Extending the Eikonal Approximation to Low Energy

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## 2 Reaction modelling

- ODCC
- Eikonal approximation

#### Coulomb breakup of <sup>15</sup>C

- 68AMeV
- 20AMeV



## Halo nuclei

Exotic nuclear structures are found far from stability In particular halo nuclei with peculiar quantal structure :

- Light, n-rich nuclei
- Low  $S_n$  or  $S_{2n}$

### Exhibit large matter radius

due to strongly clusterised structure :

neutrons tunnel far from the core and form a halo

#### One-neutron halo <sup>11</sup>Be $\equiv$ <sup>10</sup>Be + n

$${}^{15}C \equiv {}^{14}C + n$$

#### Two-neutron halo

Proton halces are also possible, but less probable





## Reactions with halo nuclei

Halo nuclei are fascinating objects but difficult to study  $[\tau_{1/2}(^{6}\text{He})=0.8 \text{ s}]$  $\Rightarrow$  require indirect techniques, like reactions

Elastic scattering

[see A. Di Pietro on Wednesday]

Breakup ≡ dissociation of halo from core by interaction with target [see A. Bonaccorso on Thursday]

Need an good understanding of the reaction mechanism i.e. an accurate theoretical description of reaction coupled to a realistic model of projectile

# Framework

Projectile (P) modelled as a two-body system : core (c)+loosely bound nucleon (f) described by

- $H_0 = T_r + V_{cf}(\boldsymbol{r})$
- $V_{cf}$  adjusted to reproduce bound state  $\Phi_0$  and resonances

Target T seen as structureless particle



*P-T* interaction simulated by optical potentials  $\Rightarrow$  breakup reduces to three-body scattering problem :

$$\left[T_R + H_0 + V_{cT} + V_{fT}\right]\Psi(\boldsymbol{r},\boldsymbol{R}) = E_T\Psi(\boldsymbol{r},\boldsymbol{R})$$

with initial condition  $\Psi(\mathbf{r}, \mathbf{R}) \xrightarrow[Z \to -\infty]{} e^{iKZ + \cdots} \Phi_0(\mathbf{r})$ 

# Continuum Discretised Coupled Channel (CDCC) Solve the three-body scattering problem :

$$\left[T_R + H_0 + V_{cT} + V_{fT}\right]\Psi(\boldsymbol{r},\boldsymbol{R}) = E_T\Psi(\boldsymbol{r},\boldsymbol{R})$$

by expanding  $\Psi$  on eigenstates of  $H_0$ 

$$\Psi(\mathbf{r}, \mathbf{R}) = \sum_i \chi_i(\mathbf{R}) \Phi_i(\mathbf{r})$$
 with  $H_0 \Phi_i = \epsilon_i \Phi_i$ 

Leads to set of coupled-channel equations (hence CC) :

$$[T_R + \epsilon_i + V_{ii}]\chi_i + \sum_{j \neq i} V_{ij}\chi_j = E_T\chi_i,$$

with  $V_{ij} = \langle \Phi_i | V_{cT} + V_{fT} | \Phi_j \rangle$ The continuum has to be discretised (hence CD) [Austern *et al.*, Phys. Rep. 154, 125 (1987)] [Tostevin, Nunes, Thompson, PRC 63, 024617 (2001)]

### Fully quantal approximation

No approximation on *P*-*T* motion, nor restriction on energy But expensive computationally (at high energies)

# Eikonal approximation

Three-body scattering problem :

$$\left[T_R + H_0 + V_{cT} + V_{fT}\right]\Psi(\boldsymbol{r},\boldsymbol{R}) = E_T\Psi(\boldsymbol{r},\boldsymbol{R})$$

with condition  $\Psi \mathop{\longrightarrow}\limits_{Z \to -\infty} e^{i K Z} \Phi_0$ 

Eikonal approximation : factorise  $\Psi = e^{iKZ}\widehat{\Psi}$ 

$$T_R \Psi = e^{iKZ} [T_R + vP_Z + \frac{\mu_{PT}}{2} v^2] \widehat{\Psi}$$

Neglecting  $T_R$  vs  $P_Z$  and using  $E_T = \frac{1}{2}\mu_{PT}v^2 + \epsilon_0$ 

$$i\hbar v \frac{\partial}{\partial Z} \widehat{\Psi}(\boldsymbol{r}, \boldsymbol{b}, Z) = [H_0 - \epsilon_0 + V_{cT} + V_{fT}] \widehat{\Psi}(\boldsymbol{r}, \boldsymbol{b}, Z)$$

solved for each *b* with condition  $\widehat{\Psi} \underset{Z \to -\infty}{\longrightarrow} \Phi_0(\mathbf{r})$ This is the dynamical eikonal approximation (DEA) [Baye, P. C., Goldstein, PRL 95, 082502 (2005)]

# <sup>15</sup>C+Pb @ 68AMeV : energy distribution



• Excellent agreement between cbcc and DEA

[P.C., Esbensen and Nunes, PRC 85, 044604 (2012)]

• Excellent agreement with experiment

[Nakamura et al. PRC 79, 035805 (2009)]

 $\Rightarrow$  Confirms the validity of the approximations

 $\ldots$  and the two-body structure of  $^{15}C$ 

<sup>15</sup>C+Pb @ 68AMeV : angular distribution



- DEA agrees well with CDCC ٩
- Though a slight shift compared to cpcc...



DEA too high and too forward due to lack of Coulomb deflection [P.C., Esbensen and Nunes, PRC 85, 044604 (2012)]

Can E-CDCC solve the problem?

#### 20AMeV

# Eikonal-CDCC (E-CDCC)

Solving the eikonal problem expanding  $\Psi$  upon  $H_0$  eigenstates  $\Phi_i(\mathbf{r})$ assuming discretised continuum [Ogata et al. PRC 68, 064609 (2003)]

$$\Psi(\boldsymbol{r},\boldsymbol{R}) = \sum_{i} \xi_{i}(\boldsymbol{b},Z) \Phi_{i}(\boldsymbol{r}) e^{i\{K_{i}Z+\eta_{i}\ln[K_{i}R-K_{i}Z]\}}$$

 $\Rightarrow$  set of coupled equations

$$\frac{\partial}{\partial Z}\xi_i(\boldsymbol{b},Z) = \frac{1}{i\hbar v_i(R)}\sum_{i'}\mathcal{F}_{ii'}(\boldsymbol{b},Z)\xi_{i'}(\boldsymbol{b},Z)e^{i(K_{i'}-K_i)Z}\mathcal{R}_{ii'}(\boldsymbol{b},Z),$$

with coupling potential

$$\mathcal{F}_{ii'}(\boldsymbol{b}, Z) = \left\langle \Phi_i \left| V_{cT} + V_{fT} - V_{C} \right| \Phi_{i'} \right\rangle_{\boldsymbol{r}}.$$

E-CDCC takes proper account of energy conservation :  $v_i(R)$  $\mathcal{R}_{ii'}(b,Z) = \frac{(K_{i'}R - K_{i'}Z)^{m_{i'}}}{(K - K - K \cdot Z)^{i\eta_i}}$  accounts for part of the Coulomb distortion





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- hybrid solution : cbcc at low L (b) and eikonal at large L (b)  $\Rightarrow$  excellent agreement with full cpcc



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- Improve eikonal using Coulomb correction :  $b \rightarrow b'$ •

# Summary and prospect

- Halo nuclei exhibit very exotic structure : core + halo
- Studied mostly through reactions
  - elastic scattering
  - breakup
- Mechanism of reactions with halo nuclei understood cpcc & eikonal agree at 70AMeV
- At 20AMeV, eikonal fails, due to Coulomb deflection But :
  - ► E-CDCC can be extended to hybrid version
    ⇒ agreement with full CDCC
  - Simple Coulomb correction works fine Can it be improved ?

How far down can we go in energy ? HIE-ISOLDE ?