

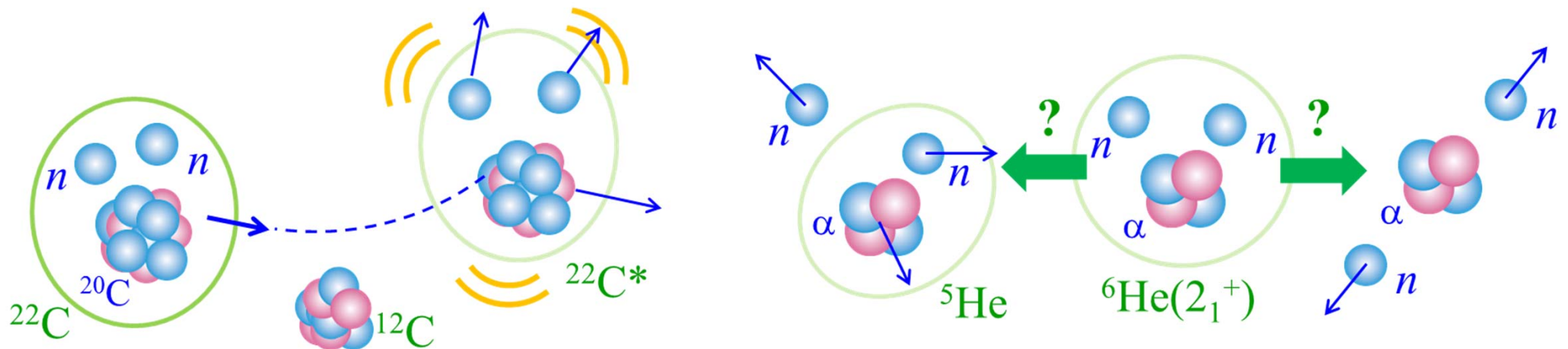
# Dynamical Studies of the Formation and Decay of Particle-Unbound States

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<sup>1</sup>RCNP, Osaka University, <sup>2</sup>RIKEN Nishina Center, <sup>3</sup>Osaka Institute of Technology,

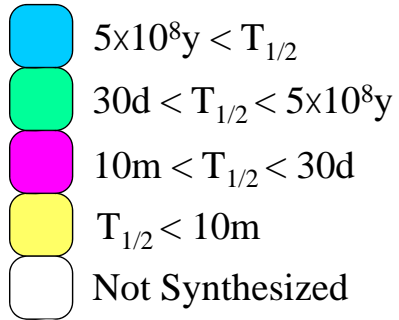
<sup>4</sup>Ichinoseki National College of Technology, <sup>5</sup>Kyushu University



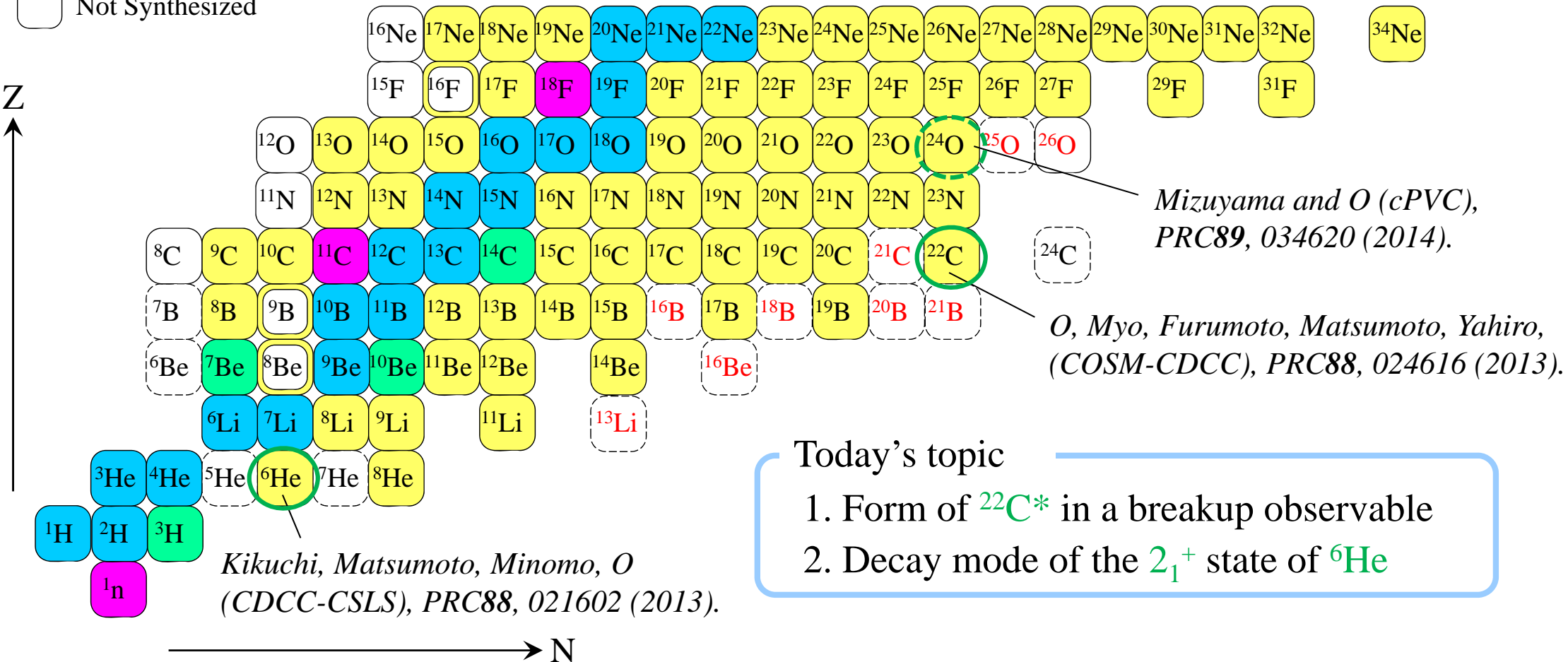
# Exploration of unbound (but not free) systems

Our Aim

*Dynamical description of Formation and Decay of unbound systems*



A: discussed in this conference (cf. talks by Gade, Hagino, Marques, Kondo, Obertelli...)



Today's topic

1. Form of  $^{22}\text{C}^*$  in a breakup observable
2. Decay mode of the  $2_1^+$  state of  $^6\text{He}$

# COSM-CDCC for $^{22}\text{C}$ breakup by $^{12}\text{C}$

## Structural part: Cluster Orbital Shell Model (COSM)

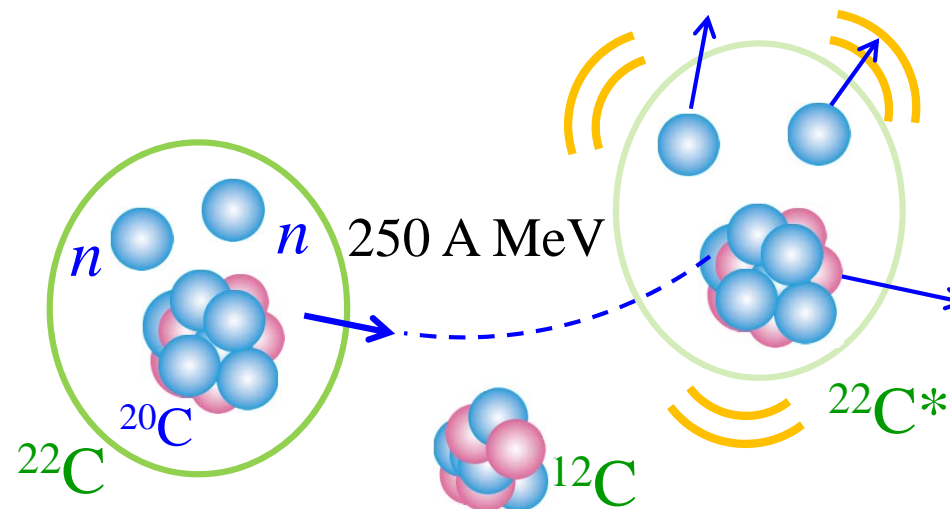
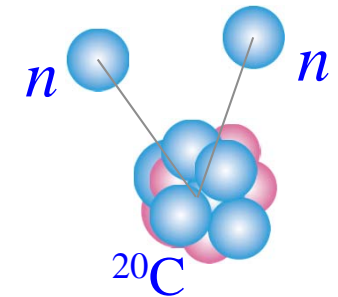
- ✓ Core + valence  $N$  system is described well.
- ✓ Pseudo states covering large space are obtained.

Details of COSM:

Y. Suzuki and K. Ikeda, PRC **38**, 410 (1988).

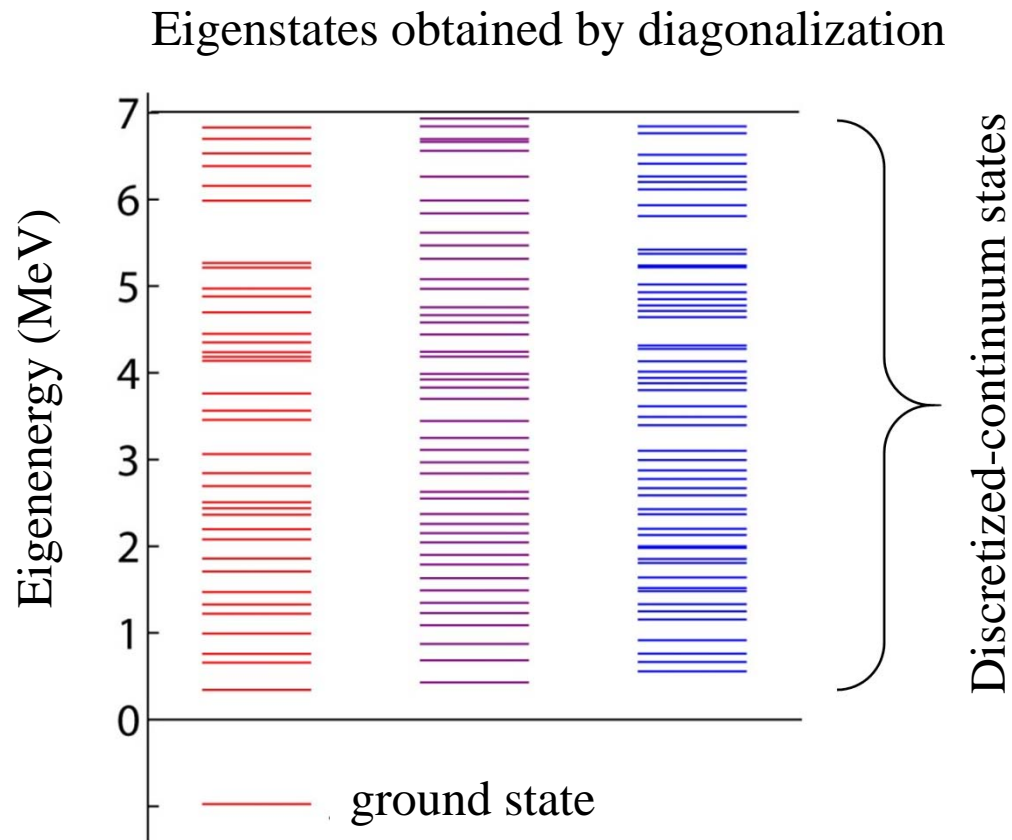
S. Aoyama *et al.*, PTP **116**, 1 (2006) [review].

T. Myo *et al.*, PL **B691**, 150 (2010) and references therein.



# COSM-CDCC for $^{22}\text{C}$ breakup by $^{12}\text{C}$

## Reaction part: Four-body CDCC



Set of the  $^{22}\text{C}$  internal wave functions  
(*basis functions* for the 4-body system)

$$\Psi^{\text{CDCC}} = \sum_{i=0}^{i_{\max}} \hat{\phi}_i \hat{\chi}_i$$

Relative motion between  $^{22}\text{C}$  and target  
(*expansion coefficients*)

Details of four-body CDCC:

T. Matsumoto *et al.*, PRC **70**, 061601(R) (2004); *ibid.* 73, 051602(R) (2006).

M. Rodriguez-Gallardo *et al.*, PRC **80**, 051601 (2009).

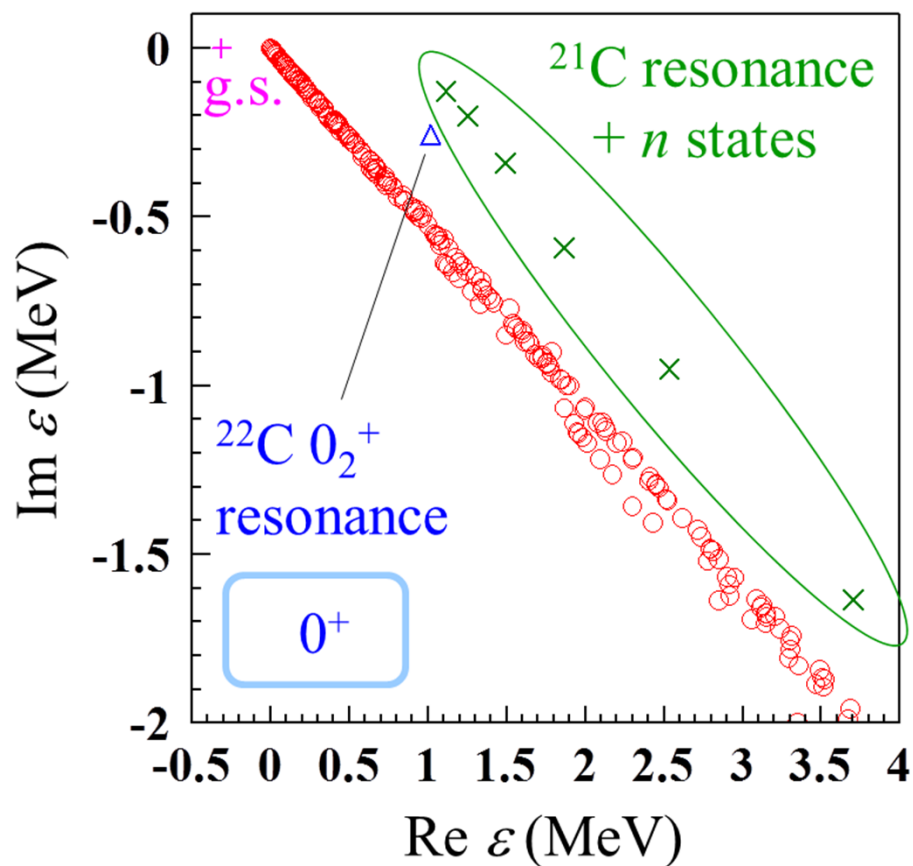
# CSM Smoothing

## (CSM: Complex-Scaling Method)

T. Matsumoto, Kato, and Yahiro, PRC **82**, 054602(R) (2010).

Eigenstates of  $H^\theta$

(complex-scaled Hamiltonian)



$$\tilde{\mathcal{T}}_i^\theta = \sum_n \langle \tilde{\phi}_i^\theta | C(\theta) | \Phi_n \rangle T_n^{\text{CDCC}}$$

index for the pseudostates  $\Phi_n$  used in CDCC

$$\frac{d\sigma}{d\epsilon} = \frac{1}{\pi} \text{Im} \sum_i \frac{\mathcal{T}_i^\theta \tilde{\mathcal{T}}_i^\theta}{\epsilon - \epsilon_i}$$

index for the eigenstates  $\phi_i^\theta$  of  $H^\theta$

# Numerical inputs

## $^{22}\text{C}$ wave function

- ✓ Minnesota force for  $n$ - $n$ , Woods-Saxon potential for  $n$ - $^{20}\text{C}$ .
- ✓  $s_{1/2}$ ,  $p_{3/2}$ ,  $p_{1/2}$ ,  $d_{5/2}$ ,  $d_{3/2}$ ,  $f_{7/2}$ ,  $f_{5/2}$ ,  $g_{9/2}$ ,  $g_{7/2}$ ,  $h_{11/2}$ , and  $h_{9/2}$  for the  $n$  s.p. orbit.
- ✓ Each orbit is described by 10 Gaussian basis functions.



D. R. Thompson *et al.*, NP **A286**, 53 (1977).

$0^+$  ground state with  $S_{2n} = 289$  keV, 604  $0^+$  and 1,385  $2^+$  PS

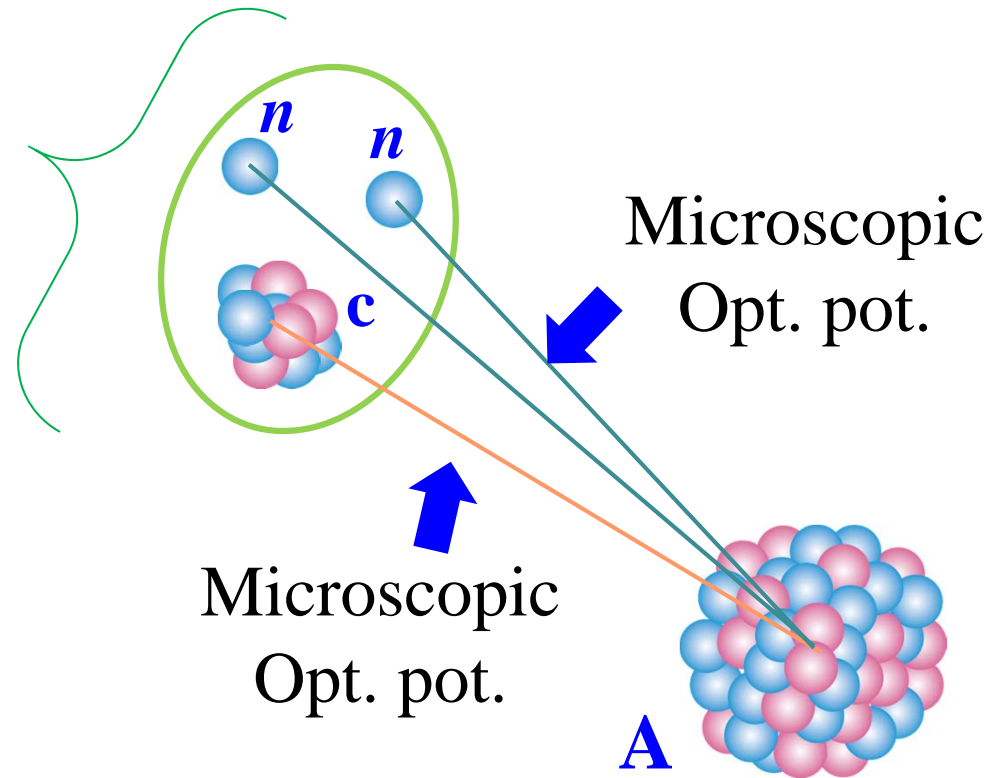
## $^{22}\text{C}$ - $^{12}\text{C}$ breakup reaction

- ✓ 77 ( $0^+$ ) + 164 ( $2^+$ ) PS below 10 MeV are included as breakup states of  $^{22}\text{C}$ .
- ✓ Distorting potentials are calculated by a microscopic folding model with CEG07 nucleon-nucleon  $g$  matrix.
- ✓ We adopt the so-called no-recoil approximation for the  $^{20}\text{C}$  core nucleus.

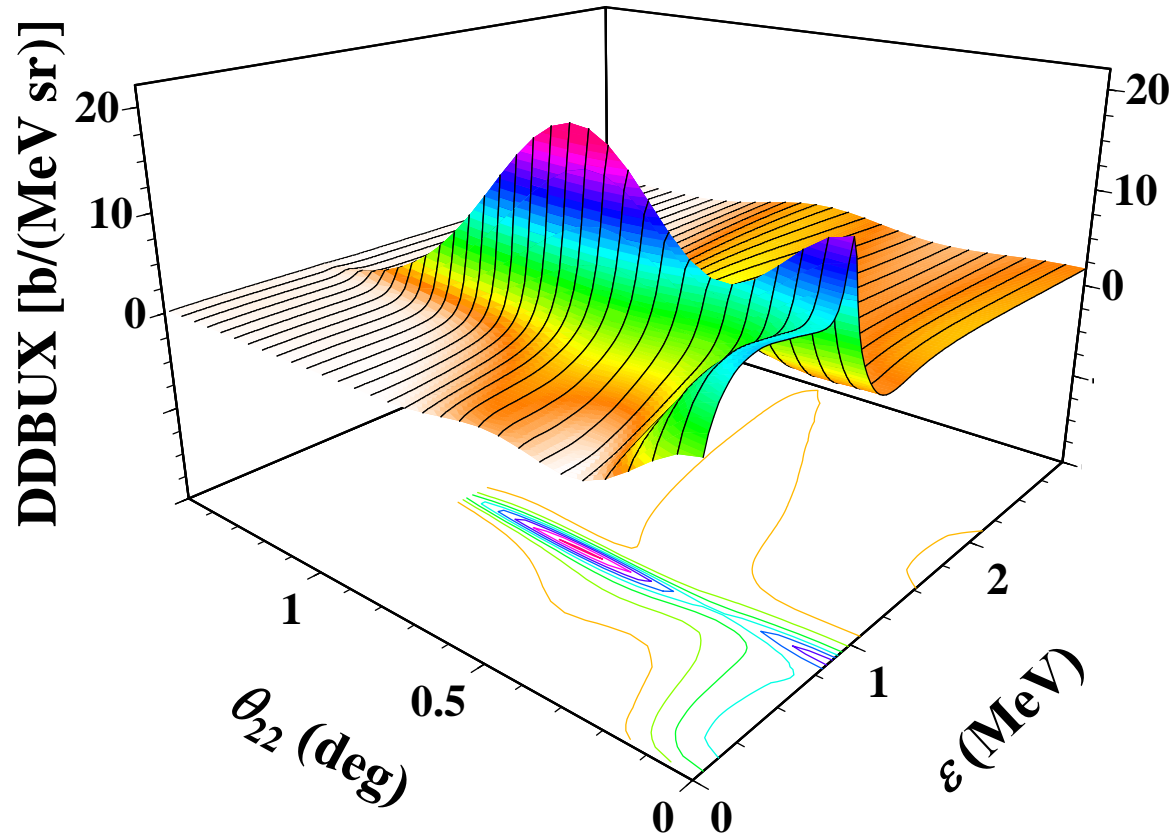
T. Furumoto *et al.*, PRC **78**, 044610 (2008).

# Microscopic CDCC

$n + n + c$  dynamics  
explicitly described



# DDBUX of $^{22}\text{C}$ by $^{12}\text{C}$



- ✓ The CSM smoothing\* is adopted to obtain the BUX.
- ✓ COSM predicts the following resonances:

$^{22}\text{C}$  resonance

$$0_2^+: 1.02 - i 0.52/2$$

$$2_1^+: 0.86 - i 0.10/2$$

$$2_2^+: 1.80 - i 0.26/2$$

$^{21}\text{C}$  resonance

$$d_{3/2}: 1.1 - i 0.10/2$$

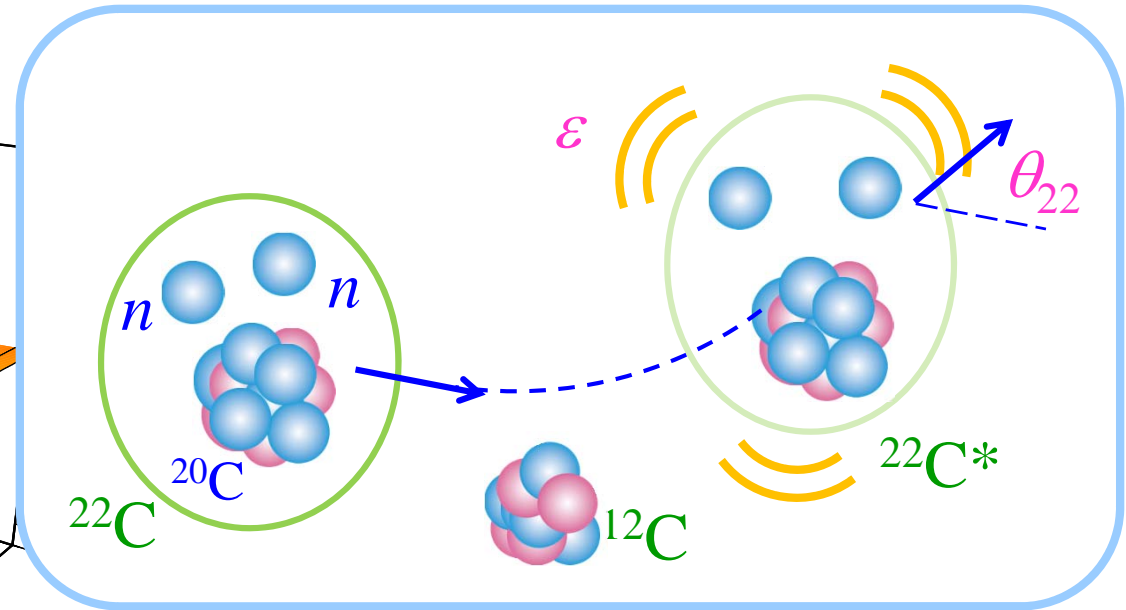
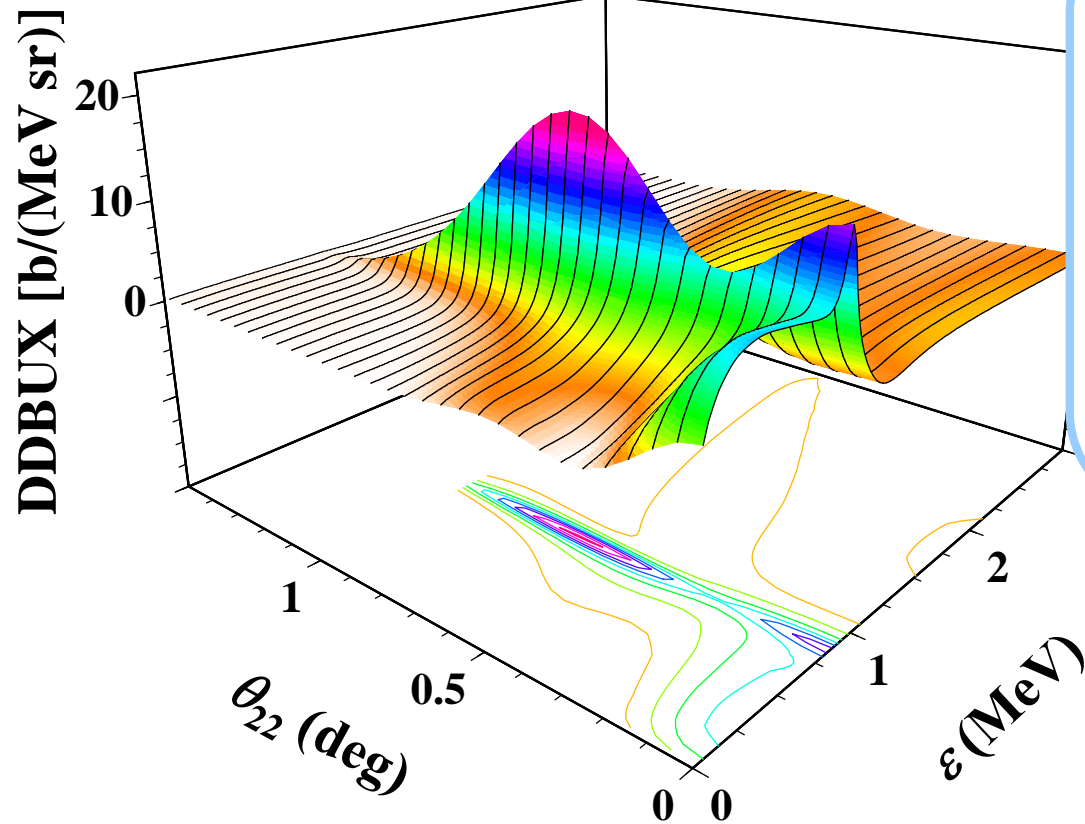


*How are these resonances observed?*

\*T. Matsumoto *et al.*, PRC **82**, 054602(R) (2010).



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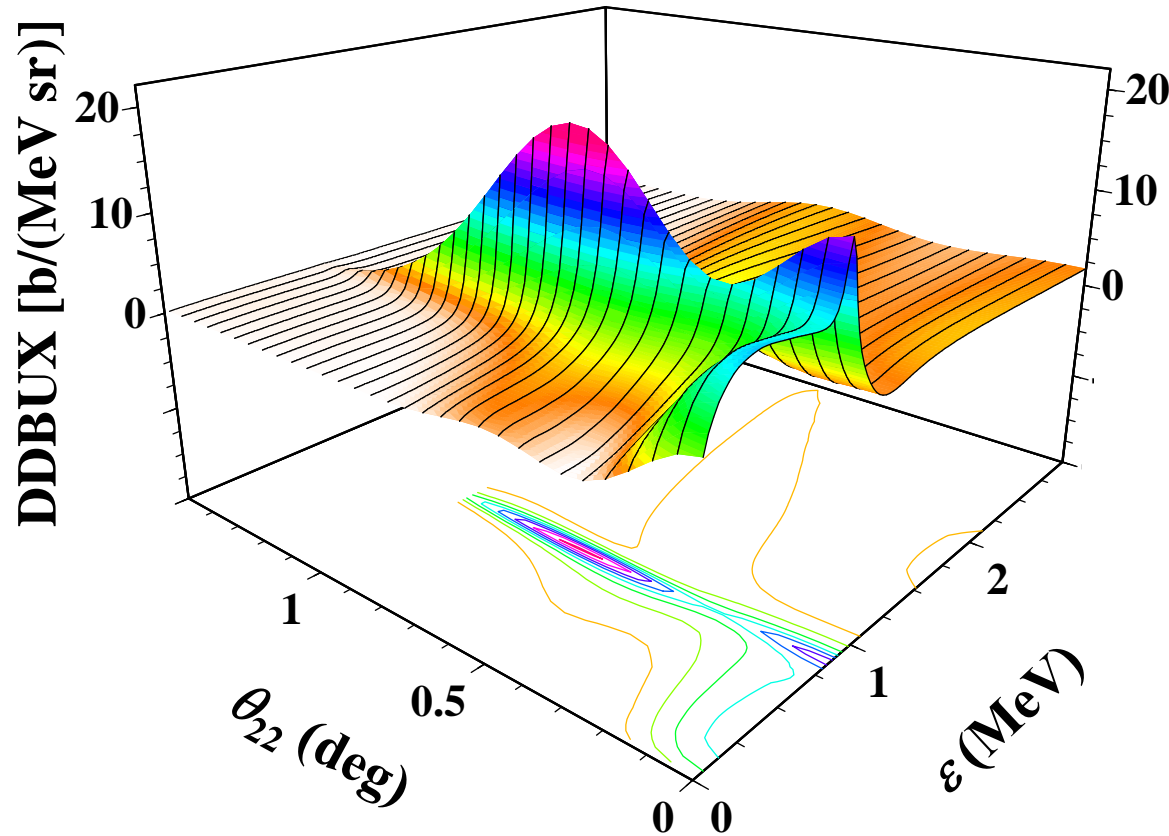
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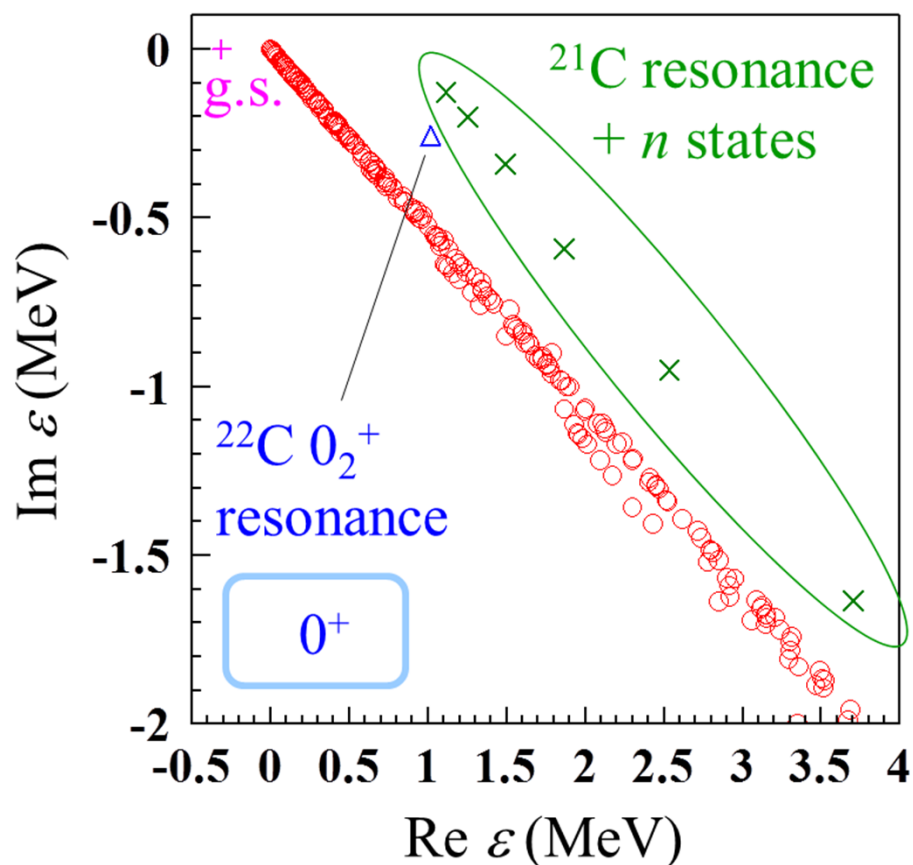
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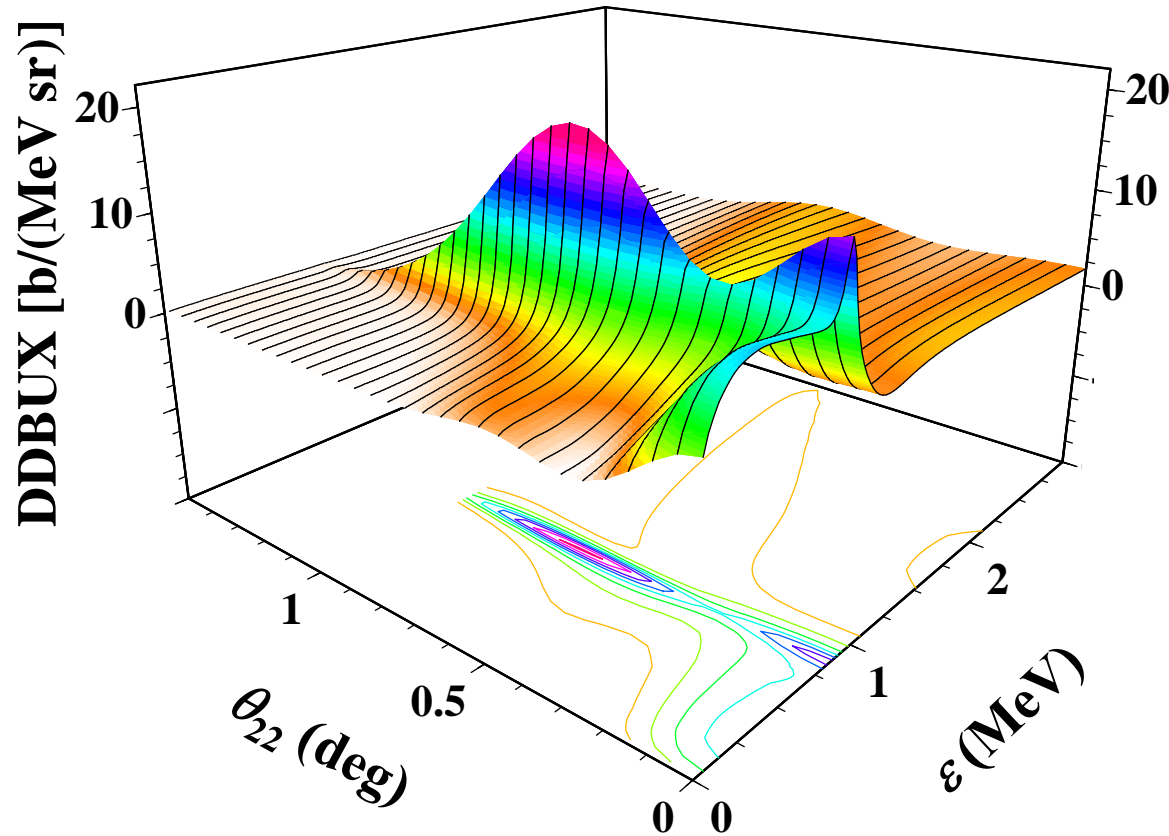
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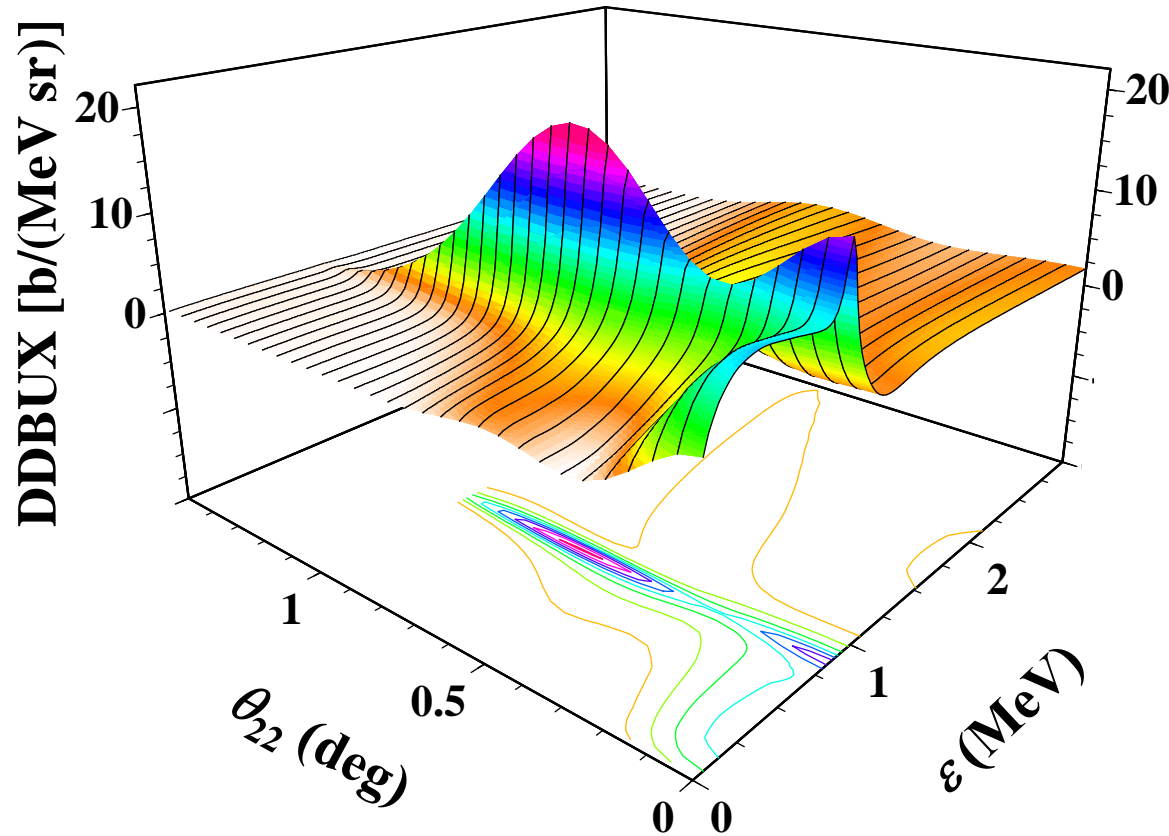
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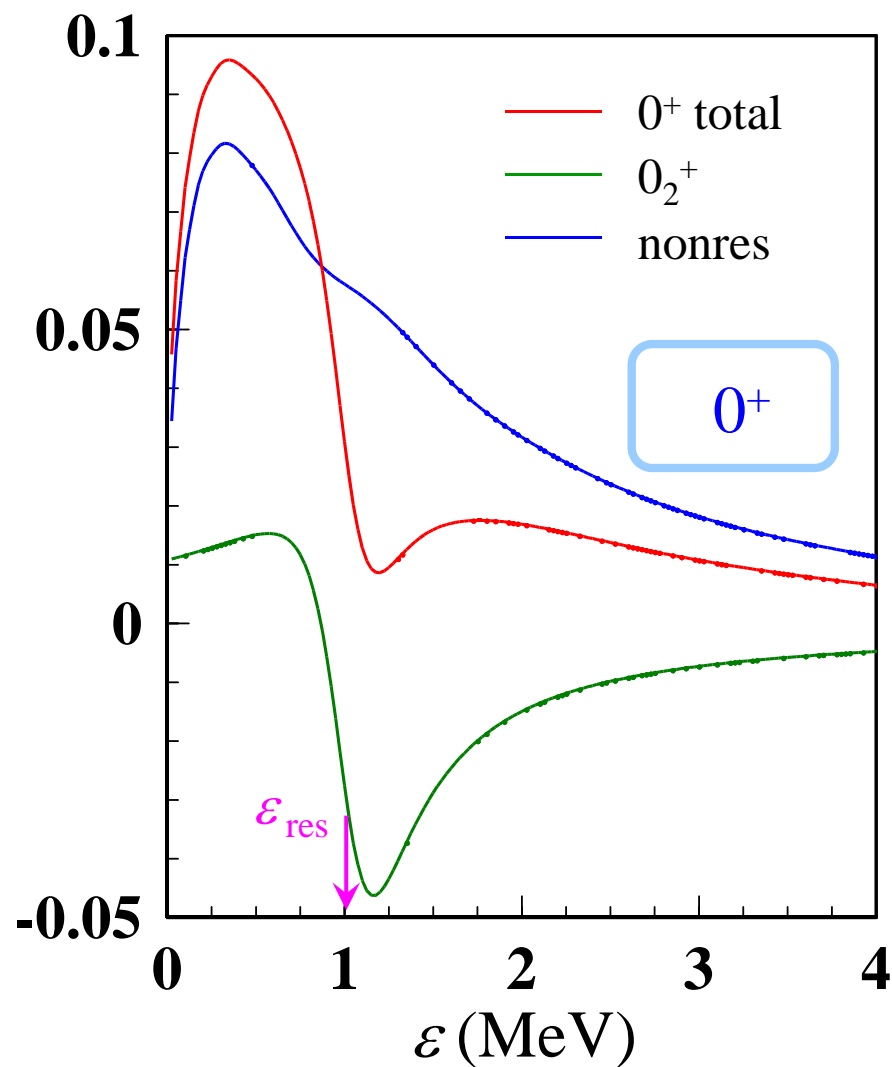
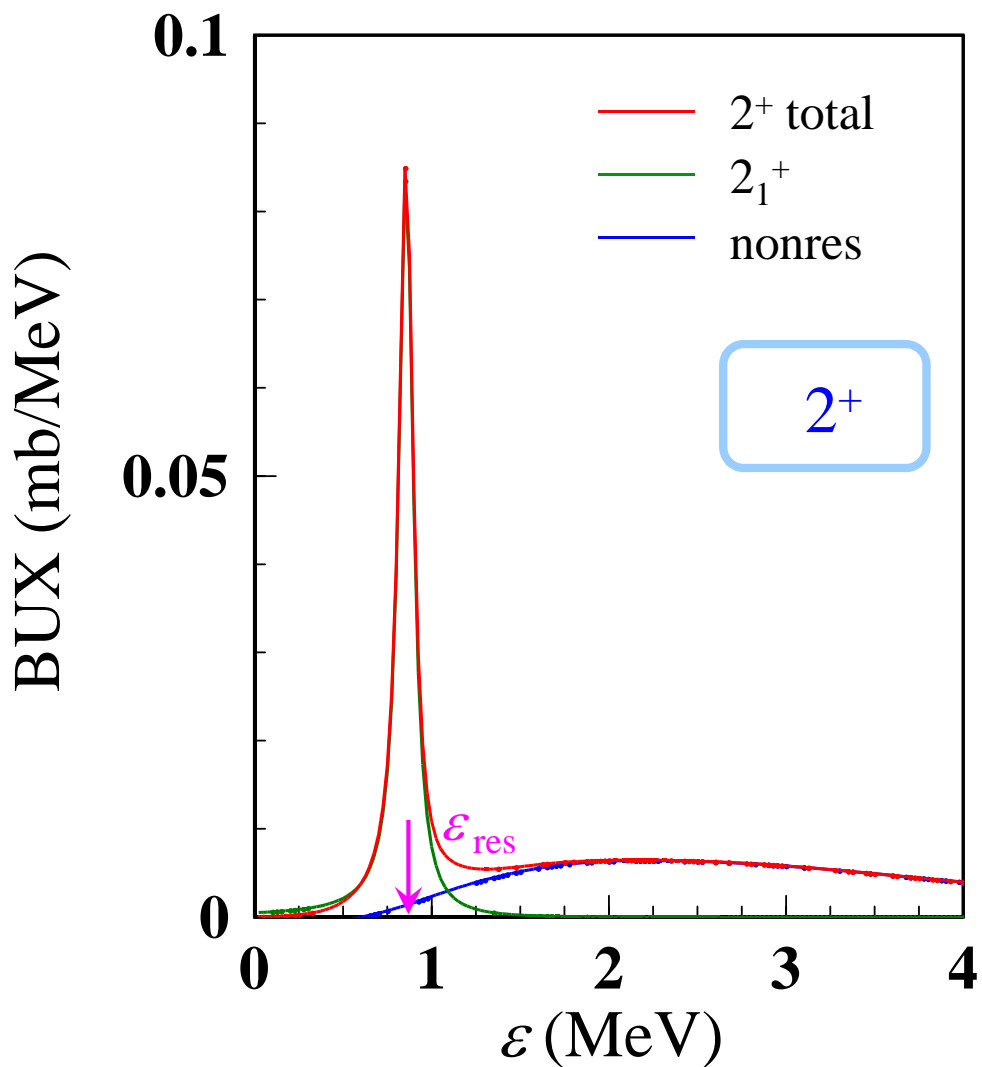
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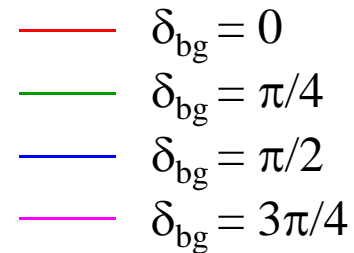
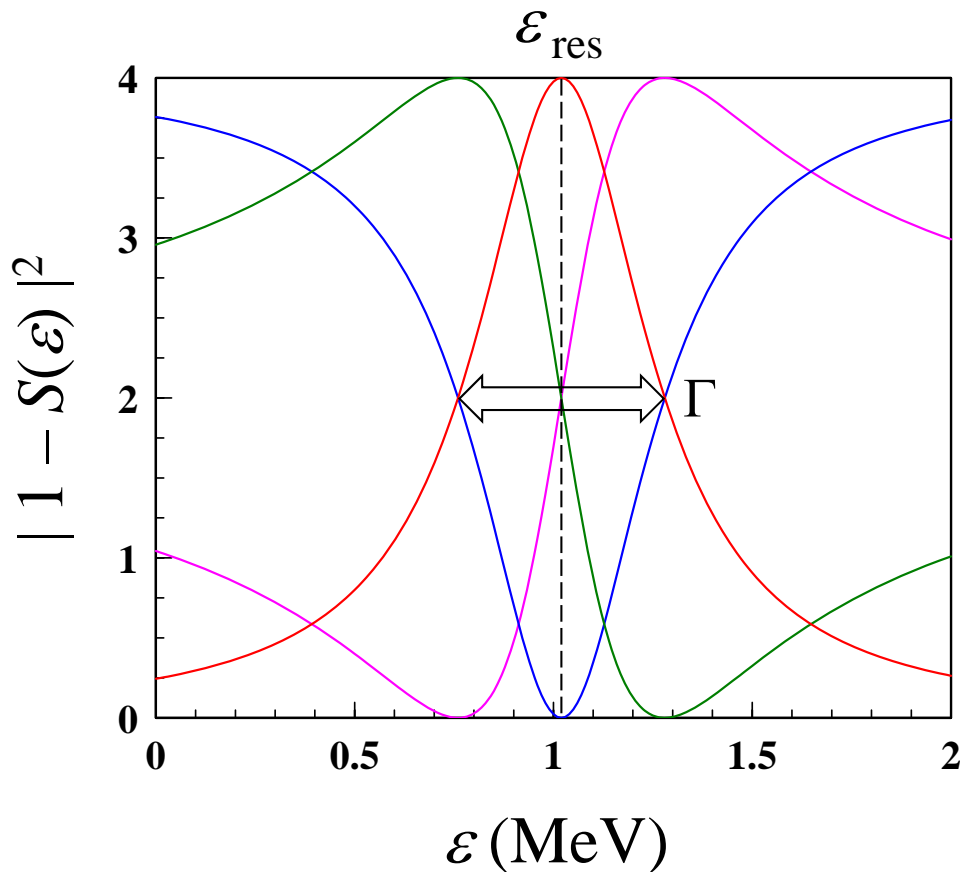
\*T. Matsumoto *et al.*, PRC **82**, 054602(R) (2010).

# Integrated BUX (0 – 0.1 deg)



- ✓ The narrow peak around 0.8 MeV is due to the  $2_1^+$  resonance of  $^{22}\text{C}$ .
- ✓ The shape of the  $0_2^+$  resonance is due to **background phase effect**.

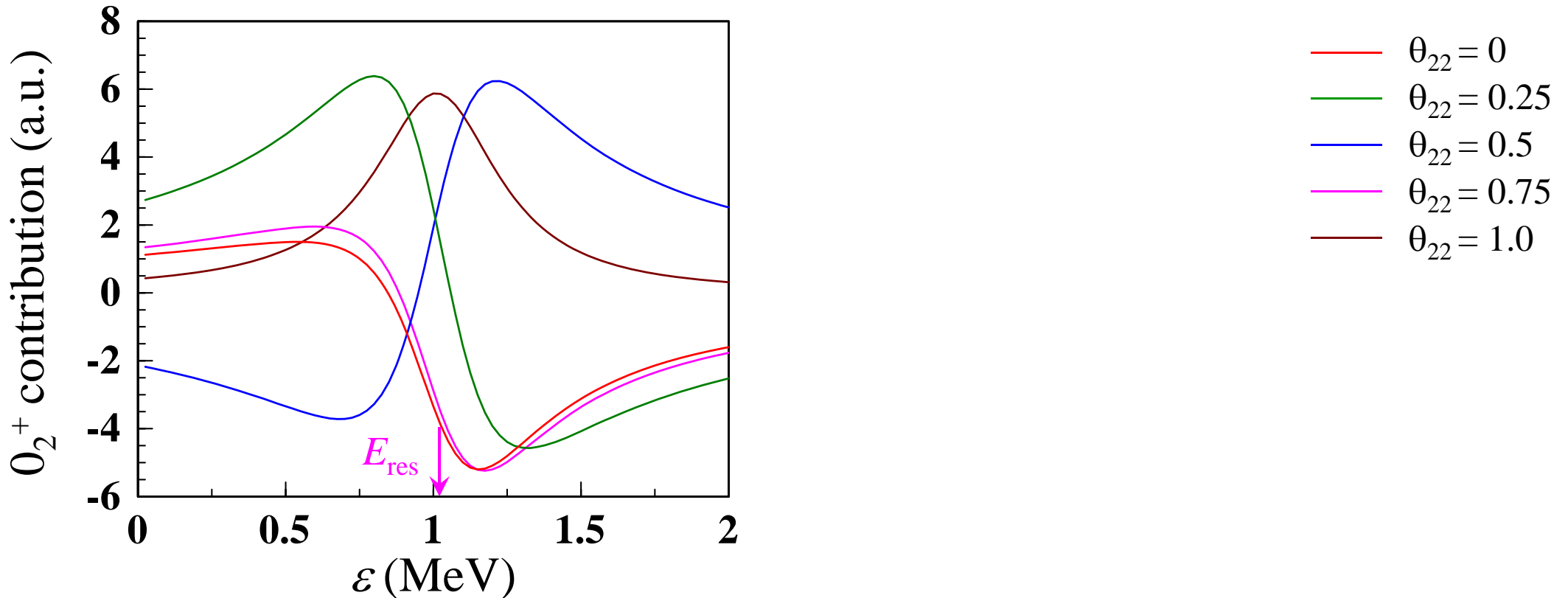
# BackGround Phase (BGP) effect



$$\begin{aligned}
 S(\epsilon) &= e^{2i\delta_{\text{bg}}(\epsilon) + 2i\delta_{\text{res}}(\epsilon)} \\
 &= e^{2i\delta_{\text{bg}}(\epsilon)} \frac{\epsilon - \epsilon_{\text{res}} - i\Gamma/2}{\epsilon - \epsilon_{\text{res}} + i\Gamma/2}
 \end{aligned}$$

- ✓ In nuclear physics, we **always** have  $\delta_{\text{bg}}$ .
- ✓ There are many examples of this effect in many research fields.
- ✓ In most cases, this effect is observed as **small changes** in the **resonance energy** and **width**.

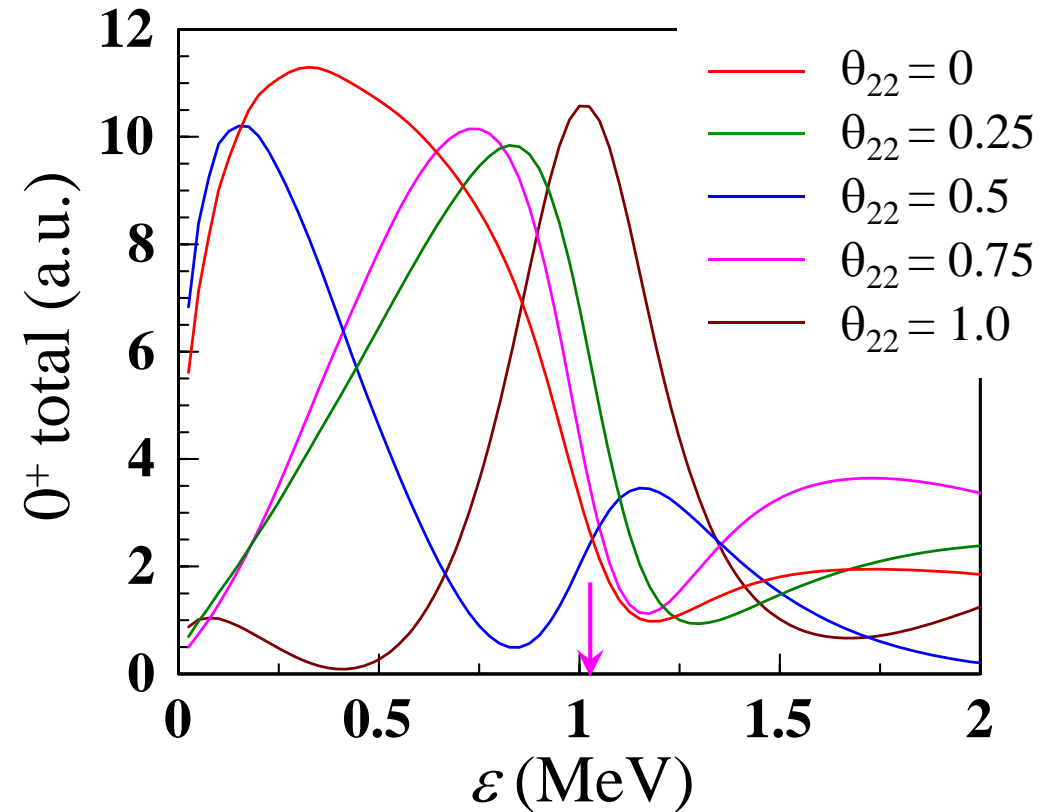
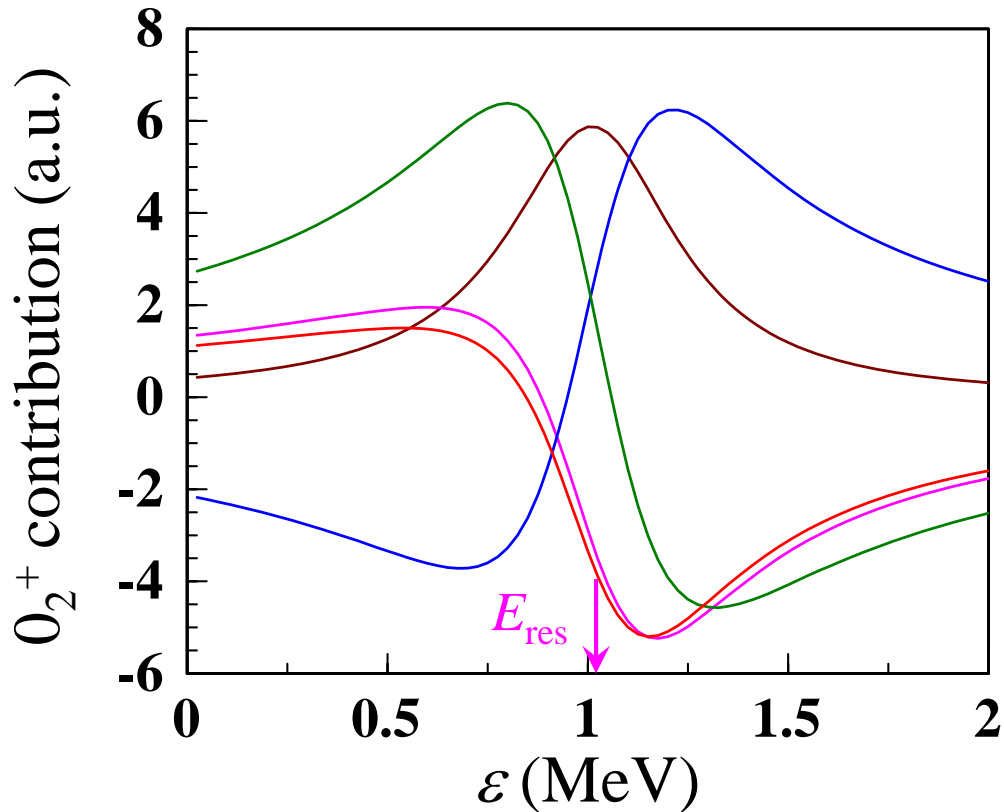
# BGP effect on the DDBUX



- ✓ The BGP effect is indeed **sizable**.
- ✓ We have a **variety of patterns** of the resonant (and  $0^+$ ) cross section.
- ✓ Appear in only the  $0^+$  state



# BGP effect on the DDBUX



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- ✓ We have a **variety of patterns** of the resonant (and  $O^+$ ) cross section.
- ✓ Appear in only the  $O^+$  state

# Summary of the 1<sup>st</sup> topic

What is the form of  $^{22}\text{C}^*$  in a breakup observable?

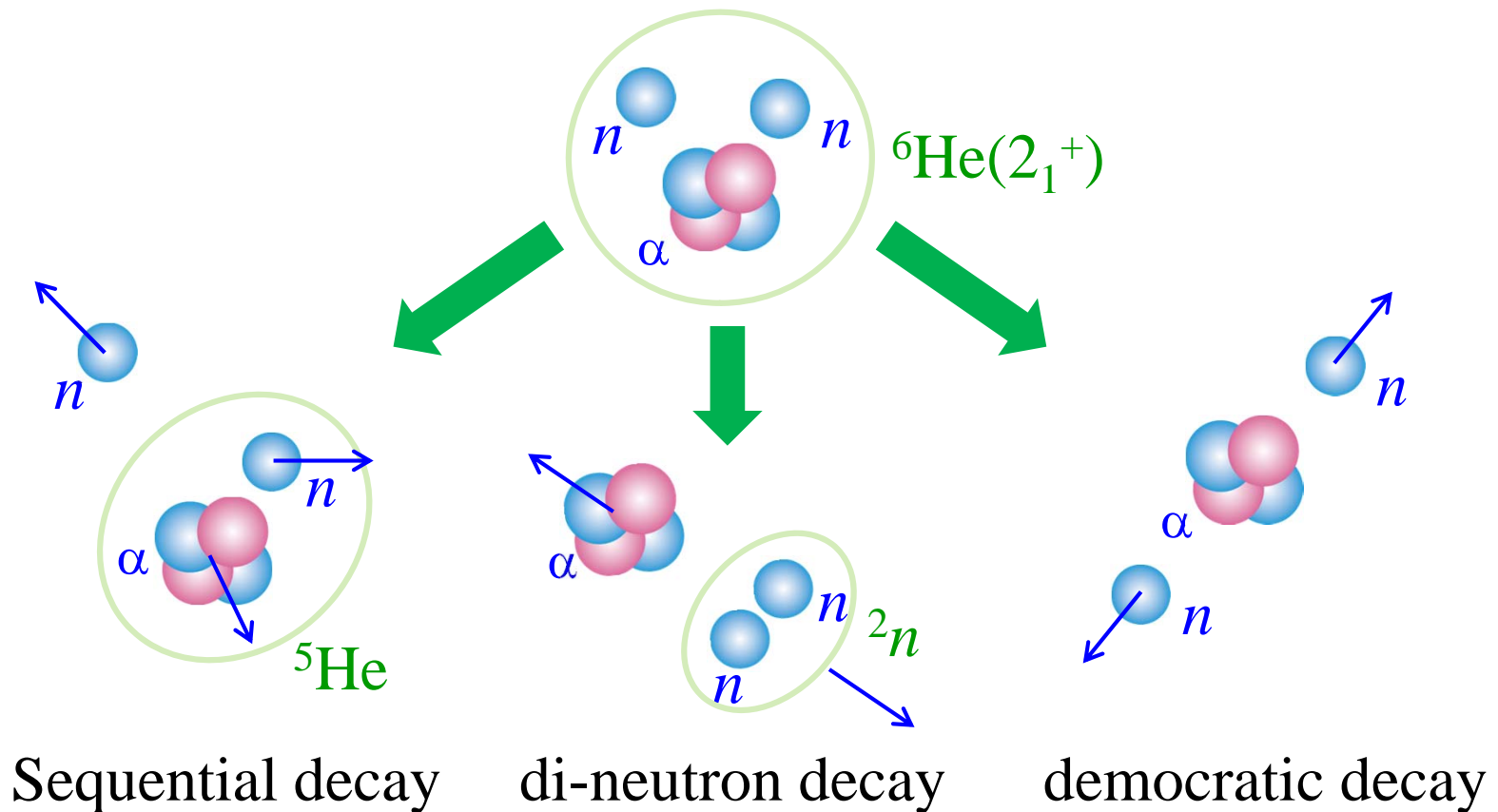
KO, Myo, Furumoto, Matsumoto, Yahiro, PRC88, 024616 (2013).

- ✓ The  $2_1^+$  state: Breit-Wigner form
- ✓ The  $0_2^+$  state: peculiar form due to the BGP effect (coexistence of the  $0^+$  resonant and nonresonant waves)
- ✓ The BGP has a strong scattering-angle dependence.
- ✓ We should be careful to identify the  $0_2^+$  state of  $^{22}\text{C}$  in the observables.

# 2<sup>nd</sup> topic

What is the decay mode of the  $2_1^+$  state of  ${}^6\text{He}$ ?

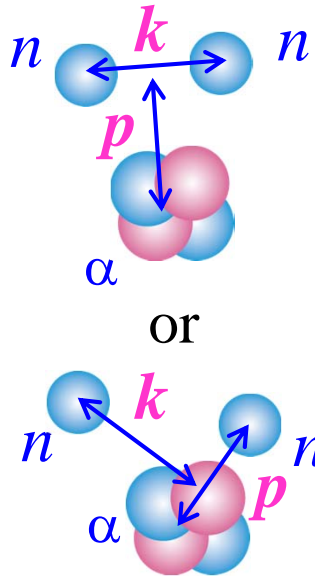
Y. Kikuchi, Matsumoto, Minomo, O, PRC88, 021602 (2013).



# CDCC-CSLS

✓ The method of Complex-Scaled solutions of the Lippmann-Schwinger Eq.

Y. Kikuchi, Myo, Takashina, Kato, Ikeda, PTP122, 499 (2009).

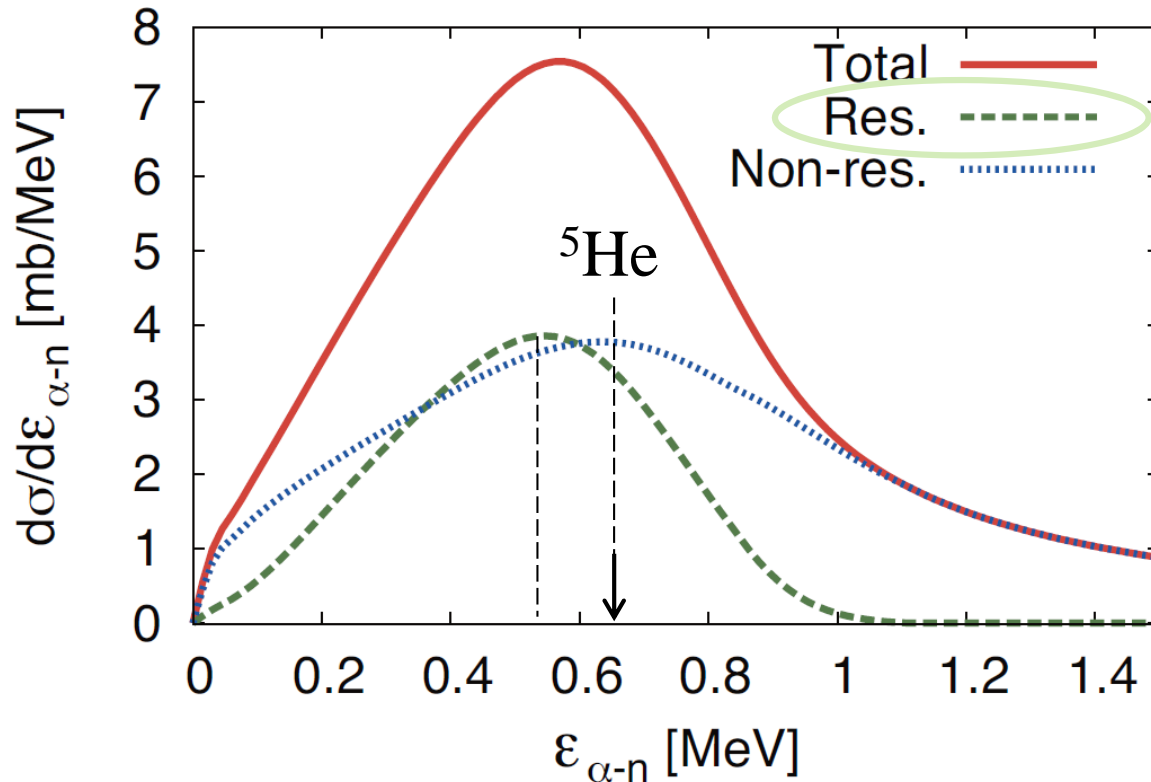
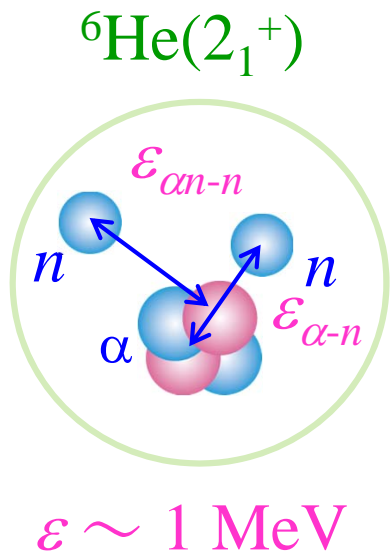


$$T(\mathbf{p}, \mathbf{k}) = \sum_n \langle \Phi^{(-)}(\mathbf{p}, \mathbf{k}) | \Phi_n \rangle T_n^{\text{CDCC}}$$

$$\equiv \sum_n f_n(\mathbf{p}, \mathbf{k}) T_n^{\text{CDCC}}$$

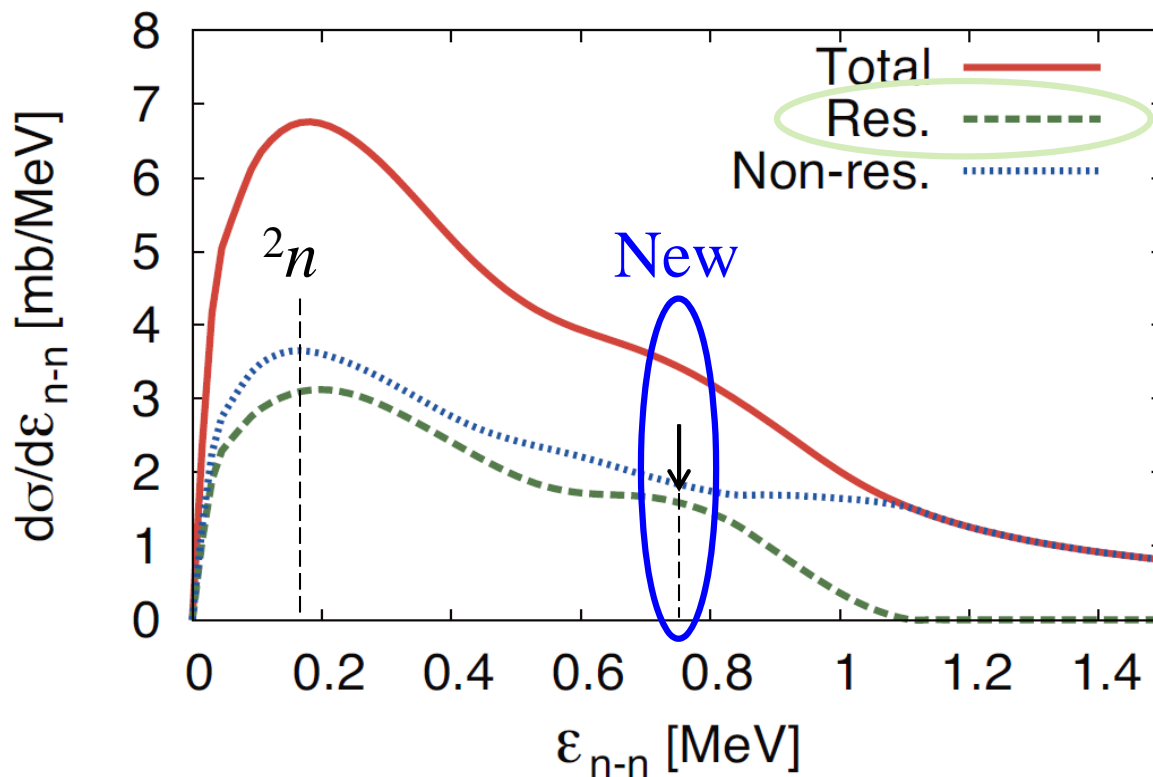
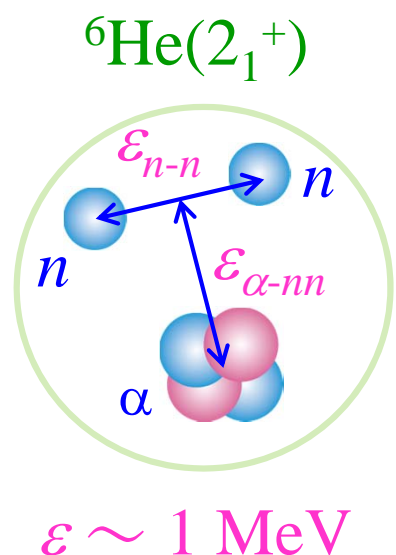
$$f_n(\mathbf{p}, \mathbf{k}) = \langle \varphi_{\text{free}}(\mathbf{p}, \mathbf{k}) | \Phi_n \rangle + \sum_i^f \langle \varphi_{\text{free}}(\mathbf{p}, \mathbf{k}) | V_{\alpha nn} C^{-1}(\theta) | \phi_i^\theta \rangle \\ \times \frac{1}{\varepsilon - \varepsilon_i^\theta} \langle \tilde{\phi}_i^\theta | C(\theta) | \Phi_n \rangle$$

# Sequential decay quenched



- ✓ When  $\varepsilon \sim 1 \text{ MeV}$  and  $\varepsilon_{\alpha-n} \sim 0.7 \text{ MeV}$ , the other neutron ( $\sim 0.3 \text{ MeV}$ ) hardly penetrates the centrifugal barrier ( $p$ -wave).
- ✓ The peak of the green line suggests the di-neutron decay or the democratic decay.

# Coexistence of two decay modes

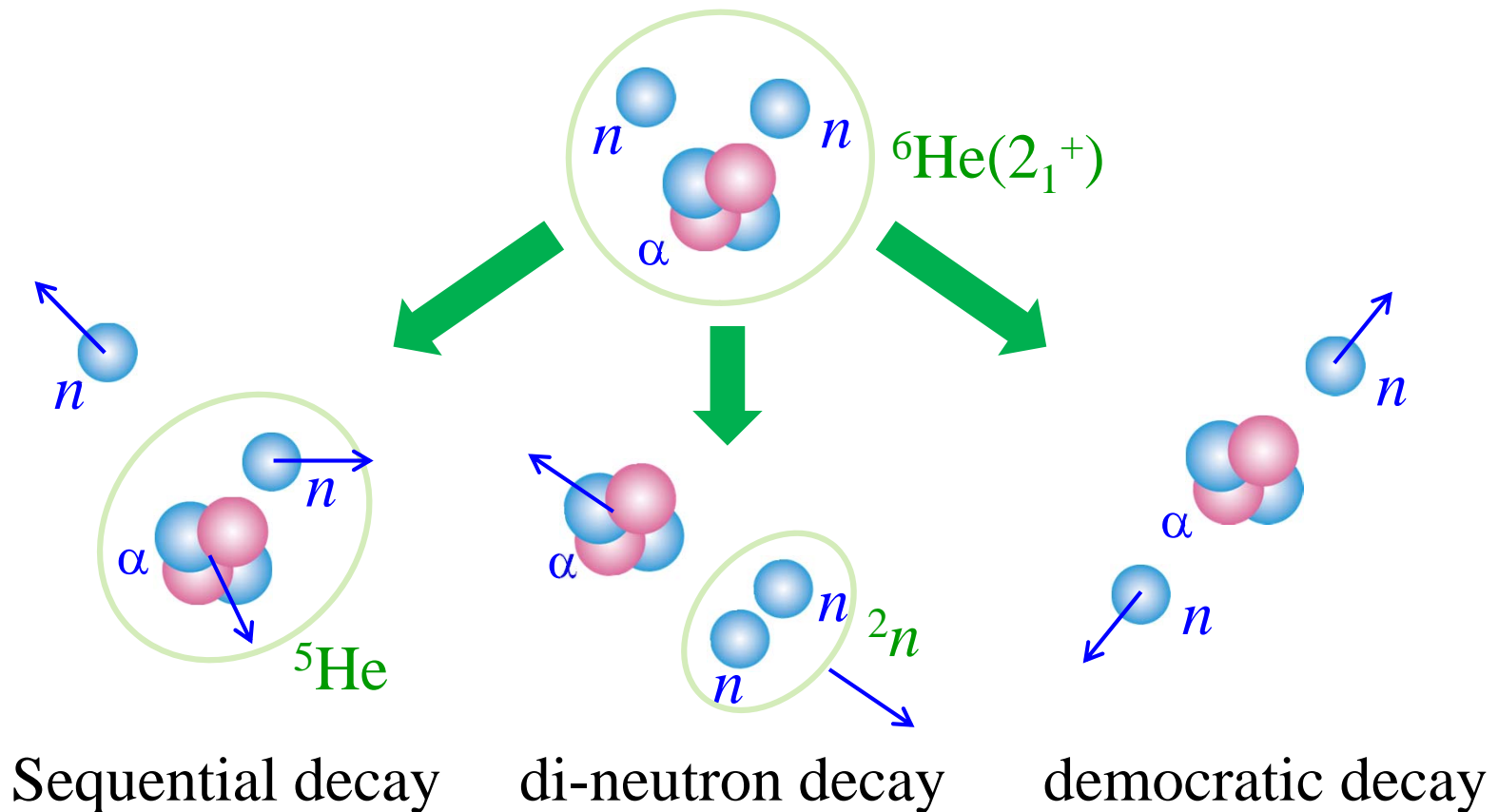


- ✓ The lower peak suggests the di-neutron decay due to the Fin. State Int. (FSI).
- ✓ The higher peak indicates the democratic decay.
- ➡ Decay of a di-neutron in the  $2_1^+$  state not due to the FSI.

# Summary of the 2<sup>nd</sup> topic

What is the decay mode of the  $2_1^+$  state of  ${}^6\text{He}$ ?

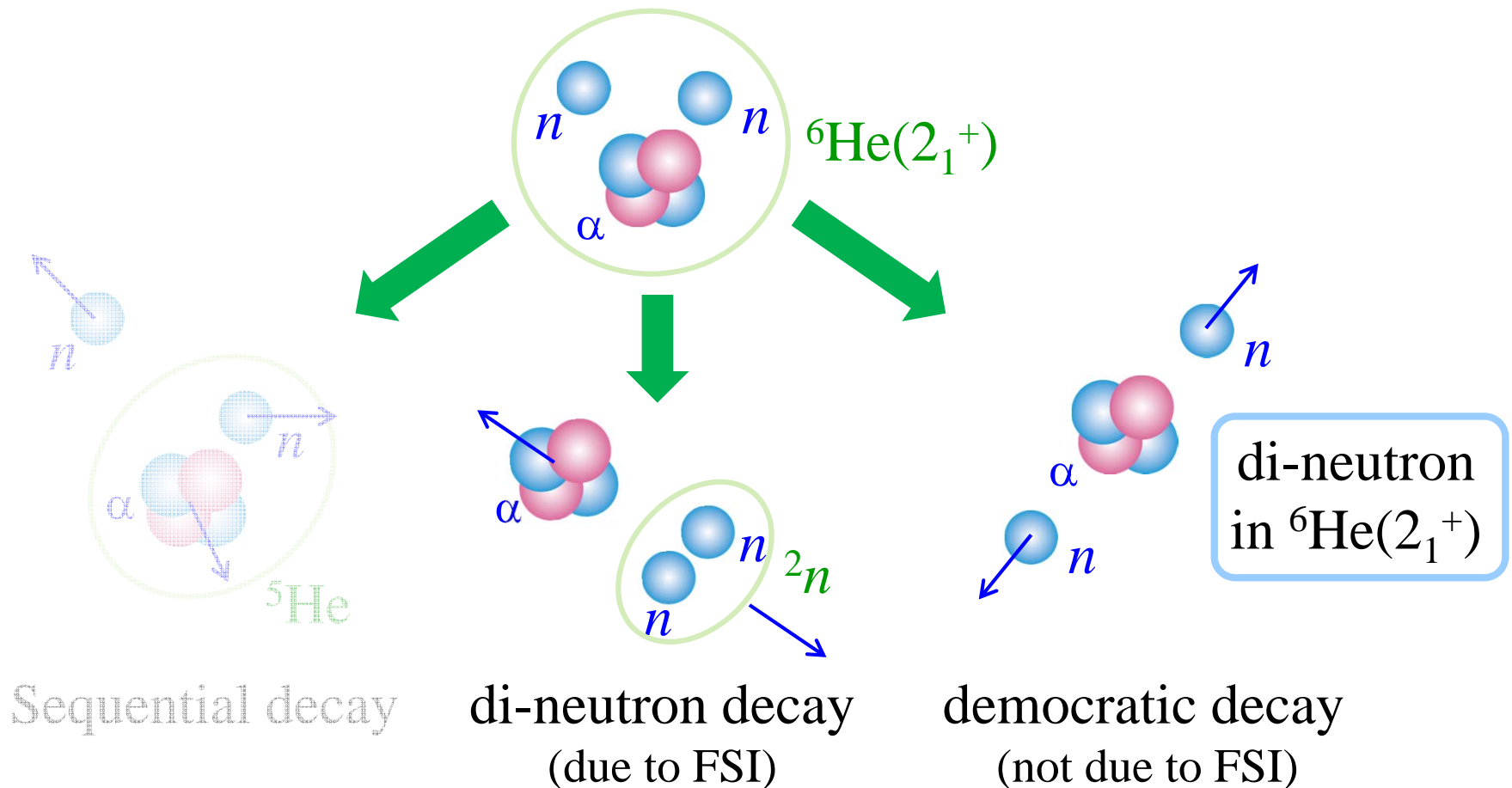
Y. Kikuchi, Matsumoto, Minomo, O, PRC88, 021602 (2013).



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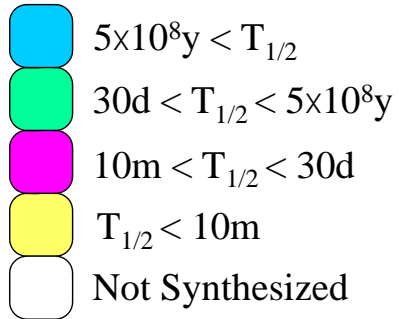




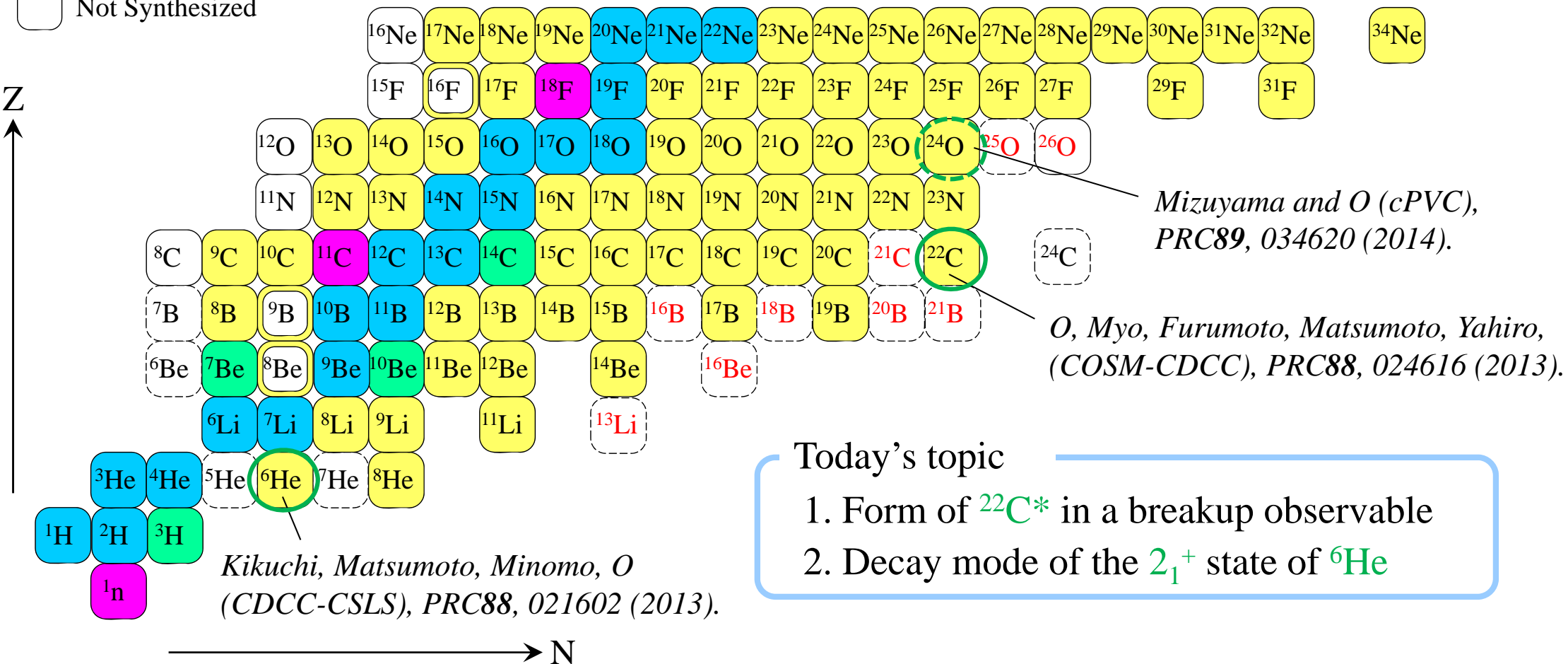
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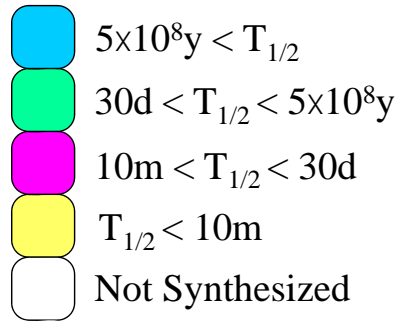
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