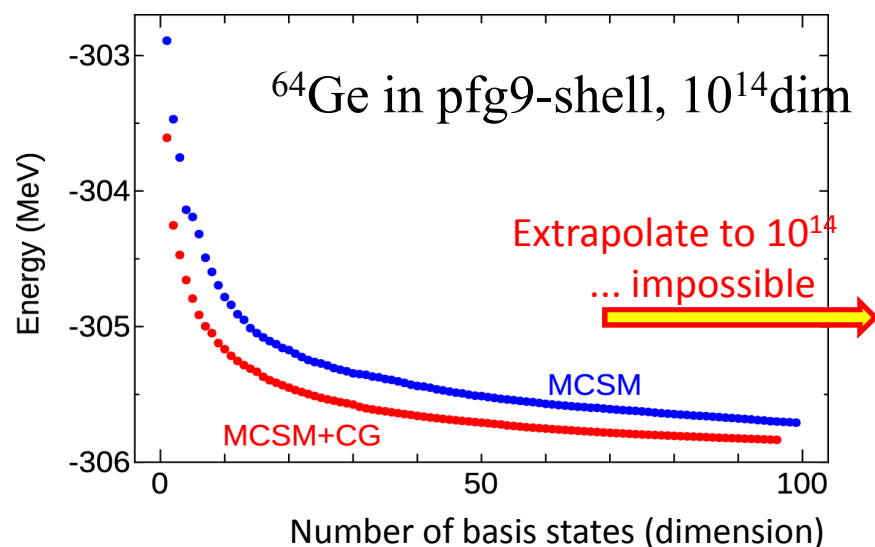


Energy variance extrapolation in the MCSM: ^{64}Ge in $pfg9$ -shell

Energy variance: $\langle \Delta H^2 \rangle = \langle H^2 \rangle - \langle H \rangle^2$

As the number of basis states increases, the approximated w.f. approaches the exact one and the energy variance approaches zero.

Extrapolate towards $\langle \Delta H^2 \rangle \rightarrow 0$



Advantages of the MCSM

- computation intensive, rather small memory access, small I/O ... good at massive parallel computation

M-scheme Lanczos method

$$|\Psi\rangle = \sum_{m=1}^{\sim 10^{10}} f_k |\phi_m\rangle$$

huge memory and disk I/O

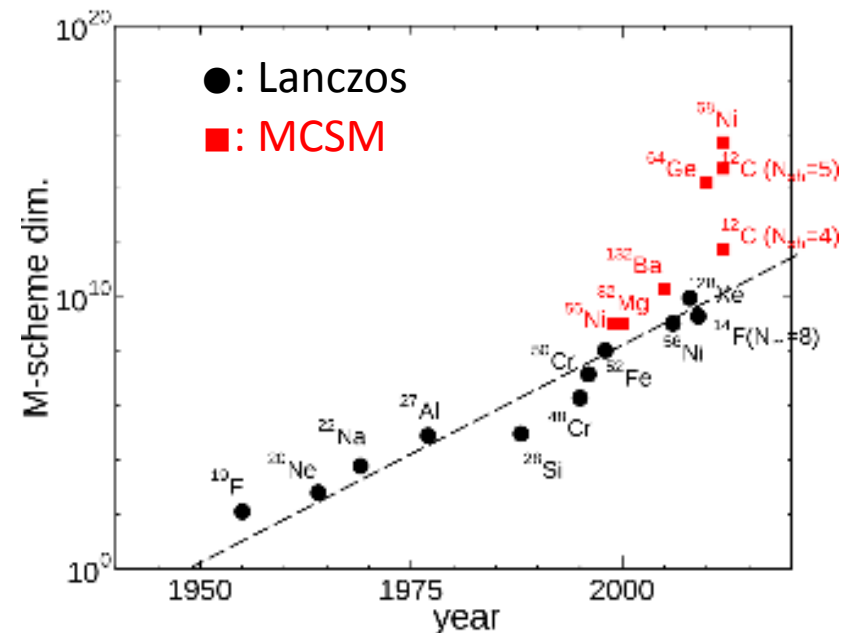
vs

MCSM wave function

$$|\Psi\rangle = \sum_{k=1}^{N_{MCSM} \sim 100} f_k P^{J,\pi} |\phi_k\rangle$$

discretized integral
of Euler's angle

- overcome the limitation of Lanczos method
- analyze “intrinsic” wave function and shape (discussed later)

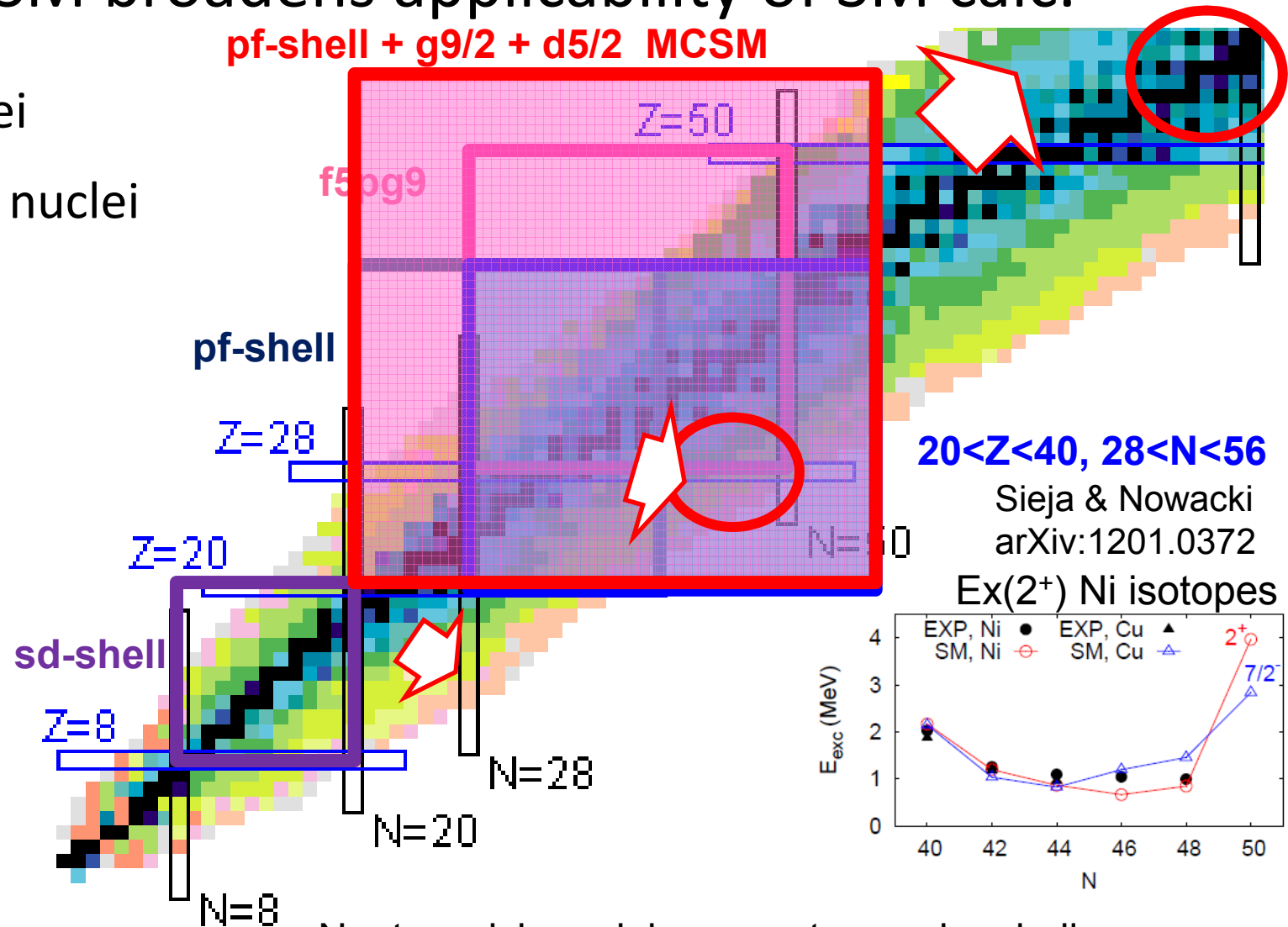


MCSM broadens applicability of SM calc.

pf-shell + g9/2 + d5/2 MCSM

heavier nuclei

neutron-rich nuclei



Neutron-rich nuclei: across two major shells

⇒ proton neutron different model space

+ truncation

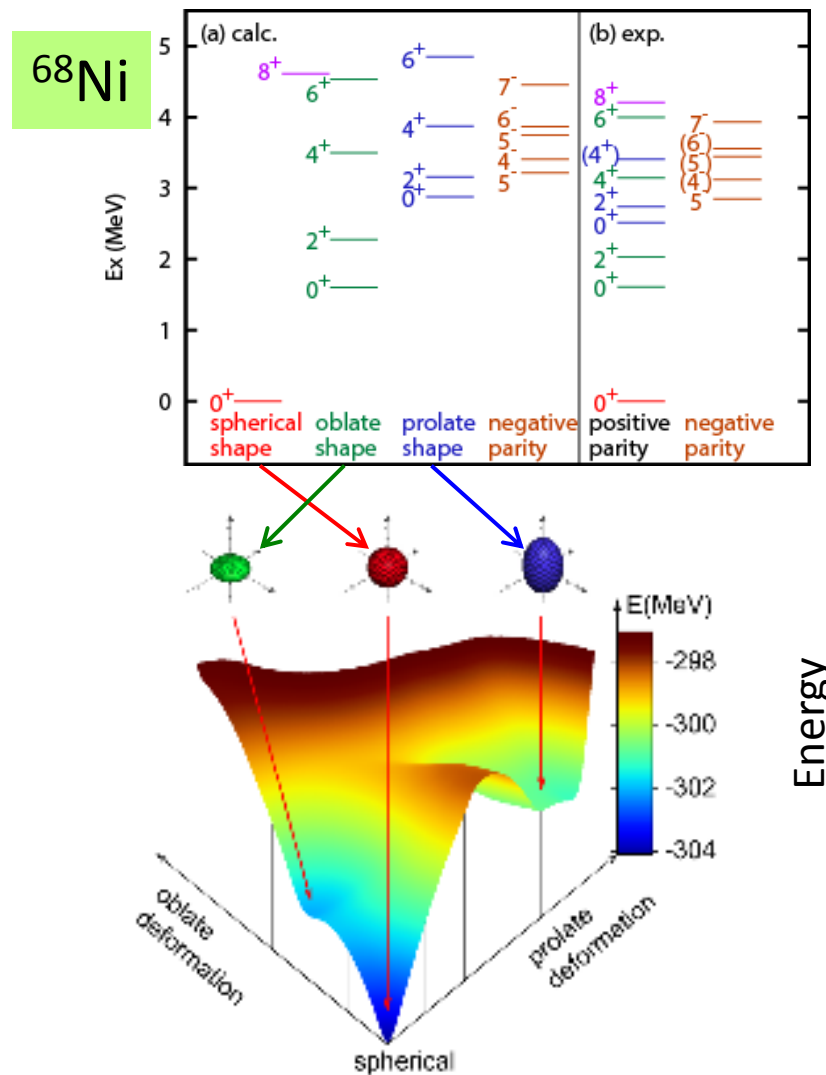
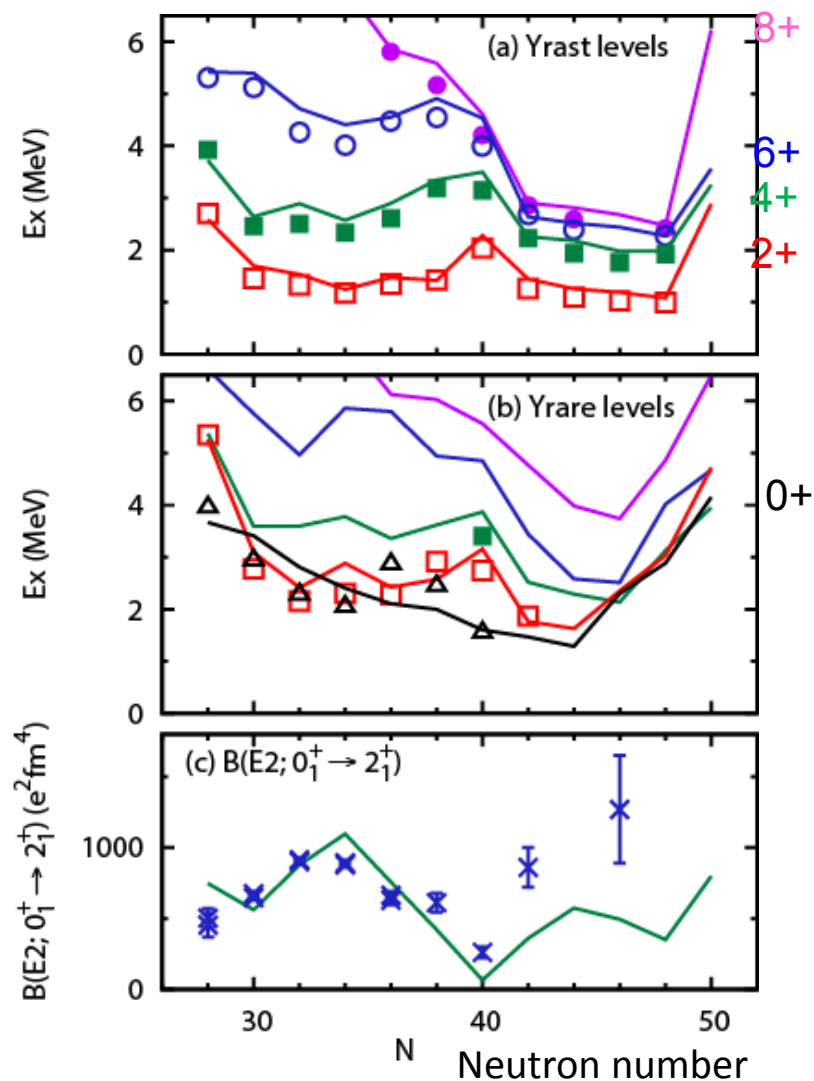
⇒ full isospin-invariant space + MCSM

Neutron-rich Ni isotopes and shell evolution

Exp.: S. Suchyta PS2-A039

pf-shell + g_{9/2}, d_{5/2} orbits w/o truncation
 ... 5.1×10^{15} M-scheme dim.

Y. Tsunoda, T. Otsuka, NS, M. Honma and Y. Utsuno,
 PRC 89, 031301(R) (2014)



^{68}Ni 0^+ states $\Leftrightarrow$ different shapes

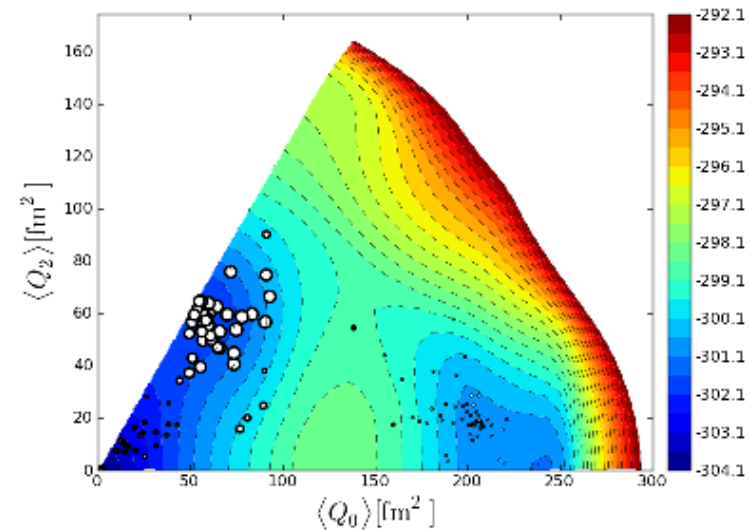
MCSM w.f. $|\Psi\rangle = \sum_{k=1}^{N_{MCSM} \sim 100} f_k P^{J,\pi} |\phi_k\rangle$

PES is calculated by Q-constraint Hartree-Fock calc.

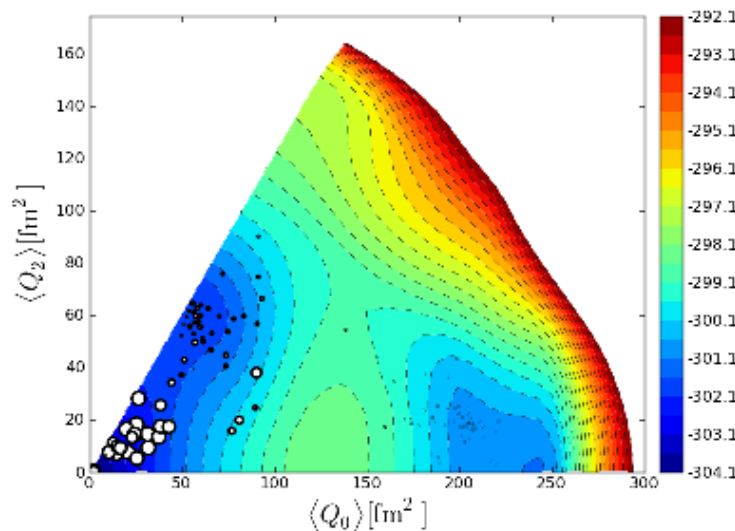
Location of circle : quadrupole deformation of $|\phi_k\rangle$

Area of circle : overlap probability $\frac{1}{N_k} |\langle \Psi | P^{J,\pi} | \phi_k \rangle|^2$

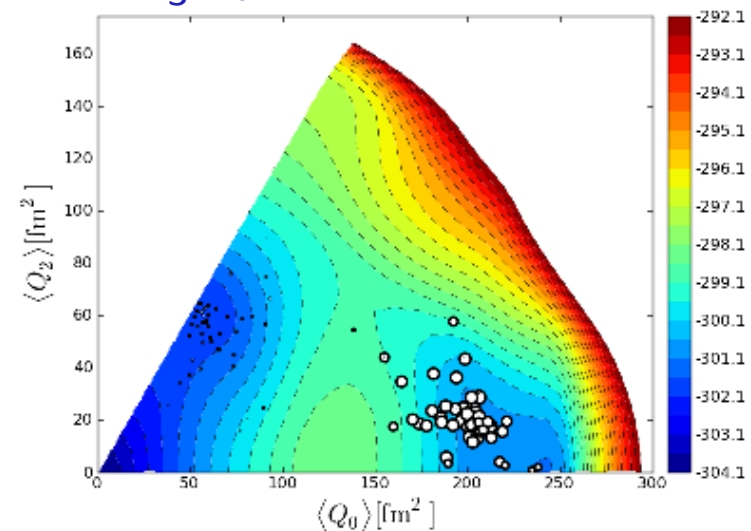
0^+_2 : oblate $\beta \sim -0.2$



0^+_1 : spherical $\beta \sim 0$



0^+_3 : prolate $\beta \sim +0.4$



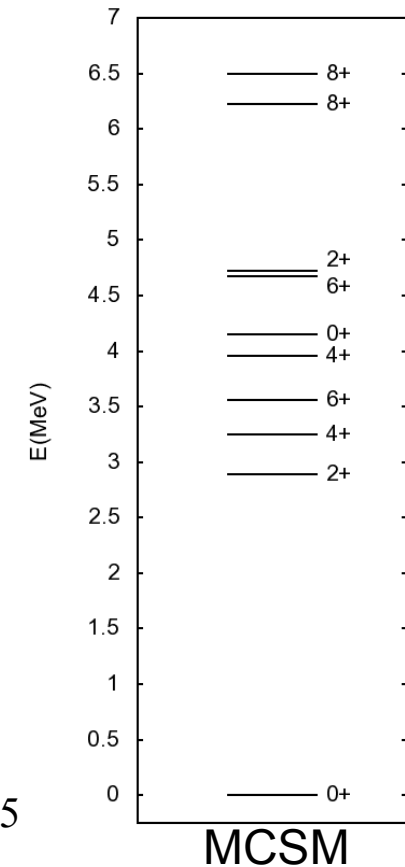
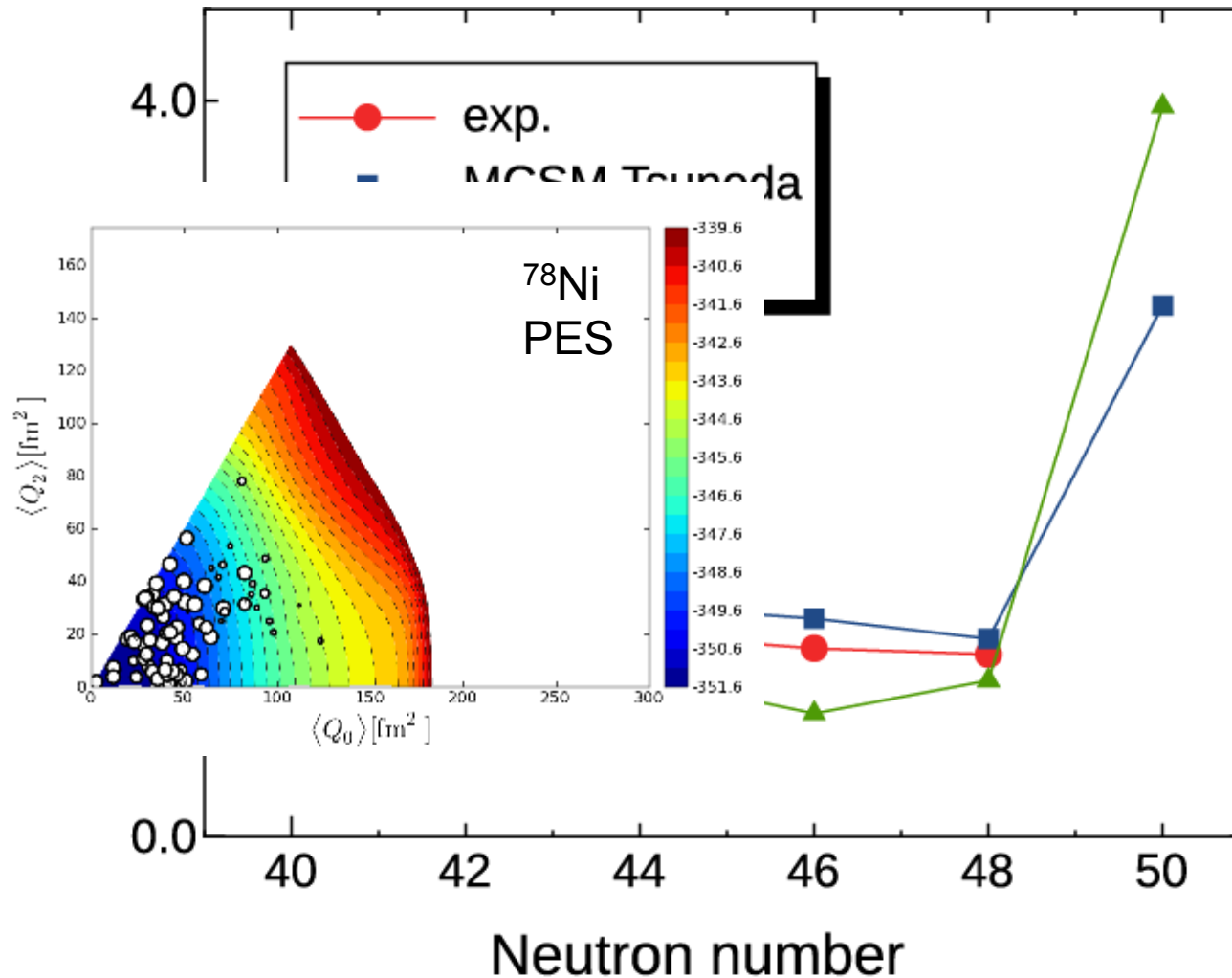
Towards ^{78}Ni : prediction

Excitation energies of Ni isotopes

$Z=28$, $N=50$

How magic?

Both SM results give large 2^+ energy



SM: Sieja and Nowacki, arXiv:1201.0373

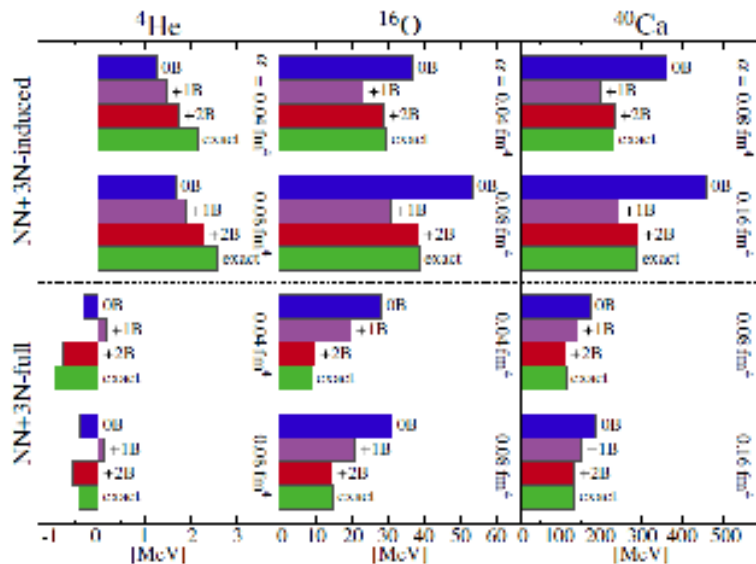
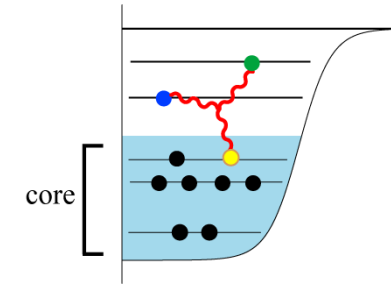
MCSM: Y. Tsunoda *et al.*, PRC 89 031301(R) (2014)

π pf, ν f5pg9d5

full pfp9d5

Towards SM calc. w/o phenomenological correction

3-body force into normal ordered 2-body force

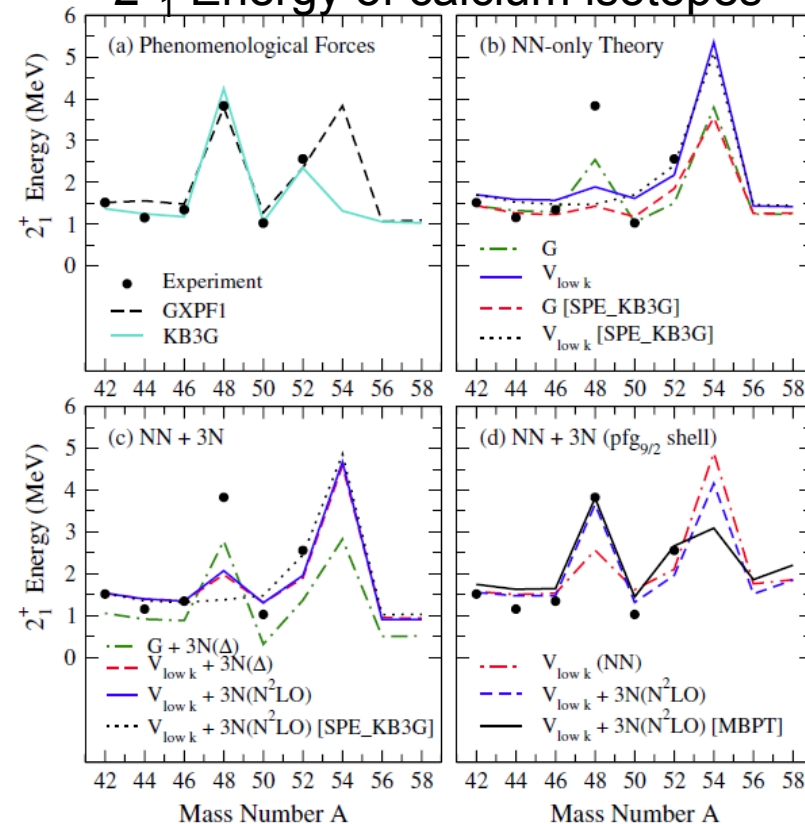


Normal ordered 2-body force is good approximation for no-core SM

Ref. R. Roth et al., PRL 109, 052501

Phenomenological fit of SM interaction
← 3-body force origin

2^+_1 Energy of calcium isotopes



SM calc. w/o phenom. correction

Normal ordered 2-body force

Ref. J. Holt, T. Otsuka, A. Schwenk and T. Suzuki, JPG 39 085111 (2012)

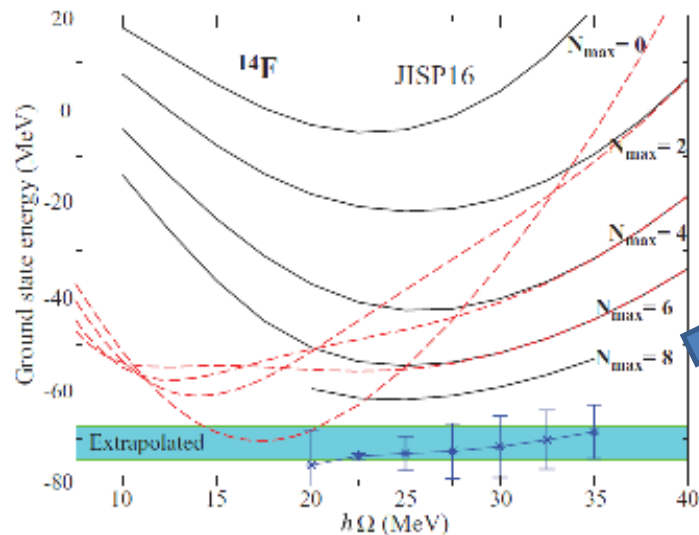
No-core shell model

Importance Truncated
NCSM

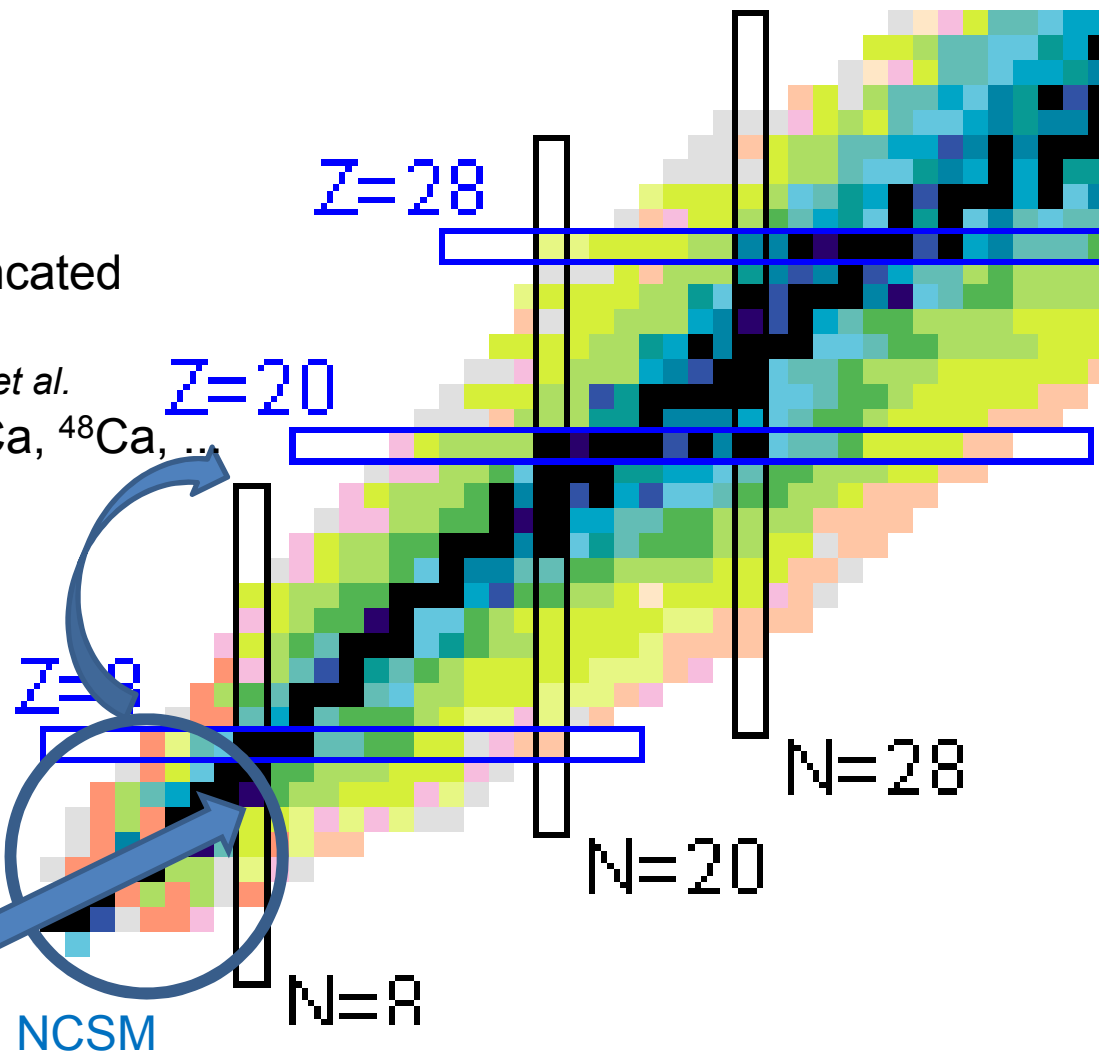
R. Roth, P. Navratil, *et al.*

^{40}Ca , ^{48}Ca , ...

Excitation energy ($N_{\max} h\omega$) truncation
Extrapolation to infinite model space



P. Maris, J. P. Vary and A. Shirokov
PRC 79, 014308 (2009)

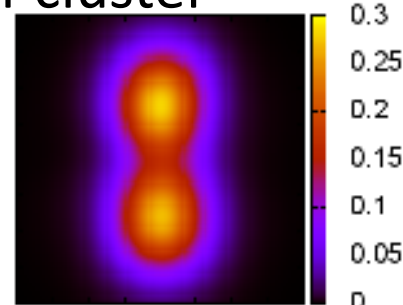


NCSM and reaction ... S. Quaglioni *et al.*

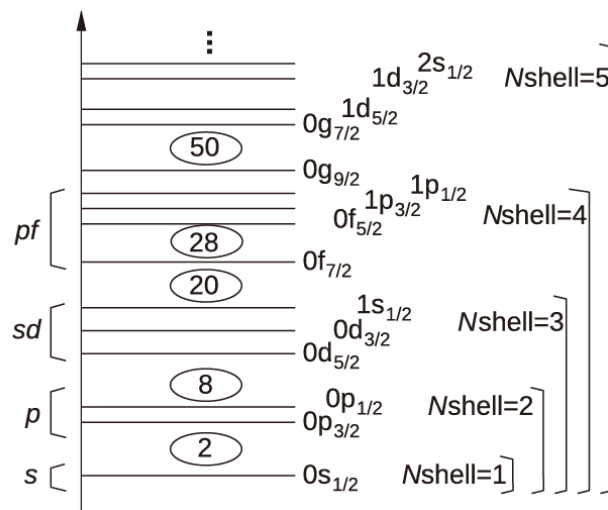
MCSM for no-core shell model calc.

- Enhance the applicability to heavier region by MCSM

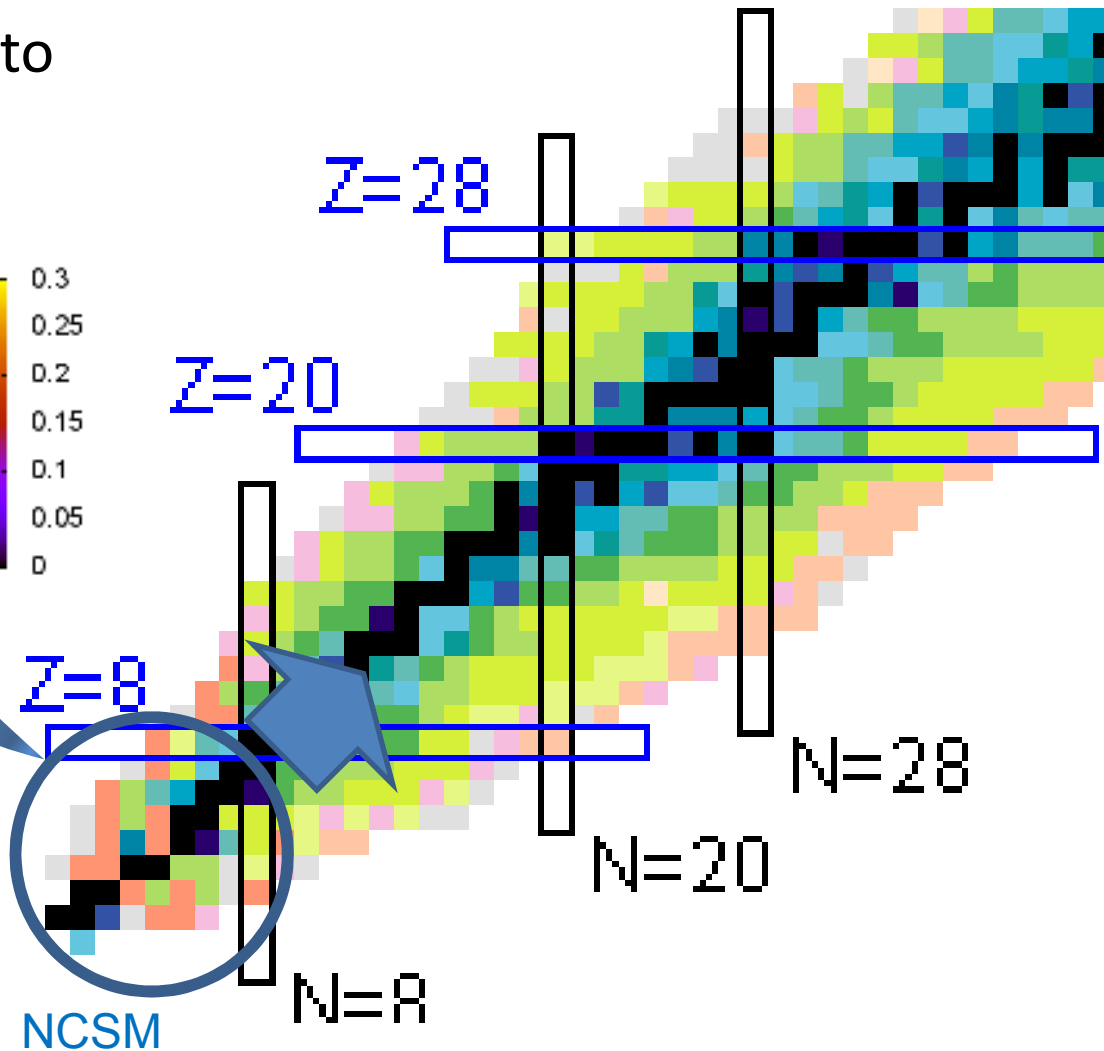
- Description of cluster structure by MCSM



- N_{shell} truncation



truncation by single-particle energy



Extrapolation to infinite model space based on EFT

$$\lambda_{sc} = \sqrt{m\hbar\omega/(N + 3/2)} = \lambda^2/\Lambda$$

$$\Lambda = \sqrt{m(N + 3/2)\hbar\omega}$$

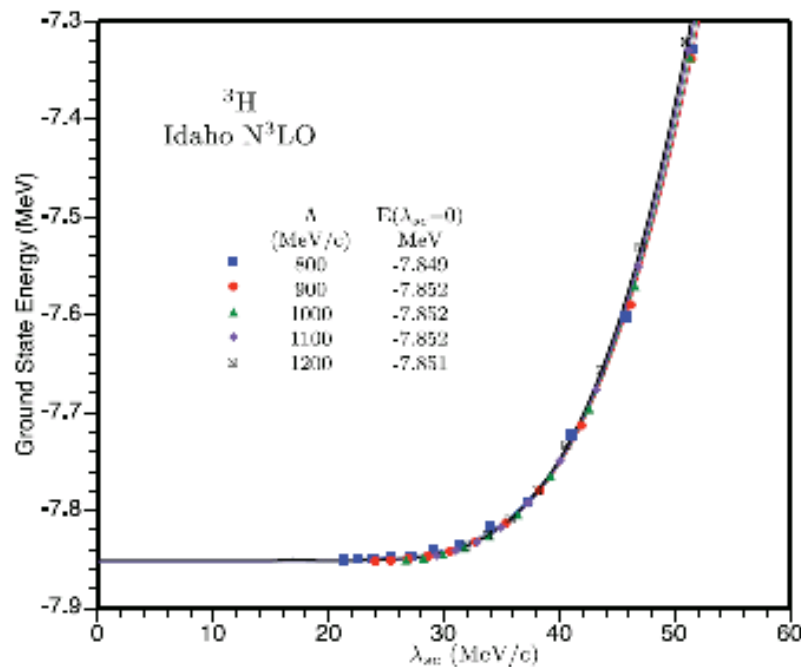
MCSM

Excitation energy (N_{nshell}) truncation

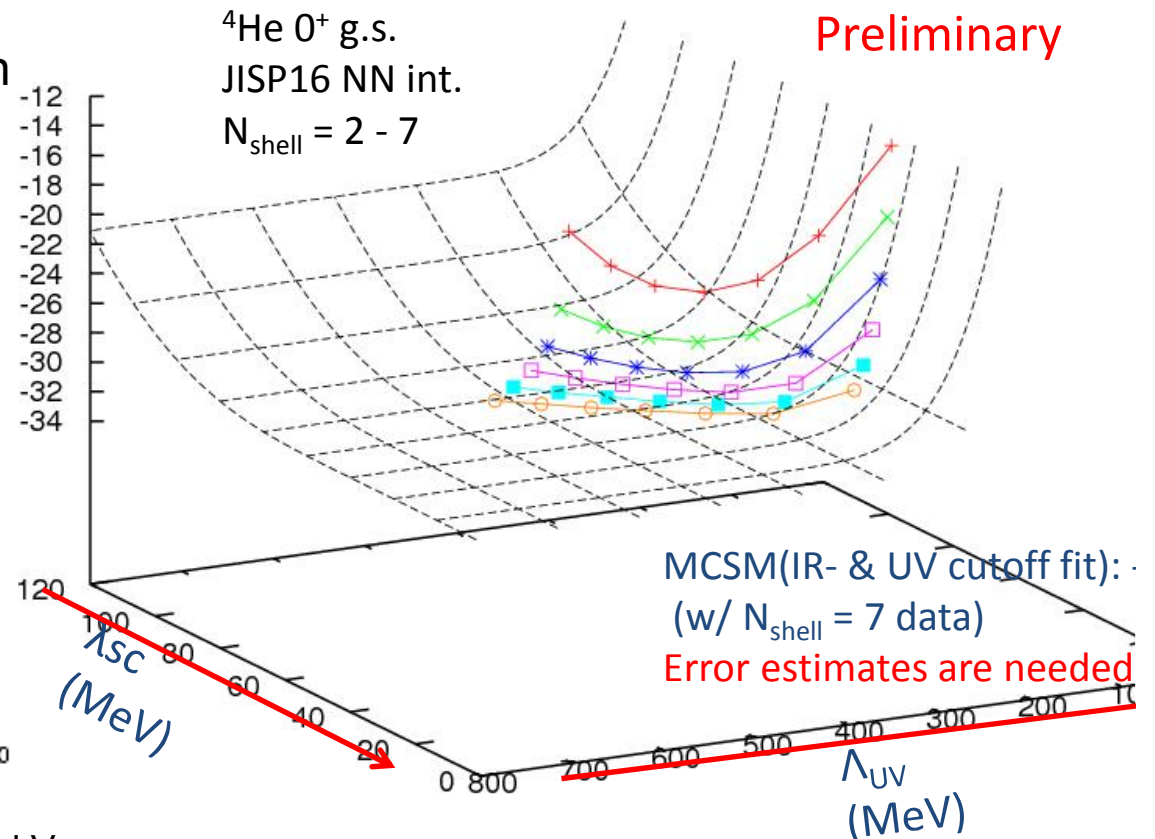
$$E(\lambda, \Lambda) = E(\lambda = 0, \Lambda = \infty) + a \exp(-b/\lambda) + c \exp(-\Lambda^2/d^2)$$

NCSM

Excitation energy ($N_{max} \hbar\omega$) truncation



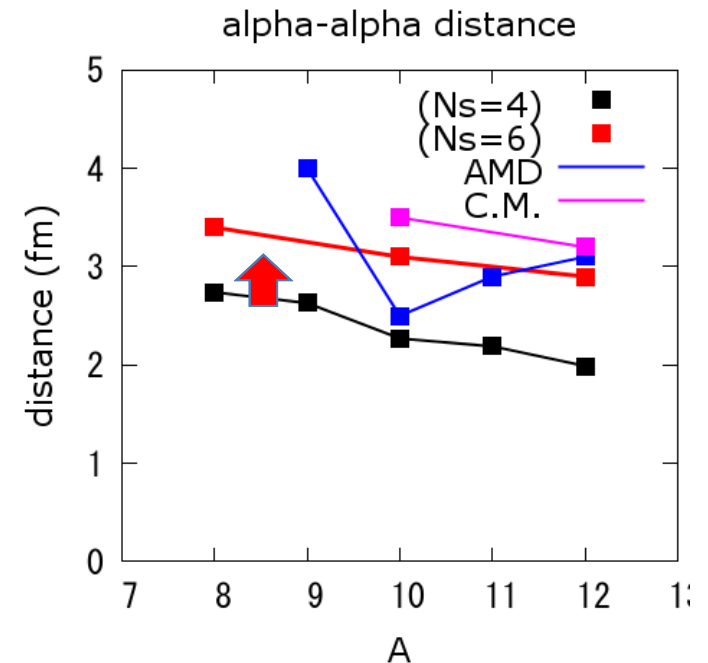
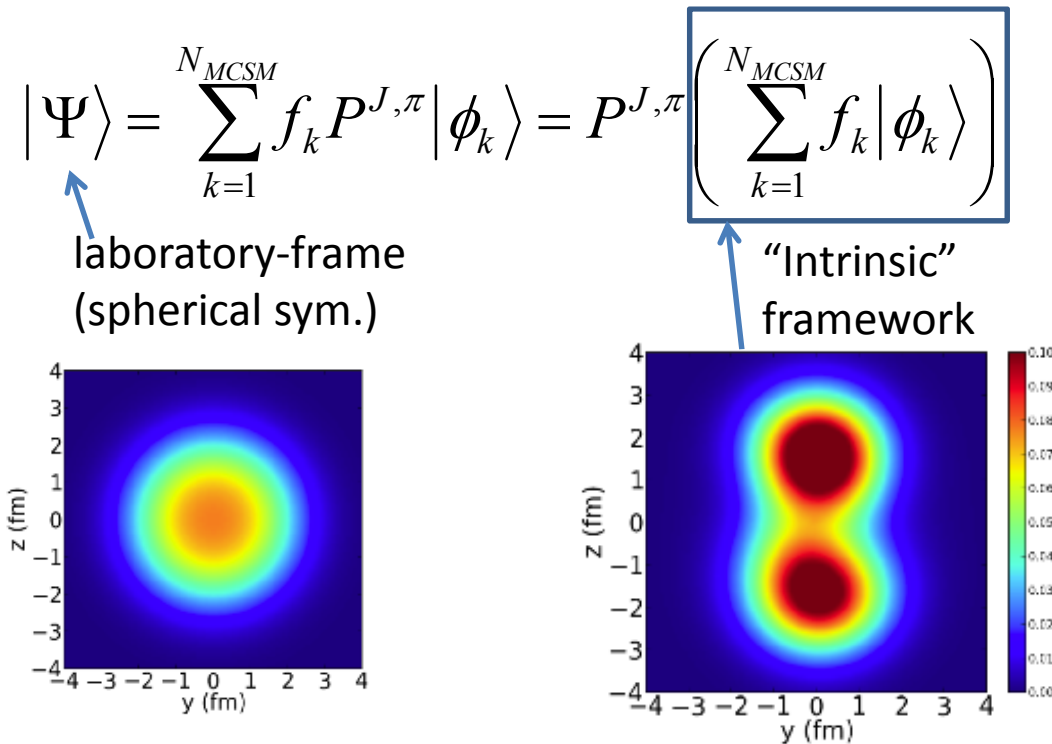
Ref. Coon, Avetian, Kruse, van Klock, Maris and Vary
PRC 86 054002 (2012)



On going: $^8\text{Be}(0^+)$, $^{12}\text{C}(0^+)$, $^{16}\text{O}(0^+)$, ...

Density profile of the MCSM wave function

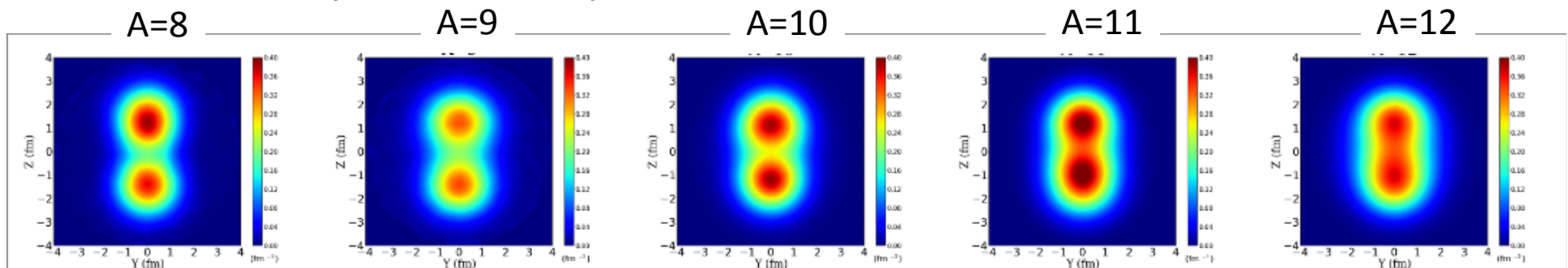
Density distribution of the ground state $J^\pi = 0^+$ of ^8Be



[AMD: Y.Enyo et al. PRC 2003]

[C.M. : M.Ito, N.Itagaki et al, PRL 2008]

Matter density of Be isotopes JISP16 interaction



Summary

- HPCI project promotes nuclear physics. In HPCI Field 5 Subject 2, we developed the “advanced” MCSM to run on K computer
 - energy-variance extrapolation, code R&D
 - unified description of neutron-rich Ni isotopes, shell evolution and shape coexistence
 - test the feasibility of no-core MCSM and propose a new method to draw the density profile
- MCSM purvey analysis of “intrinsic” shapes and their mixing, which enlighten cluster structure and shape coexistence.

Collaborators

HPCI project, Field 5

- | | |
|------------------------------------|--------------------------|
| • Takaharu Otsuka (Tokyo) | Co-chair |
| • Takashi Abe (Tokyo) | Poster AS2-A018 |
| • Shuichiro Ebata (Hokkaido) | Poster PS2-B022 |
| • Michio Honma (Aizu) | |
| • Yoritaka Iwata (CNS, Tokyo) | Poster PS2-A061 |
| • Takahiro Mizusaki (Senshu Univ.) | |
| • Tomoaki Togashi (CNS, Tokyo) | Poster PS2-A035 |
| • Naofumi Tsunoda (CNS, Tokyo) | |
| • Yusuke Tsunoda (Tokyo) | Poster PS2-A037 |
| • Yutaka Utsuno (JAEA) | Oral Jun/3 |
| • Tooru Yoshida (CNS, Tokyo) | Poster PS2-A005 |
| | |
| • James Vary (Iowa) | |
| • Pieter Maris (Iowa) | ICNT workshop, next week |

Thank you !