RIKEN Nishina Center Monthly Colloquium (14th February 2012)

重イオン核反応のための微視的相互作用モデル

Microscopic interaction models for HI reactions

ー現状と今後の展望ー

- present status & future perspective -

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- 1. Introduction
- *2. Microscopic theory for nucleus-nucleus interaction with a new complex G-matrix interaction, CEG07*
 - I. Application to proton-nucleus scattering
 - II. Application to heavy-ion (HI) scattering/reactions
 ⇒ Importance of repulsive three-body force effect
- *3. Attractive-to-repulsive transition of HI optical potential around E/A=200~300 MeV*
- 4. Global optical potential for exotic heavy ions
- *5. future perspectives*

1. Introduction

- Understanding the interactions between composite nuclei (AA interactions), starting from NN interaction :
 - ✓ one of the most fundamental subjects in nuclear physics



- one of the key issues to understand various nuclear reactions:
 - optical potentials: elastic scattering
 - distorting potentials as doorway to various reactions (inelastic, transfer, knockout, breakup ···)
- ✓ important to survey unknown nuclear structures/reaction of <u>unstable nuclei</u> far from stability lines (N>>Z, Z>>N), for which
 - few/no elastic-scattering data & phenom. potential information is available.

Uopt(R) = Vopt(R) + i Wopt(R) : complex potential

- Phenomenolocical optical potentials:
 needs Exp. Data (elastic scattering) to determine potential parameters (e.g. Woods-Saxon form)
 - ✓ optical potential for heavy-ion systems (AA) has large ambiguity in depth & shape due to strong absorption (in most cases)
 - ✓ → only sensitive to potential at nuclear surface







How deep is the nucleus-nucleus potential at short distances ?



✓ Can we probe the depth ?

 In general, it is very difficult to probe the <u>central depth</u> of the H.I. potentials, due to <u>strong absorption</u>.

Can we probe H.I. potential at short distances?

→ Yes, we can! (at least for light heavy-ions)

by the measurements of refractive scattering at <u>high-q region</u> (backward) , such as <u>nuclear-rainbow</u> phenomena.



By the way,

Q: Why do we need to know the <u>central depth</u> of the potential ?

A: We can study the property of <u>high-density nuclear matter</u>, such as that in <u>neutron stars</u>, in laboratory experiments.



Hubble Space Telescope Wide Field Planetary Camera 2 赤色巨星・超巨星における 元素合成の終焉(Fe, Ni)

- → 重力崩壊
- → 高密度核物質/超新星爆発
- → 中性子星

ハイペロン星/クォーク星

高密度核物質の性質、 特に、 非圧縮率 (Incompressibility) が重要

 $\boldsymbol{K_{\rm nm}} = 9\rho_0^2 \frac{\boldsymbol{d}^2 (\boldsymbol{E} / \boldsymbol{A})}{\boldsymbol{d}^2 \boldsymbol{\rho}}$



How can we probe the property of high-density nuclear matter at \rho > \rho_0?





2

 ρ/ρ_0

з

20

0

-20



 If the nuclear matter is soft, the central depth of the potential may become deep. But, good quality of exp. data are not always in our hands.

→ We need a microscopic theory that explains & predicts

- correct <u>depth & shape</u> of **heavy-ion optical potentials**, (hopefully , of both the <u>real</u> and <u>imaginary</u> parts)
- ✓ including **unstable nuclei** (n-rich & p-rich isotopes)
- ✓ correct <u>energy dependence</u> over the wide range of incident energy, up <u>to a few hundred MeV/u</u>

starting from bare NN interaction in free space

2. Microscopic theory for nucleus-nucleus interaction with a new complex G-matrix interaction, CEG07

Breuckner Theory (G-matrix theory)

- 無限核物質中での有効相互作用を導出する理論
- 密度pの核物質中で、Bethe-Goldstone方程式を解く
 ⇒媒質効果(Pauli effects, Binding effect etc.)を考慮した ladder diagram をすべて足しあげる。

$$G(\omega) = V + \sum_{\boldsymbol{q}_1, \boldsymbol{q}_2} V \frac{Q(\boldsymbol{q}_1, \boldsymbol{q}_2)}{\omega - e(q_1) - e(q_2) + i\varepsilon} G(\omega)$$

 $Q(oldsymbol{q}_1,oldsymbol{q}_2)$: Pauli-Projection Operator

Complex G-matrix interaction (CEG07)

T.Furumoto, Y. Sakuragi and Y. Yamamoto, Phys.Rev.C 78 (2008) 044610

"<u>ESC04</u>" : the latest version of Extended Soft-Core force designed for *NN*, *YN* and *YY* systems

Th. Rijken, Y. Yamamoto, Phys.Rev.C 73 (2006) 044008

<u>References:</u>

- √ T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC<u>78</u> (2008) 044610,
- ✓ T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC<u>79</u> (2009) 011601(R),
- ✓ T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC<u>80</u> (2009) 044614
- ✓ T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC82 (2010) 029908(E)
- ✓ T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC<u>82</u> (2010) 044612

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Λ

- 1. Three-body attraction (三体引力)
 - Fujita-Miyazawa diagram
 - important at <u>low density</u> region

2. Three-body repulsion (三体斥力)

- originated the triple-meson correlation
- important at <u>high-density</u> region

In the ESC04 model

 \Rightarrow density-dependent effective two-body force

saturation curve in nuclear matter with G-matrix interaction (CEG07)



有限核への適用 ⇒ 無限系で求めた有効核力を、有限系の核子密度でfolding



I. Application to proton-nucleus scattering





Target

$$U_{C}(\mathbf{R}) = \int \rho(\mathbf{r}) T_{D}(\mathbf{R}, \mathbf{r}; k_{F}, E) d\mathbf{r}$$
$$+ \int \rho(\mathbf{R}, \mathbf{r}) T_{EX}(\mathbf{R}, \mathbf{r}; k_{F}, E) \exp(i\mathbf{k}_{0} \cdot \mathbf{s}) d\mathbf{r}$$
$$= V_{C}(\mathbf{R}) + iW_{C}(\mathbf{R})$$

Complex G-matrix interaction (CEG07)

$$T_{D,EX} = T_{D,EX}^{(real)} + i T_{D,EX}^{(imag)}$$

Single folding Potential
(LS part)
$$T^{(LS)}(s) \qquad r$$
Proton
$$T^{(LS)}(s) \qquad r$$

$$Target$$

$$U_{LS}(\mathbf{R})\ell \cdot \sigma = \sum_{i} \int \varphi_{i}^{*}(\mathbf{r})T_{D}^{(LS)}(\mathbf{R},\mathbf{r};k_{F},E)\mathbf{L} \cdot \mathbf{S}\varphi_{i}(\mathbf{r})d\mathbf{r}$$

$$+ \sum_{i} \int \varphi_{i}^{*}(\mathbf{r})T_{EX}^{(LS)}(\mathbf{R},\mathbf{r};k_{F},E)\mathbf{L} \cdot \mathbf{S}\varphi_{i}(\mathbf{R})\exp(i\mathbf{k}_{0} \cdot \mathbf{s})d\mathbf{r}$$

$$= (V_{LS}(\mathbf{R}) + iW_{LS}(\mathbf{R}))\ell \cdot \sigma$$

Complex G-matrix interaction (CEG07)

$$T_{D,EX}^{(LS)} = T_{D,EX}^{(LS,real)} + iT_{D,EX}^{(LS,imag)}$$

Renormalization of the imaginary part strength



So, we renormalize (suppress) the imaginary part strength $V(\mathbf{R}) + \underline{iN_W}W(\mathbf{R}) + \left(V_{LS}(\mathbf{R}) + \underline{iN_W}W_{LS}(\mathbf{R})\right)\ell \cdot \sigma$

Renormalized factor N_W is fixed to reproduce measured total reaction cross sections

$$V(\mathbf{R}) + i N_W W(\mathbf{R}) + (V_{LS}(\mathbf{R}) + i N_W W_{LS}(\mathbf{R})) \ell \cdot \sigma$$



CEG07a(two-body only) $N_W = 0.60$ CEG07b(TBR+TBA) $N_{W} = 0.60$ **CEG07c**(**TBR+TBA+** ω) $N_W = 0.65$ **CEG86** (two-body only) $N_{W} = 0.80$

CEG07 folding-model cal. of proton scattering by ¹²C



T.Furumoto, Y. Sakuragi and Y. Yamamoto, Phys.Rev.C 78 (2008) 044610

CEG07 folding-model cal. of proton scattering by ⁴⁰Ca



T.Furumoto, Y. Sakuragi and Y. Yamamoto, Phys.Rev.C 78 (2008) 044610

CEG07 folding-model cal. of proton scattering by ⁹⁰Zr & ²⁰⁸Pb



T.Furumoto, Y. Sakuragi and Y. Yamamoto, Phys.Rev.C 78 (2008) 044610

<u>II. Application to heavy-ion scattering/reractions</u>



有限核への適用 ⇒ 無限系で求めた有効核力を、有限系の核子密度でfolding



有限核への適用 ⇒ 無限系で求めた有効核力を、有限系の核子密度でfolding



T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC<u>79</u> (2009) 011601(R),
 T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC<u>80</u> (2009) 044614
 T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC82 (2010) 029908(E)
 T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC<u>82</u> (2010) 044612

$$\underline{Double folding Potential}_{with complex-G (CEG07)}
 U(\mathbf{R}) = \int \rho_1(\mathbf{r}_1) \rho_2(\mathbf{r}_2) g_D(\mathbf{s}; \rho, E) d\mathbf{r}_1 d\mathbf{r}_2
 projectile(P)
 target(T)
 + \int \rho_1(\mathbf{r}_1, \mathbf{r}_1 - \mathbf{s}) \rho_2(\mathbf{r}_2, \mathbf{r}_2 + \mathbf{s}) g_{EX}(\mathbf{s}; \rho, E) \exp\left[i\frac{\mathbf{K} \cdot \mathbf{s}}{M}\right] d\mathbf{r}_1 d\mathbf{r}_2
 = V_{DFM}(\mathbf{R}) + iW_{DFM}(\mathbf{R})$$

Complex G-matrix interaction (CEG07)

$$g_{D,EX} = g_{D,EX}^{(real)} + i g_{D,EX}^{(imag)}$$

T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC<u>79</u> (2009) 011601(R),
 T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC<u>80</u> (2009) 044614
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$$\frac{Double folding Potential}{with complex-G (CEG07)}$$

$$U(\mathbf{R}) = \int \rho_1(\mathbf{r}_1) \rho_2(\mathbf{r}_2) g_D(\mathbf{s}; \rho, E) d\mathbf{r}_1 d\mathbf{r}_2$$

$$projectile(P)$$

$$target(T)$$

$$+ \int \rho_1(\mathbf{r}_1, \mathbf{r}_1 - \mathbf{s}) \rho_2(\mathbf{r}_2, \mathbf{r}_2 + \mathbf{s}) g_{EX}(\mathbf{s}; \rho, E) \exp\left[i\frac{\mathbf{K} \cdot \mathbf{s}}{M}\right] d\mathbf{r}_1 d\mathbf{r}_2$$

$$= V_{DFM}(\mathbf{R}) + iW_{DFM}(\mathbf{R})$$

✓ Renormalization factor for the imaginary part

$$\rightarrow U_{DFM} = V_{DFM} + i N_W W_{DFM}$$

¹⁶O+¹⁶O: bench-mark system to test DFM

Frozen-density approx. (FDA) $\rho = \rho_1 + \rho_2$

¹⁶O+¹⁶O Elastic Scattering at *E/A* = 70 MeV : DFM with CEG07 ⇒ decisive effect of Three-body force (mainly TBR) is clearly observed !



¹⁶O+¹⁶O: bench-mark system to test DFM

Frozen-density approx. (FDA) $\rho = \rho_1 + \rho_2$

¹⁶O+¹⁶O Elastic Scattering at *E/A* = 70 MeV : DFM with CEG07 ⇒ decisive effect of Three-body force (mainly TBR) is clearly observed !



¹²C + ¹²C elastic scattering

 $N_{W} = 0.5$

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¹⁶O + ¹²C, ²⁸Si, ⁴⁰Ca



T.Furumoto, Y. Sakuragi, Y. Yamamoto, (Phys. Rev. C 80, 044614 (2009)
<u>Applications of microscopic FMP to</u>

1. reaction calculations (CC, CDCC etc.)

2. scattering of unstable nuclei

Microscopic Coupled Channel (MCC) with CEG07

Coupled Channel equation

$$\begin{bmatrix} T_R + U_{\alpha\alpha}(\mathbf{R}) - E_{\alpha} \end{bmatrix} \chi_{\alpha}(\mathbf{R}) = -\sum_{\beta \neq \alpha}^N U_{\alpha\beta}(\mathbf{R}) \ \chi_{\beta}(\mathbf{R})$$
$$U_{\alpha\beta}(\mathbf{R}) = \int \rho_{00}^{(P)}(\mathbf{r}_1) \rho_{\alpha\beta}^{(T)}(\mathbf{r}_2) v_{NN}(\mathbf{s};\rho, E) d\mathbf{r}_1 d\mathbf{r}_2 = V_{\alpha\beta} + iW_{\alpha\beta}$$
transition density **CEG07**



¹⁶O + ¹⁶O inelastic scattering studied by a complex *G-matrix interaction* (@ E/A=70 MeV)

by M.Takashina, T. Furumoto, Y.Sakuragi PRC 81, 047605 (2010)

CC cal. with complex-G (CEG07) $U_{ii}^{DFM}(\mathbf{R}) = V_{ii}^{DFM}(\mathbf{R}) + iN_{W}W_{ii}^{DFM}(\mathbf{R})$



^{9,11}Li + ¹²C "quasi-elastic" scattering

• ⁹Li density : proton, neutron \Rightarrow single Gaussian form

$$\binom{R_{\rm r.m.s}^{p}}{R_{\rm r.m.s}^{n}} = 2.18 \,({\rm fm})$$

 $R_{\rm r.m.s}^{n} = 2.39 \,({\rm fm})$

$$\bullet^{11} \text{Li density : } {}^{9} \text{Li} + \text{di-neutron model}$$

$$\rho^{({}^{11}\text{Li})}(r) = \left\langle \psi_0(R) \mid \rho^{({}^{9}\text{Li})}(\vec{r} - \frac{2}{11}\vec{R}) + \rho^{(2n)}(\vec{r} + \frac{9}{11}\vec{R}) \mid \psi_0(R) \right\rangle$$

$$R_{\rm r.m.s} = 3.16 \, ({\rm fm})$$

Y.Hirabayashi, S.Funada and Y. Sakuragi (Proceedings of International Symposium on Structure and Reactions of Unstable Nuclei, pp227-pp232 (1991))



^{9,11}Li + ¹²C "quasi-elastic" scattering



 $U_{ij}^{DFM}(\mathbf{R}) = V_{ij}^{DFM}(\mathbf{R}) + iN_{W}W_{ij}^{DFM}(\mathbf{R})$

9,11Li + 12C "quasi-elastic" scattering E/A ~ 60 MeV



Exp. data : J. J. Kolata et al., (Phys. Rev. Lett. 69 (1993) 2631

⁶Li elastic scattering with ⁶Li $\rightarrow \alpha$ +d break-up

 \Rightarrow <u>Continuum-Discretized Coupled-Channels (CDCC) method</u>





$$U_{ij}^{DFM}(\mathbf{R}) = V_{ij}^{DFM}(\mathbf{R}) + iN_{W}W_{ij}^{DFM}(\mathbf{R})$$

Y. Sakuragi, M, Ito, Y. Hirabayashi, C. Samanta (Prog. Theor. Phys. 98 (1997) 521)

elastic scattering of ⁶Li by ¹²C, ²⁸Si at E/A = 53 MeV

CDCC cal. with complex-G (CEG07) folding model

 $U_{ij}^{DFM}(\mathbf{R}) = V_{ij}^{DFM}(\mathbf{R}) + iN_{W}W_{ij}^{DFM}(\mathbf{R})$



3. attractive-to-repulsive transition of the nucleus-nucleus potentials with the increase of collision energy



原子核間相互作用が E/A=100~400 MeVの 間で引力→斥力に転移

丁度、RIBFで多くの核反応実験 が行われるエネルギー領域

T. Furumoto, Y. Sakuragi, Y. Yamamoto, Phys. Rev. C <u>82</u>, 044612 (2010)

◆ <u>NN interaction :</u>

long-range attraction
 short-range repulsive core

nucleon-nucleus (NA) interaction :

attractive at low energies (E<200 MeV)
 wine-bottle-bottom (WBB) around transitional energies repulsive at high energies (E>500 MeV)



L.G.Arnold, (Phys.Rev.C25(1982)936



similar behavior to NA int. f(*d*-*A*) ~ *f*(*p*-*A*)+ *f*(*n*-*A*)



nucleon-nucleus (NA) interaction :

attractive at low energies (E<200 MeV)
 wine-bottle-bottom (WBB) around transitional energies repulsive at high energies (E>500 MeV)

Y. Sakuragi, M. Tanifuji, NPA560, 945(1993).



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Q: How about in nucleus-nucleus systems?





Q: How about optical potential for heavy ions?

- A: according to the predictions of microscopic theory,
- ✓ <u>attractive-to-repulsive transition</u> occurs ?
 - → Yes, but thus far we have no experimental evidence.
- ✓ if so, in what energy region?
 - \rightarrow the transition occurs around E/A = 300~400 MeV
- ✓ how can we observe the transition, if it really occurs?

→ measure the evolution of elastic scattering <u>angular distribution</u> with increasing energy in the energy range of $E/A = 200 \sim 400 \text{ MeV}.$

✓ what are the new ingredients we can learn, if we observe the transition?

 → ① repulsive three-body force (TBF) in nuclear medium
 & ② tensor force effects besides the genuine repulsive core of NN int.

¹²C + ¹²C elastic scattering at $E/A = 100 \sim 400 \text{ MeV}$

 \succ real potential becomes repulsive around *E*/A = 300 \sim 400 MeV





¹²C + ¹²C elastic scattering at $E/A = 100 \sim 400 \text{ MeV}$



S-matirx elements of the ${}^{12}C + {}^{12}C$ elastic scattering at $E/A = 100 \sim 400$ MeV with CEG07b (with TBF effects)



NN tensor force plays an essential role in the attractive-to-repulsive transition of the A-A potentials



spin(S) and isospin(T)
components VST
of folding potential

★ (S,T) = (0,0) and (0,1) do not include the tensor force.

★ (S,T) = (1,0) and (1,1) components include the tensor force,

T. Furumoto, Y. Sakuragi, Y. Yamamoto, PRC<u>82</u> (2010) 044612

¹²C + ¹²C elastic scattering at $E/A = 100 \sim 400 \text{ MeV}$



Summary & Conclusion of part 1

complex G-matrix folding model with a <u>new G-matrix</u> <u>CEG07</u> predicts that

- <u>attractive-to-repulsive transition</u> occurs also in <u>heavy-ion</u> optical potentials *around* <u> $E/A = 300 \sim 400 MeV$ </u> → *but, no experimental evidence* → **BIG CHALLENGE**!
- ✓ can be observed by *measuring the energy-evolution of elastic* scattering <u>angular distribution</u> in the energy range of $E/A = 200 \sim 400 \text{ MeV}.$
- ✓ new ingredients we have learnt are the important roles of
 - *I repulsive three-body force (TBF) in nuclear medium*
 - ② tensor force effects

 \checkmark

4. Global optical potential for heavy ions systems including exotic nuclei

Global potential for projectiles of unstable nuclei up to driplines

Global optical potential for nucleus-nucleus systems from 50 MeV/u to 400 MeV/u T. Furumoto, W. Horiuchi, M. Takashina, Y. Yamamoto, Y. Sakuragi (submitted to PRC, Feb.2012)





for which

few/no elastic-scattering data & phenom. potential information is available.



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Global optical potential for nucleus-nucleus systems from 50 MeV/u to 400 MeV/u

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Global parameterization of the CEG07 folding-model potentials ✓ projectiles : Z = 6 (C isotope) ~ 20 (Ca isotope) (even-even) ✓ targets : ¹²C ~ ²⁰⁸Pb (closed or sub-closed shell nuclei) ✓ energy range : E/A = 30 ~ 400 MeV

Folding-model potential with CEG07a, CEG07b

$$\begin{split} U_{\rm D}(R) &= \int \rho_1(r_1)\rho_2(r_2)v_{\rm D}(s;\rho,E/A)dr_1dr_2 \\ &= \int \{\rho_1^{(\rm p)}(r_1)\rho_2^{(\rm p)}(r_2)v_{\rm D}^{(\rm pp)}(s;\rho,E/A) + \rho_1^{(\rm p)}(r_1)\rho_2^{(\rm n)}(r_2)v_{\rm D}^{(\rm pn)}(s;\rho,E/A) \\ &+ \rho_1^{(\rm n)}(r_1)\rho_2^{(\rm p)}(r_2)v_{\rm D}^{(\rm np)}(s;\rho,E/A) + \rho_1^{(\rm n)}(r_1)\rho_2^{(\rm n)}(r_2)v_{\rm D}^{(\rm nm)}(s;\rho,E/A)\}dr_1dr_2, \end{split}$$

$$\begin{split} U_{\rm EX}(R) &= \int \rho_1(r_1,r_1+s)\rho_2(r_2,r_2-s)v_{\rm EX}(s;\rho,E/A)\exp\left[\frac{ik(R)\cdot s}{M}\right]dr_1dr_2 \\ &= \int \{\rho_1^{(\rm p)}(r_1,r_1+s)\rho_2^{(\rm p)}(r_2,r_2-s)v_{\rm EX}^{(\rm pp)}(s;\rho,E/A) + \rho_1^{(\rm p)}(r_1,r_1+s)\rho_2^{(\rm n)}(r_2,r_2-s)v_{\rm EX}^{(\rm pp)}(s;\rho,E/A) \\ &+ \rho_1^{(\rm n)}(r_1,r_1+s)\rho_2^{(\rm p)}(r_2,r_2-s)v_{\rm EX}^{(\rm np)}(s;\rho,E/A) + \rho_1^{(\rm n)}(r_1,r_1+s)\rho_2^{(\rm n)}(r_2,r_2-s)v_{\rm EX}^{(\rm np)}(s;\rho,E/A) \\ &\times \exp\left[\frac{ik(R)\cdot s}{M}\right]dr_1dr_2, \end{split}$$

Globally-parameterized density ("Sao Paolo density")

L. C. Chamon, B. V. Carlson, L. R. Gasques, D. Pereira, C. D. Conti, M. A. Alvarez, M. S. Hussein, M. A. C. Ribeiro, E. S. Rossi, Jr., et al., Phys. Rev. C 66, 014610 (2001).

Globally-parameterized density ("Sao Paolo density") L.C.Chamon et al., PRC66, 014601 (2001)



FIG. 3. The R_0 parameter obtained for charge distributions extracted from electron scattering experiments and for theoretical densities obtained from Dirac-Hartree-Bogoliubov calculations.

Globally-parameterized density ("Sao Paolo density") L.C.Chamon et al., PRC66, 014601 (2001)



FIG. 2. Equivalent diffuseness values obtained for charge distributions extracted from electron scattering experiments and for theoretical densities obtained from Dirac-Hartree-Bogoliubov calculations.

Global parameterization of the CEG07 folding-model potentials

T. Furumoto, W. Horiuchi, M. Takashina, Y. Yamamoto, Y. Sakuragi (submitted to PRC, 2012)

✓ projectiles : Z = 6 (C isotope) ~ 20 (Ca isotope) (even-even)
 ✓ targets : ¹²C ~ ²⁰⁸Pb (closed or sub-closed shell nuclei)
 ✓ energy range : E/A = 30 ~ 400 MeV

$$V_F(R) = \sum_{n=1}^{10} \left\{ \alpha_n \exp\left(-\frac{R^2}{\gamma_n^2}\right) \right\},$$

$$W_F(R) = \sum_{n=1}^{10} \left\{ \beta_n \exp\left(-\frac{R^2}{\gamma_n^2}\right) \right\},$$

$$\alpha_n = \alpha_n (A_p, Z_p, A_t, E/A),$$

$$\beta_n = \beta_n (A_p, Z_p, A_t, E/A),$$

$$\gamma_n = 0.45 \left(\frac{n+8}{18}\right) (A_p^{1/3} + A_t^{1/3} + 1)$$











<u>summary</u>

- 1. Introduction
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Summary

- We have proposed a <u>new complex G-matrix</u> ("CEG07"),
 ✓ derived from <u>ESC04(extended soft-core) NN force</u>
 - ✓ include three-body force (TBF) effect
 - ✓ calculated up to higher density (about twice the normal density)
- We have applied DFM with <u>new complex G-matrix</u> ("CEG07") to nucleus-nucleus (AA) elastic/inelastic scattering & breakup
- CEG07 is successful for <u>nucleus-nucleus elastic scattering</u>
 - ✓ reproduce <u>cross section</u> data for ¹²C, ¹⁶O elastic scattering by ¹²C, ¹⁶O, ²⁸Si, ⁴⁰Ca targets at various energies.
 - ⇒ decisive role of Three-body <u>repulsive</u> force effect
 ✓ We also demonstrated possible applications to nuclear reactions (inelastic/breakup) including unstable nuclei
- The HI optical potential shows attractive-to-repulsive transition around E/A=200~300 MeV
- We constructed Global potentials for projectiles of unstable nuclei up to driplines, based on the microscopic CEG07 folding potentials.

A brief history of the double-folding model (DFM) study of HI optical potentials, before CEG07
Microscopic / semi-microscopic models :

✓ starting from <u>NN interactions</u> (V_{NN})

G-matrix with scattering b.c.

✓ VNN : effective NN interaction in nuclear medium

- should have proper density-dependence (ρ-dep) consistent with nuclear saturation properties
- ✓ should have proper energy-dependence (E-dep)
- ✓ should be complex (real-part + imaginary part)

However, no such ideal effective VNN exists so far !

<u>Simple M3Y</u> (1975~1985)

no density-dependence

real part only (add a phenom. imag. pot) zero-range exchange term \checkmark

$$v_{\rm NN}(\mathbf{r}) = 7999 \frac{e^{-4r}}{4r} - 2134 \frac{e^{2.5r}}{2.5r} - \hat{J}_{00}\delta(\mathbf{r})$$



Projectile

$$v_{\rm NN}(\mathbf{r}) = 7999 - \frac{1}{4r} - 2134 - \frac{1}{2.5r} - J_{00}\delta(\mathbf{r})$$

⇒ too deep at short distances, but gives

a reasonable strength at nuclear surface



Target

- due to strong absorption for Heavy lons (HI) ⇒ sensitive only to nuclear surface
- ⇒ "Successful" for low-energy (E/A<30 MeV) scattering</p> of heavy-ion (HI) projectiles with $A_p < 40$ [G.R.Satchler and W.G.Love, *Phys.Rep.55*, 183(1979)]



Introduction of density-dependence

DDM3Y-ZR (with <u>zero-range</u> exchange term)

$$v_{NN}(E,\rho;\mathbf{s}) = g(E,\mathbf{s})f(E,\rho)$$

 $f(E,\rho) = C(E) [1 + \alpha(E)e^{-\beta(E)\rho}]$

- ⇒ greatly reduce the potential strength at short distances
- ⇒ reproduce refractive phenomena, such as nuclear-rainbow (eg.⁴He+A, ¹⁶O+¹⁶O)

