Chiral Symmetry in Strongly Interacting Matter



- Spontaneously Broken Chiral Symmetry and Nuclear Chiral (Thermo-)Dynamics
- QCD interface with Nuclear Physics:

Chiral Effective Field Theory

Nuclear Equation of State

in the context of the QCD Phase Diagram

- Density and Temperature Dependence of the Chiral Condensate
- Astrophysical Constraints from Neutron Stars in Binaries



Prelude: PHASES and STRUCTURES of QCD



Phases and Symmetry Breaking Pattern

QCD Thermodynamics at zero baryon chemical potential











Low-energy / low temperature limit of QCD is realized as a Chiral Effective Field Theory of (weakly interacting) Nambu-Goldstone Bosons (PIONS)

CHIRAL EFFECTIVE FIELD THEORY

Gasser & Leutwyler Weinberg Ecker ... many others

• LOW-ENERGY QCD: Effective Field Theory of weakly interacting Nambu-Goldstone Bosons (PIONS) representing QCD at scales $Q << 4\pi f_{\pi} \sim 1 \, {
m GeV}$





Low-Energy Expansion: CHIRAL PERTURBATION THEORY







 \bigcirc

successfully applied to:

PION-PION scattering

PION-NUCLEON scattering

PION photoproduction and COMPTON scattering on the NUCLEON

Iong range NUCLEON-NUCLEON interaction

NUCLEAR MATTER and NUCLEI













Nucleon-Nucleon Interaction





Region I. Classical region, $r \leq 1.5\kappa^{-1}$, $(\kappa^{-1}$ is the pion Compton wave length) where the one-pion-exchange potential dominates and the quantitative behavior of the potential has been established.

Region II. Dynamical region, $0.7\kappa^{-1} \leq r \leq 1.5\kappa^{-1}$, where the two-pion-exchange potential competes with and exceeds the one-pion-exchange potential. The recoil effect is also appreciable in this region. The qualitative behavior, however, has been clarified.

Region III. *Phenomenological region*, $r \leq 0.7\kappa^{-1}$, where exist so many complicated effects, e.g., the relativistic effect, the isobar effect, the effect of new particles, etc., that at present we may have no means but some phenomenological treatment to fit with experiments.

M. Taketani, S. Nakamura, M. Sasaki Prog. Theor. Phys. **6** (1951) 581



NUCLEAR INTERACTIONS from **CHIRAL EFFECTIVE FIELD THEORY**

Weinberg

Bedaque & van Kolck Bernard, Epelbaum, Kaiser, Meißner; ...



Systematically organized HIERARCHY



NN Scattering Phase Shifts from **CHIRAL EFFECTIVE FIELD THEORY**



quantitatively accurate at same level of precision as best phenomenological potentials

CHIRAL EFFECTIVE FIELD THEORY

at work in nuclear few-body systems

example: elastic nd scattering



Explicit $\Delta(1230)$ DEGREES of FREEDOM

Large **spin-isospin polarizability** of the nucleon







N. Kaiser, S. Gerstendörfer, W.W., NPA637 (1998) 395 N. Kaiser, S. Fritsch, W.W., NPA750 (2005) 259

ΜΑΜΙ

(2001)

500

400

600

500

400

300

200

100

0

-100

-200

100

200

300

 \mathbf{E}_{γ} (MeV)

 $\sigma_{3/2}$ - $\sigma_{1/2}$ (µbarn)

C. Ordonez, L. Ray, U. van Kolck, PRL 72 (1994) 1982

$$\int_{N} \frac{\pi}{\pi} \int_{N} \frac{1}{\sqrt{N}} V_{c}(\mathbf{r}) = -\frac{9 g_{A}^{2}}{32\pi^{2} f_{\pi}^{2}} \beta_{\Delta} \frac{e^{-2\mathbf{m}_{\pi}\mathbf{r}}}{\mathbf{r}^{6}} P(\mathbf{m}_{\pi}\mathbf{r}) \qquad \longleftrightarrow \qquad \int_{N} \frac{\pi}{\pi} \int_{N} \frac{\pi}{\sqrt{N}} \frac{\pi}{\sqrt{N}} \frac{1}{\sqrt{N}} \frac{1$$

Explicit $\Delta(1230)$ **DEGREES of FREEDOM** (contd.)





Improved convergence



Important pieces of the CHIRAL NUCLEON-NUCLEON INTERACTION



N. Kaiser, S. Gerstendörfer, W.W.: Nucl. Phys. A 637 (1998) 395

CENTRAL ATTRACTION from TWO-PION EXCHANGE

• $\Delta(1232)$ • π •



Short distance:

NN POTENTIAL from LATTICE QCD

N. Ishii, S. Aoki, T. Hatsuda: Phys. Rev. Lett. 99 (2007) 022001



4 NUCLEAR MATTER and QCD PHASES



- momentum scale: Fermi momentum
- NN distance:
- energy per nucleon:
- compression modulus:

$$f k_F \simeq 1.4 \,\, fm^{-1} \sim 2m_\pi$$

 $f d_{NN} \simeq 1.8 \,\, fm \simeq 1.3 \,\, m_\pi^{-1}$
 $E/A \simeq -16 \,\, MeV$
 $K = (260 \pm 30) \,\, MeV \sim 2m_\pi$

CHIRAL DYNAMICS and the NUCLEAR MANY-BODY PROBLEM

N. Kaiser, S. Fritsch, W.W. (2002 - 2005)

$${\color{black} Small}\\ {\color{black} scales:} \qquad {\color{black} k_F}\sim 2\,m_\pi\sim M_\Delta-M_N<<4\pi\,f_\pi$$

• **PIONS** (and **DELTA** isobars) as **explicit** degrees of freedom





IN-MEDIUM CHIRAL PERTURBATION THEORY

In-medium nucleon propagator:

 $\frac{\mathbf{i}}{\gamma \cdot \mathbf{p} - \mathbf{M}_{\mathbf{N}} + \mathbf{i}\epsilon} - 2\pi(\gamma \cdot \mathbf{p} + \mathbf{M}_{\mathbf{N}})\delta(\mathbf{p}^2 - \mathbf{M}_{\mathbf{N}}^2)\theta(\mathbf{p}_0)\theta(\mathbf{k}_{\mathbf{F}} - |\vec{\mathbf{p}}|)$

- Loop expansion in ChPT \longleftrightarrow Systematic expansion of ENERGY DENSITY $\mathcal{E}(\mathbf{k}_{\mathbf{F}})$ in powers of Fermi momentum [modulo functions $\mathbf{f}_{\mathbf{n}}(\mathbf{k}_{\mathbf{F}}/\mathbf{m}_{\pi})$]
 - Finite nuclei \longleftrightarrow energy density functional
- Nuclear **thermodynamics**: compute **free energy density**



NUCLEAR MATTER



- $E_0/A = -16 \, MeV \ , \ \rho_0 = 0.16 \, fm^{-3} \ , \ K = 290 \, MeV$
- Realistic (complex, momentum dependent) single-particle potential ... satisfying Hugenholtz van Hove and Luttinger theorems (!)



Asymmetry energy $A(k_F^0) = 34 \, MeV$



Landau parameters



NUCLEAR THERMODYNAMICS



S. Fritsch, N. Kaiser, W.W.: Nucl. Phys. A 750 (2005) 259

PHASE DIAGRAM of NUCLEAR MATTER



PHASE DIAGRAM of NUCLEAR MATTER

Trajectory of **CRITICAL POINT** for **asymmetric matter**



... determined almost entirely by **isospin** dependent **pion** exchange dynamics











NUCLEAR MANY-BODY CALCULATIONS

... using NN and NNN interactions from Chiral Effective Field Theory

No-Core-Shell-Model results for ¹⁰B,¹¹B, ¹²C and ¹³C @ N²LO



Navratil et al., PRL 99 (2007) 042501



systematic improvements with inclusion of 3-body interactions





DENSITY FUNCTIONAL STRATEGIES

... constrained by (chiral) symmetry breaking pattern of Low-Energy QCD

Examples (part I)

Strategy :

 Calculate physics at long and intermediate distances using nuclear chiral effective field theory

Fix short distance constants (contact interactions) e.g. in Pb region

Predict systematics for all other nuclei





Examples (part II)





P. Finelli, N. Kaiser, D. Vretenar, W.W.: Nucl. Phys. A735 (2004) 449, A770 (2006) 1

Examples (part III):



... more applications

Gamow-Teller beta decays ¹⁴C

B-

interesting case:

anomalously long lifetime (5739 y) enables radiocarbon dating

Theoretically **not** understood on the basis of **two-nucleon** interactions only



Solution: chiral effective interaction including three-body force J.W. Holt, N. Kaiser, W.W.: Phys. Rev. C79 (2009) 054331, Phys. Rev. C81 (2010) 024002

Spin-orbit interactions



Role of 2nd order tensor force from pion exchange and three-body interactions

N. Kaiser: Phys. Rev. C68 (2003) 054001; N. Kaiser and W.W.: Nucl. Phys. A804 (2008) 60

In-medium Chiral SU(3) dynamics and **hypernuclei**

Weak Λ -nuclear spin-orbit coupling

N. Kaiser, W.W.: Phys. Rev. C71 (2005) 015203 P. Finelli, N. Kaiser, D.Vretenar, W.W.: Phys. Lett. B658 (2007) 90; Nucl. Phys. A831 (2009) 163





Anomalously long beta decay lifetime of $^{14}\mathrm{C}$

Early history: B. Janovici, I. Talmi : Phys. Rev. 95 (1954) 289 (Role of **tensor force**)



Known lifetime of 5730 years enables radiocarbon dating

 But theoretical description using realistic nucleon-nucleon interactions overestimates the GT strength

Idea: Derive a density-dependent two-nucleon force from the leading-order chiral three-nucleon force



Large suppression of GT strength at p₀ due to chiral 3NF!

J.W. Holt, N. Kaiser, W.W.: Phys. Rev. C79 (2009) 054331, Phys. Rev. C81 (2010) 024002



6 CHIRAL CONDENSATE at finite DENSITY

 $(\mathbf{T} = \mathbf{0})$

Hellmann - Feynman theorem: $\langle \Psi | \bar{\mathbf{q}} \mathbf{q} | \Psi \rangle = \langle \Psi | \frac{\partial \mathcal{H}_{\mathbf{QCD}}}{\partial \mathbf{m}_{\mathbf{q}}} | \Psi \rangle = \frac{\partial \mathcal{E}(\mathbf{m}_{\mathbf{q}}; \rho)}{\partial \mathbf{m}_{\mathbf{q}}}$





CHIRAL CONDENSATE: DENSITY DEPENDENCE

Symmetric Nuclear Matter



Substantial **change** of **symmetry breaking scenario** between chiral limit $m_q=0$ and physical quark mass $m_q\sim 5\,{
m MeV}$

Nuclear Physics would be very different in the chiral limit !



No indication of first order chiral phase transition in the range $ho \leq 2
ho_0, ~~{f T} \leq 100~{f MeV}$

CHIRAL CONDENSATE: DENSITY DEPENDENCE Neutron Matter



- Qualitative difference between **nuclear** and **neutron** matter
- Chiral limit: logarithmic singularity as $\, {f m}_{\pi}
 ightarrow 0$
- Important: realistic treatment of two-body and three-body correlations in extrapolations to high-density matter





identify leading collective degrees of freedom (→ order parameters)

quarks as **quasiparticles** with dynamically generated masses



"wrong" degrees of freedom at low temperature, non-zero baryon density?

PHASE DIAGRAM (contd.)

Issues: Critical Point(s)

Diquarks and Color Super Conducting Phases



Chemical Potential μ

M. Tachibana, N. Yamamoto, T. Hatsuda, G. Baym (2006-09)





PHASE DIAGRAM (contd.)

PNJL model calculations

N. Bratovich, T. Hell, W.W. (2010)



Quarks **cannot** be the relevant **active quasiparticles** at low temperatures and baryon chemical potentials $\mu_{
m B} \sim 1~{
m GeV}$

PHASE DIAGRAM (contd.)



Major challenge: design QCD phase diagram in accordance with known realistic constraints from hadronic and nuclear physics



8 Constraints from Astrophysical Observations

Constraints on **Equation of State** at **HIGH DENSITY**

(2 - 10 times the density of normal nuclear matter)







Supernovae

to

Neutron Stars







NEUTRON STARS and the EQUATION OF STATE of DENSE BARYONIC MATTER

J. Lattimer, M. Prakash: Astrophys. J. 550 (2001) 426 Phys. Reports 442 (2007) 109

Mass-Radius Relation



Oulook: New Constraints from NEUTRON STARS in BINARIES

... using additional information from observables such as apparent surface and flux during cooling phase of burst



Implications for the EQUATION of STATE









Summary and Conclusions

- Interface of Low-Energy QCD and Nuclear Physics: Nuclear Chiral (Thermo-) Dynamics really works !
- Importance of Two-Pion Exchange in-medium processes in combination with Pauli principle
- Three-Nucleon Forces are natural part of nuclear chiral dynamics
 → important stabilizer of nuclear matter at higher density
 - Magnitude of **Chiral Condensate** remains non-zero in the range

 $ho \leq 2
ho_0, \ \mathbf{T} \leq \mathbf{100} \ \mathbf{MeV}$

→ corridor of **Spontaneous Chiral Symmetry Breaking** persists along baryon density axis

